

Summer school on Quantum Chaos
Sponsored by IAMP, EMS & ESI
Introduction: Zeev Rudnick

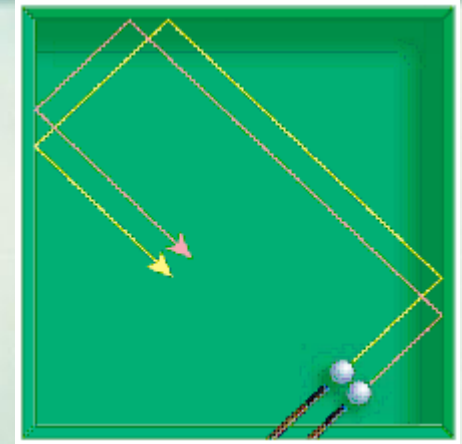
What is Quantum Chaos ?

- “Quantum chaos” = study of energy levels and stationary states of quantum systems in the semiclassical limit $\hbar \rightarrow 0$.
- This week we will explore some aspects of subject.

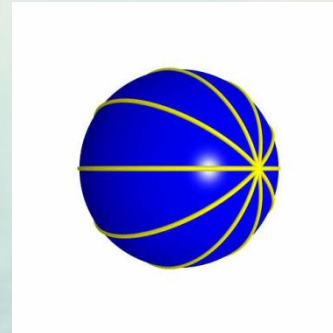
Models of classical mechanics

planar billiards:

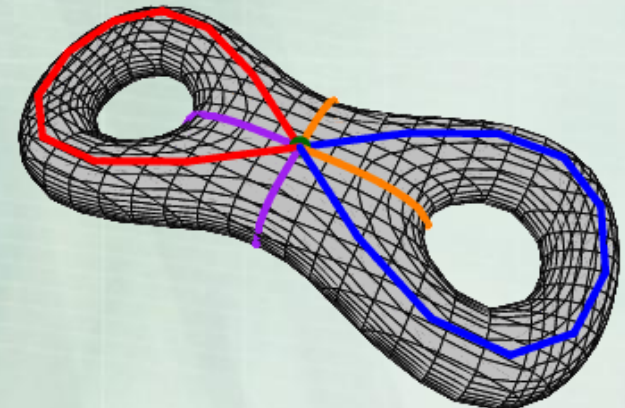
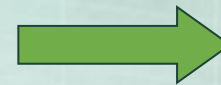
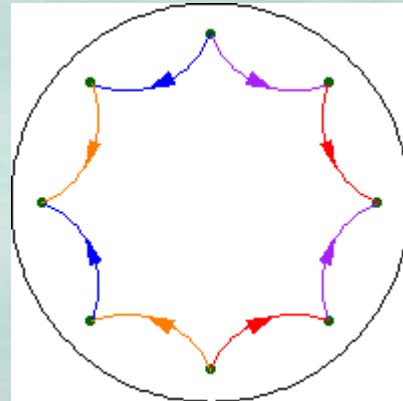
angle of reflection = angle of incidence



Geodesic flow on a surface



Geodesic flow on negatively curved surface



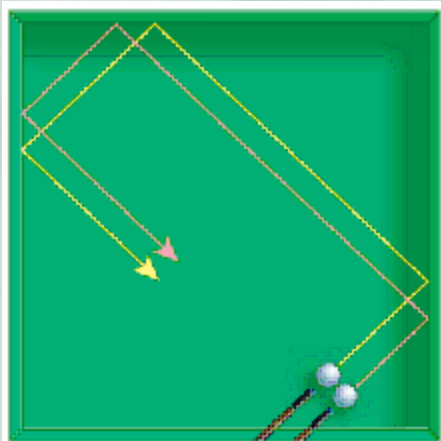
Regularity vs. chaos in classical dynamics

see courses by Nonnenmacher, Riviere

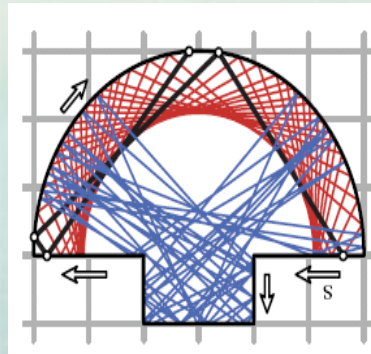
Classification of (conservative, Hamiltonian) dynamical systems, e.g; planar billiards, geodesic flows on surfaces

Regular (integrable):

- A full set of constants of motion
- dynamics confined to invariant tori in phase space.
- Linear separation of trajectories



...mixed ...



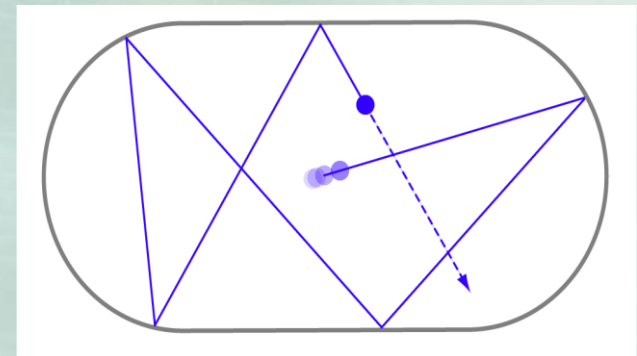
... pseudo-integrable



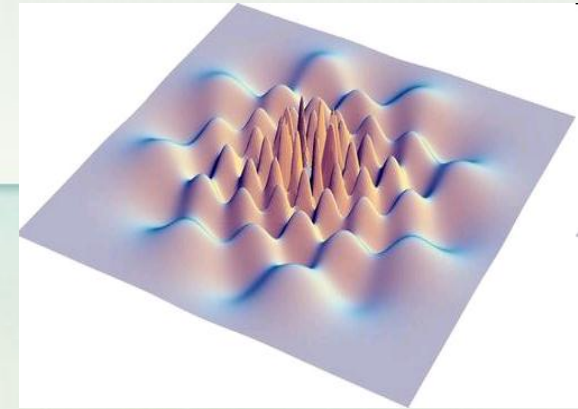
Rudnick

Chaotic:

- Typical orbits densely cover all of available phase space (ergodicity)
- Exponential divergence of nearby trajectories (hyperbolicity).



Quantum mechanics



A particle at time t is described by its wave function $\Psi(q,t)$

$|\Psi(q,t)|^2$ = probability density of particle in state Ψ

Time evolution is described by **Schrödinger's equation** :

$$i\hbar \frac{\partial \Psi}{\partial t} = -\frac{\hbar^2}{2} \Delta \Psi \quad \hbar = 1.054 \times 10^{-34} \text{ J-s}$$

Stationary states: $\Psi(q,t) = \psi(q)e^{-itE/\hbar}$ with $\psi(q)$ an eigenfunction of Δ

$$-\frac{\hbar^2}{2} \Delta \psi = \mathbf{E} \psi \quad \mathbf{E} = \text{energy level}$$

The semiclassical limit $\hbar \rightarrow 0$ & the correspondence principle:

“classical mechanics is a special case of quantum mechanics”.

If so, then:

How is the dichotomy “regular vs. chaotic” manifested in Quantum Mechanics ?

Statistics of eigenfunctions

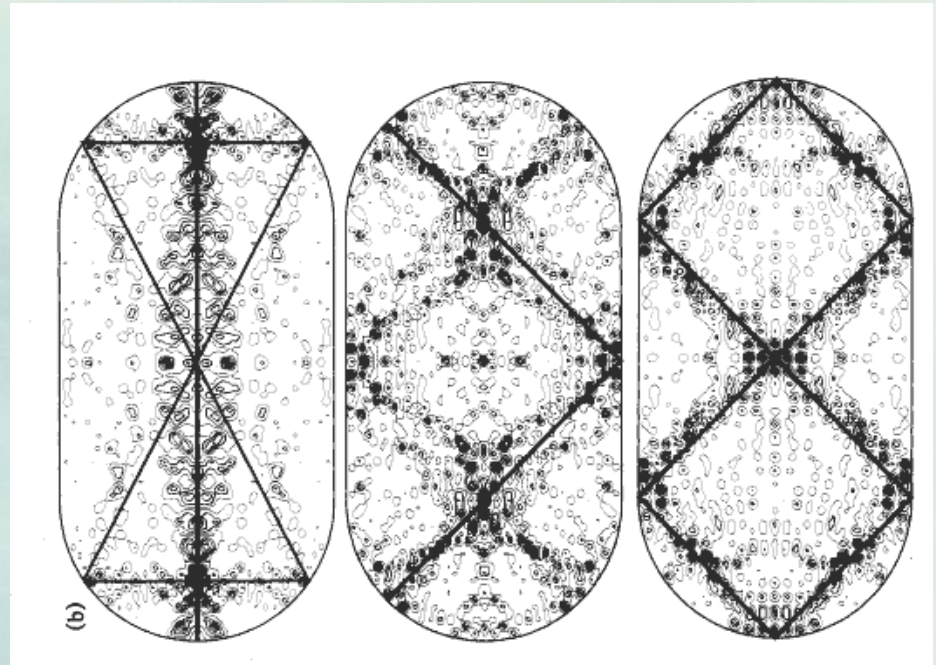
The semi-classical eigenfunction hypothesis of M.V. Berry and A. Voros (~ 1977):

“Each semi-classical eigenstate has a Wigner function concentrated on the region explored by a typical orbit over infinite times”.

In particular, for chaotic systems, “The wave functions cover phase space uniformly”

However.....“scars” were found by Heller and by McDonald & Kauffman (1984-88)

Scars: Concentration of eigenfunctions on unstable periodic orbits (controversial)



E. Heller: Scarred stadium mode

A mathematical formulation

see Hassel's course

For a particle with wave function ψ , the expectation values of its position coordinates q_1 , q_2 are given by

$$\langle q_1 \rangle_\psi := \int q_1 |\psi(q_1, q_2)|^2 dq_1 dq_2$$

“Likewise”, for any classical observable $a(q,p)$ of position $q=(q_1, q_2)$ and momentum $p=(p_1, p_2)$, one can define a (pseudo-differential) operator $Op(a)$ so that the expected value of the observable a “at the state ψ ” is the diagonal matrix element $\langle Op(a)\psi, \psi \rangle$

A possible interpretation of the statement that “wave functions cover phase space uniformly” is that the matrix elements converge to the classical average of a :

$$\langle Op(a)\psi, \psi \rangle \xrightarrow{E_\psi \rightarrow \infty} \iint a(q, p) dp dq$$

Quantum Ergodicity (Schnirelman)

Schnirelman (1974): For a Riemannian manifold M with ergodic geodesic flow, “most” eigenfunctions cover phase space uniformly: if u_n is an ONB consisting of eigenfunctions of the Laplacian, then there is a subsequence of density one s.t. for all observables $a(p,q)$

$$\langle Op(a)u_n, u_n \rangle \xrightarrow{n \rightarrow \infty} \int_{S^*M} a$$

Zelditch (1987), Colin de Verdiere (1985)

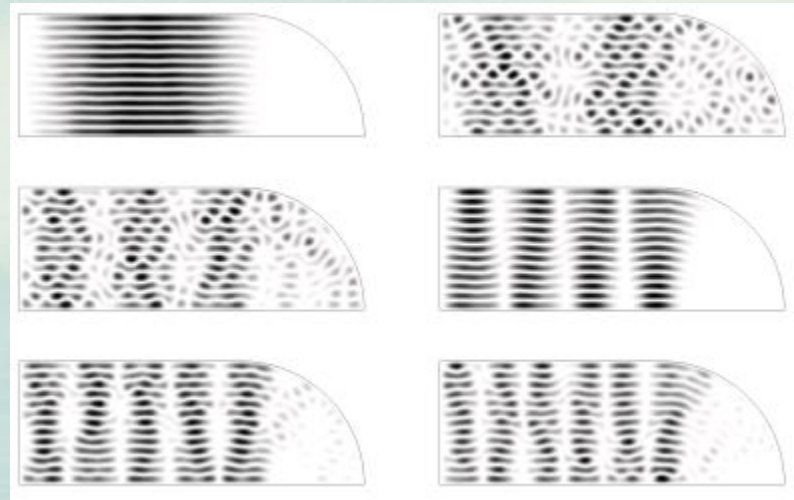
see Hassel’s course

One interpretation of “Scars” is as possible exceptional subsequences

ZR & Sarnak (1994): Conjecture that for negatively curved manifolds, no exceptional subsequence - **Quantum Unique Ergodicity (QUE)**.

see Einsiedler’s course

Bouncing ball modes in the stadium billiard



Hassel (2010): BBM exist for “most” stadia

Nodal lines of eigenfunctions

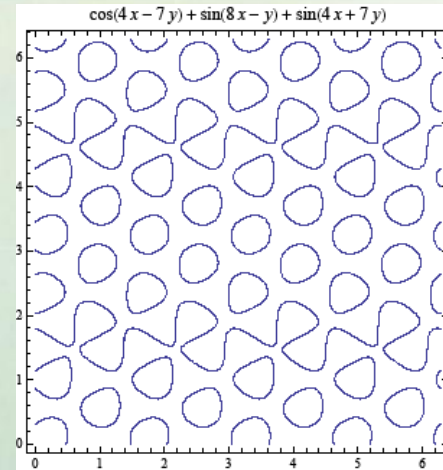
The **nodal line** of φ is the set $\varphi^{-1}(0)=\{x: \varphi(x)=0\}$.

Question: How do nodal lines of eigenfunctions vary as we increase the eigenvalue ?

see Zelditch's talk

Random wave models –used to predict features such as the statistics of the number of nodal domains of eigenfunctions

see Sodin's course



nodal lines of eigenfunctions on the flat torus

Statistics of the eigenvalues

A major insight of Quantum Chaos the statistics of the energy spectrum falls into a few universality classes, described by simple statistical models, depending on the coarse classification of the dynamics of the classical limit of the system.

One popular statistical measure is the level spacing distribution $P(s)$:= limiting distribution of the normalized gaps δ_j between adjacent levels

$$\delta_n := \frac{E_{n+1} - E_n}{\text{mean spacing}}$$


The diagram shows a horizontal axis with six energy levels labeled $E_1, E_2, E_3, E_4, E_5, E_6$ from left to right. Each level is marked with a green star. Below the axis, double-headed arrows indicate the gaps between adjacent levels: E_1 to E_2 , E_2 to E_3 , E_3 to E_4 , E_4 to E_5 , and E_5 to E_6 . The gaps are of varying widths, with the largest gap between E_5 and E_6 .

$$\frac{1}{N} \# \{n \leq N : \delta_n < x\} \xrightarrow{N \rightarrow \infty} \int_0^x P(s) ds$$

Statistical models: Poisson vs RMT

Statistics of the energy levels in the local regime can be compared to some simple probabilistic models:

a) Uncorrelated levels: Take the levels E_n to be independent, uniform in $[0,1]$ (homogeneous Poisson process on the line with intensity 1).

Here the level spacing distribution $P(s)=\exp(-s)$

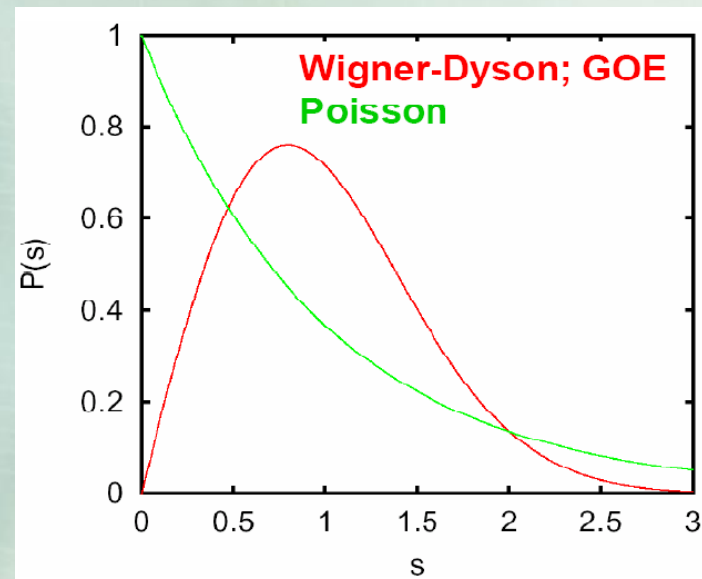
b) E_n are the eigenvalues of a **random** symmetric matrix (GOE)

$N \times N$ symmetric matrices $H=H^T$,

matrix elements=independent real Gaussians

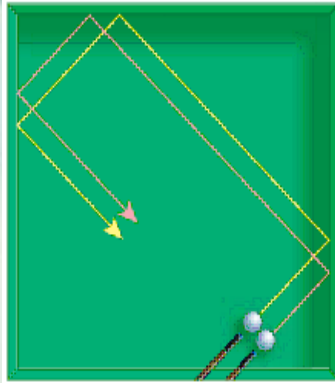
$P(s)$ was computed by Gaudin and Mehta (1960's)

see Keating's course



Integrable vs. chaotic: universality conjectures for the level spacing distribution $P(s)$

Berry & Tabor 1977: For integrable dynamics expect* Poisson statistics



see Marklof's talk

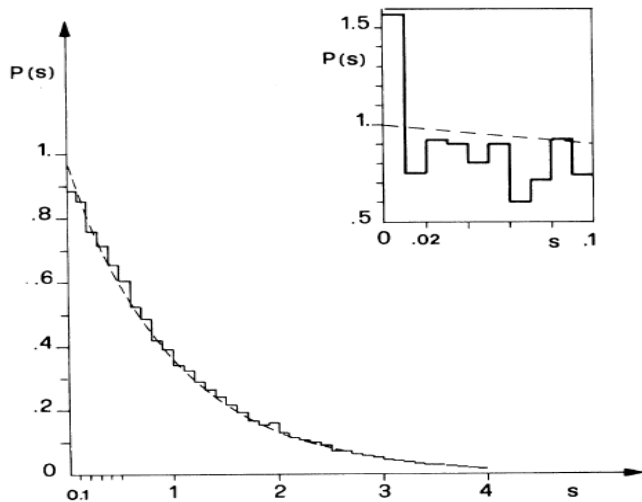
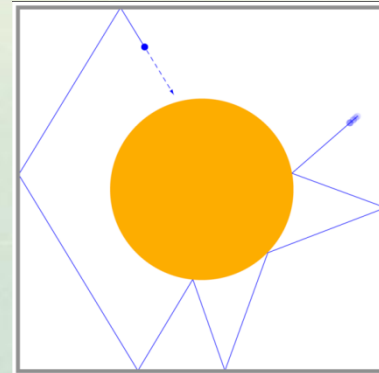


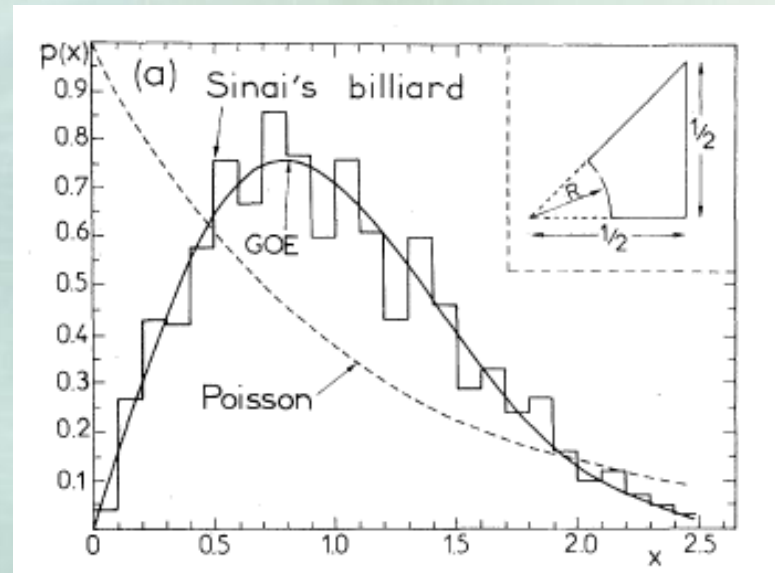
FIG. 1. Level-spacing distribution obtained from the first 100 000 levels (1) with $\alpha = \pi/3$. The dotted line is the Poisson distribution $P(s) = e^{-s}$.

rectangular billiard, aspect ratio $= \sqrt{\pi/3}$

Bohigas, Giannoni & Schmit (1984): for chaotic dynamics expect* GOE statistics



see Keating's course, Muller's talk



*="generically"

Have a good week!

