Solution and properties of the resummed BK equation and KGBJS equation

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Krzysztof Kutak, D. T., arXiv:1306.3369 [hep-ph]

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$$\begin{aligned} \mathcal{A}(x,k^2,p) &= \mathcal{A}_0(x,k^2,p) \\ &+ \int \frac{\mathrm{d}^2 \bar{\mathbf{q}}}{\pi \bar{q}^2} \int_{x/x_0}^{1-Q_0/\bar{q}} \mathrm{d}z \,\theta(p-z\bar{q}) P_{gg}(z,k,\bar{q},p) \mathcal{A}\left(\frac{x}{z},{k'}^2,\bar{q}\right) \end{aligned}$$

$$P_{gg}(z,k,\bar{q},p) = \bar{\alpha}_s \Delta_s(p,z\bar{q}) \left(\frac{\Delta_{ns}(z,k,\bar{q})}{z} + \frac{1}{1-z}\right)$$

CCFM evolution – angular ordering



$$\mathcal{A}(x,k^2,p) = \mathcal{A}_0(x,k^2,p) + \int \frac{\mathrm{d}^2 \bar{\mathbf{q}}}{\pi \bar{q}^2} \int_{x/x_0}^{1-Q_0/\bar{q}} \mathrm{d}z \,\theta(p-z\bar{q}) P_{gg}(z,k,\bar{q},p) \mathcal{A}\left(\frac{x}{z},{k'}^2,\bar{q}\right)$$

 $\bar{\boldsymbol{\xi}} > \xi_{i+1} > \xi_i$

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CCFM evolution – low x limit



$$\mathcal{A}(x,k^2,p) = \mathcal{A}_0(x,k^2,p) + \bar{\alpha}_s \int \frac{\mathrm{d}^2 \bar{\mathbf{q}}}{\pi \bar{q}^2} \int_{x/x_0}^1 \frac{\mathrm{d}z}{z} \,\theta(p-z\bar{q})\Delta_{ns}(z,k,\bar{q})\mathcal{A}\left(\frac{x}{z},k^{\prime 2},\bar{q}\right)$$

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High energy factorization (k_T -factorization) applied to DIS

Working with BFKL/BK:

$$F_2(x,Q^2) = \sum_q e_q^2 \int \frac{\mathrm{d}k}{k} \sigma_{F_2} \left(k^2, Q^2, m_q^2\right) \mathcal{F}(x,k^2)$$

Working with CCFM/KGBJS:

$$F_2(x,Q^2) = \sum_q e_q^2 \int \frac{\mathrm{d}k}{k} \int_x^1 \frac{\mathrm{d}z}{z} \hat{\sigma}_{F_2}\left(z,k^2,Q^2,m_q^2\right) \mathcal{A}(\frac{x}{z},k,p)$$

$$p^2 = \frac{Q^2}{z(1-z)}$$

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Reformulated Balitsky-Kovchegov evolution equation

A new form introduces a new scale μ to sum low-q emissions:

$$\begin{split} \phi(x,k) &= \tilde{\phi}_0(x,k) \\ &+ \bar{\alpha}_s \int_x^1 \frac{\mathrm{d}z}{z} \int \frac{\mathrm{d}^2 q}{\pi q^2} \theta(q^2 - \mu^2) \Delta_R(z,k,\mu) \\ &\left(\phi(\frac{x}{z}, |k+q|^2) - \frac{q^2}{\pi R^2} \delta(q^2 - k^2) \phi^2(\frac{x}{z},q^2) \right) \end{split}$$

with

$$\Delta_R(z,k,\mu) = \exp\left(-\bar{\alpha}_s \log\frac{1}{z}\log\frac{k^2}{\mu^2}\right)$$

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Kutak, Golec-Biernat, Jadach, Skrzypek, JHEP 1202 (2012) 117

Reformulated BK – sensitivity to μ



The new form works down to $k = \mu$.

Reformulated BK – consistency check



BK evolution – nonlinear behavior

Following Kutak, Stasto, Eur. Phys. J. C 41 (2005) 343, define:



 $\ln Q_s$ grows linearly with $C \ln \frac{1}{x}$

BK vs KGBJS

The BK equation:

$$\begin{split} \phi(x,k) &= \tilde{\phi}_0(x,k) \\ &+ \bar{\alpha}_s \int_{x/x_0}^1 \frac{\mathrm{d}z}{z} \int \frac{\mathrm{d}^2 \mathbf{q}}{\pi q^2} \theta(q^2 - \mu^2) \Delta_R(z,k,\mu) \\ &\left(\phi(\frac{x}{z},k'^2) - \frac{q^2}{\pi R^2} \delta(q^2 - k^2) \phi^2(\frac{x}{z},q^2)\right) \end{split}$$

promoted to a nonlinear CCFM-based equation:

$$\begin{aligned} \mathcal{E}(x,k^2,p) &= \mathcal{E}_0(x,k^2,p) \\ &+ \bar{\alpha}_s \int_{x/x_0}^1 \frac{\mathrm{d}z}{z} \int \frac{\mathrm{d}^2 \bar{\mathbf{q}}}{\pi \bar{q}^2} \,\theta(p-z\bar{q}) \Delta_{ns}(z,k,\bar{q}) \\ &\left(\mathcal{E}\left(\frac{x}{z},k'^2,\bar{q}\right) - \frac{\bar{q}^2}{\pi R^2} \delta(\bar{q}^2-k^2) \,\mathcal{E}^2(\frac{x}{z},\bar{q}^2,\bar{q}) \right) \end{aligned}$$

Kutak, Golec-Biernat, Jadach, Skrzypek, JHEP 02 (2012) 117 Kutak, JHEP 12 (2012) 33

KGBJS – solution



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Krzysztof Kutak, D. T., arXiv:1306.3369 [hep-ph]

KGBJS – solution

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KGBJS - solution



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Solutions of KGBJS and BK

For $p \to \infty$, KGBJS becomes similar to BK.

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Solutions of KGBJS and BK

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KGBJS - the effect of running coupling

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KGBJS – the effect of running coupling

KGBJS - limits of linear behavior

Contour lines indicate shape of the $Q_s(x, p)$ dependecy:

$$\beta(x, Q_s(x, p), p) = const.$$

Qualitatively similar effect to the saturation lines observed by Avsar, Stasto, JHEP 1006 (2010) 112.

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KGBJS - limits of linear behavior

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KGBJS - limits of linear behavior

For large p, the growth of the saturation scale is "liberated".

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KGBJS – limits of linear behavior

One may also consider the $P_s(x,k)$ dependecy:

$$\beta\left(x,k,P_s(x,k)\right) = const$$

Note that the quadratic term of KGBJS is:

$$-\frac{1}{\pi R^2} \int_x^{x_0} \frac{\mathrm{d}w}{w} \theta(p-zk) P_{gg}(z,k,k) \mathcal{E}^2(w,k,k)$$

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Thank you.

Future work:

 \blacktriangleright include Δ_s and full splitting function in the calculations

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- test variants of kinematical constraints
- evolve high-energy-factorizable gluon density directly
- study phenomenological implications

The Balitsky-Kovchegov evolution equation

Integral form in momentum space:

$$\begin{split} \phi(x,k) &= \phi_0(x,k) \\ &+ \bar{\alpha}_s \int_x^1 \frac{\mathrm{d}z}{z} \int_0^\infty \frac{\mathrm{d}l^2}{l^2} \left(\frac{l^2 \phi(\frac{x}{z},l^2) - k^2 \phi(\frac{x}{z},k^2)}{|k^2 - l^2|} + \frac{k^2 \phi(\frac{x}{z},k)}{\sqrt{4l^4 + k^4}} \right) \\ &- \bar{\alpha}_s \int_x^1 \frac{\mathrm{d}z}{z} \phi^2(\frac{x}{z},k) \end{split}$$

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CCFM evolution – Δ_{ns}

$$\begin{split} \mathcal{A}(x,k^2,p) &= \mathcal{A}_0(x,k^2,p) \\ &+ \bar{\alpha}_s \int \frac{\mathrm{d}^2 \bar{\mathbf{q}}}{\pi \bar{q}^2} \int_{x/x_0}^1 \frac{\mathrm{d}z}{z} \,\theta(p-z\bar{q}) \Delta_{ns}(z,k,\bar{q}) \mathcal{A}\left(\frac{x}{z},k^{'2},\bar{q}\right) \\ \Delta_{ns} &= \exp\left(-\bar{\alpha}_s \ln\frac{1}{z} \ln\frac{k^2}{zq^2}\right) \quad \text{for} \quad k^2 > zq^2 \end{split}$$

In this work, calculations were done with Δ_{ns} after Kwiecinski, Martin, Sutton, PRD 52 (1995) 1445:

$$\Delta_{ns}(z,k,q) = \exp\left(-\bar{\alpha}_s \ln \frac{z_0}{z} \ln \frac{k^2}{z_0 z q^2}\right)$$

with

$$z_0 = \begin{cases} z & \text{ for } \frac{k}{q} < z \\ \frac{k}{q} & \text{ otherwise} \\ 1 & \text{ for } \frac{k}{q} > 1 \end{cases}$$