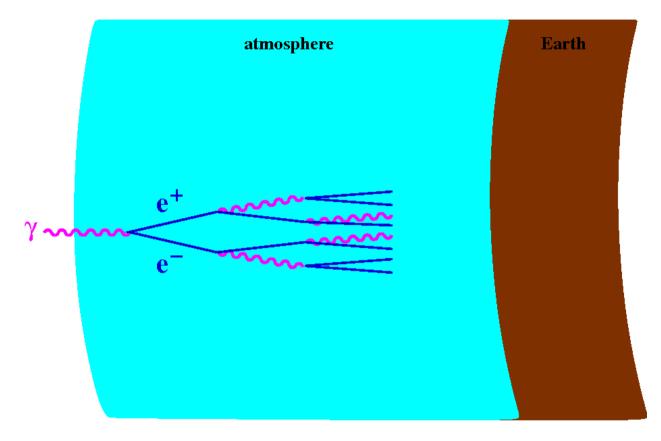
The Landau-Pomeranchuk-Migdal Effect 1953 – today

Cosmic Rays to Quark-Gluon Plasmas to String Theory

Peter Arnold University of Virginia

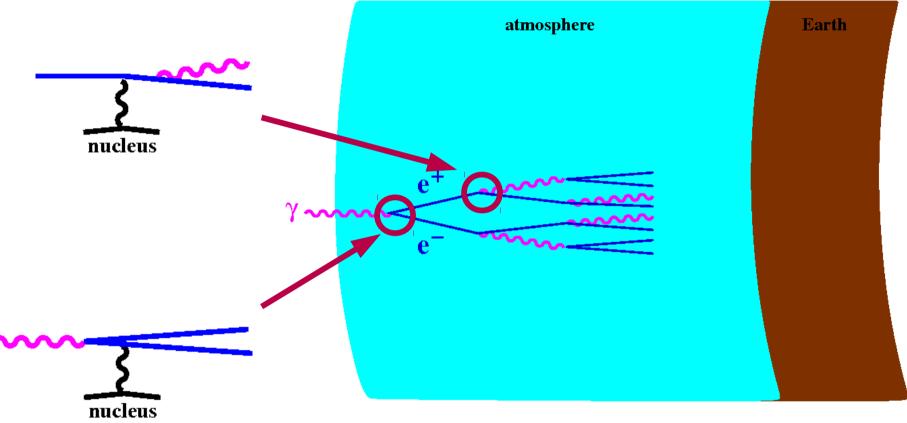
Cosmic Ray Showering

High energy particles traveling through matter lose energy via successive bremsstrahlung and pair production:



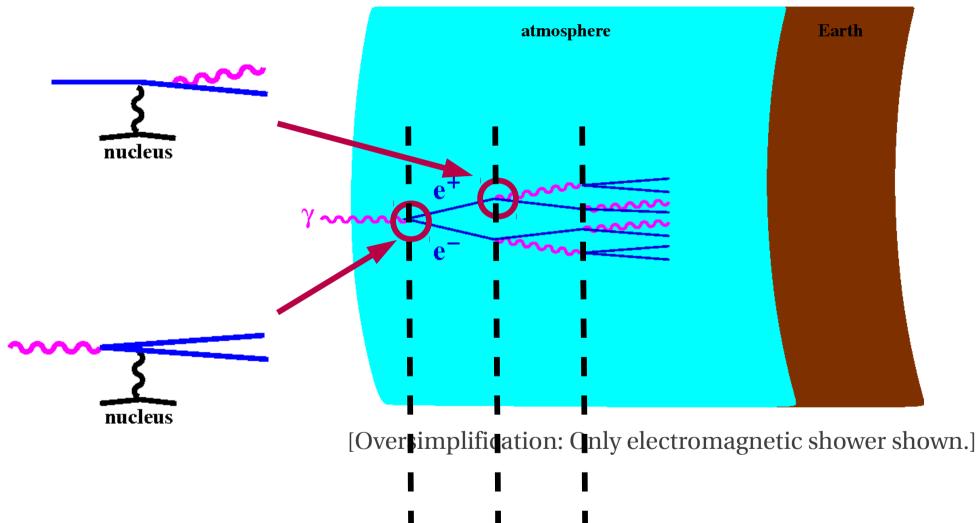
[Oversimplification: Only electromagnetic shower shown.]

High energy particles traveling through matter lose energy via successive bremsstrahlung and pair production:



[Oversimplification: Only electromagnetic shower shown.]

High energy particles traveling through matter lose energy via successive bremsstrahlung and pair production:





Landau-Pomeranchuk-Migdal (LPM) effect^{2/27}

What is the LPM Effect?

A coherence effect that complicates calculations of bremsstrahlung or pair production when a very high energy particle scatters from a medium.

Places it comes up in QED

- Very high energy cosmic rays showering in the atmosphere.
- Certain beam dump experiments designed to measure the LPM effect.

Places it comes up in QCD

• Energy loss of high energy jets in a quark-gluon plasma.

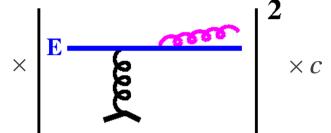


- Complete leading-order calculations of the viscosity and other transport coefficients of a weakly-coupled quark-gluon plasma.
- Coherence in high-energy, small-*x* scattering from a large nucleus: color glass condensate.

The LPM Effect



brem rate ~ $n\sigma v$ ~ (density of scatterers) ×



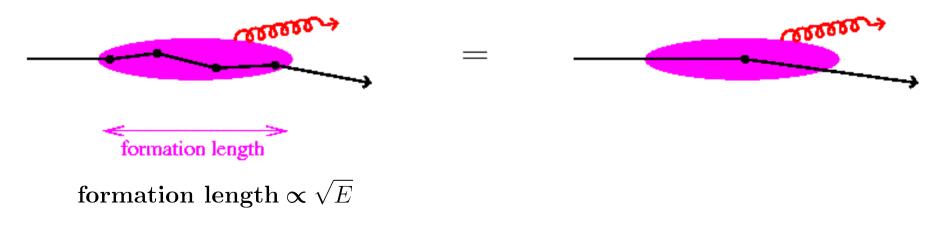
Result: high-energy brem rate independent of *E*.

<u>Problem</u>

At very high energy,

probabilities of brem from successive scatterings no longer independent;

brem from several successive (small angle) collisions not very different from brem from one collision.



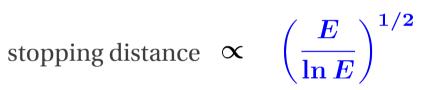
Result: a reduction of the naive brem rate.

Example: stopping distance (in a infinite medium)

If LPM effect ignored:

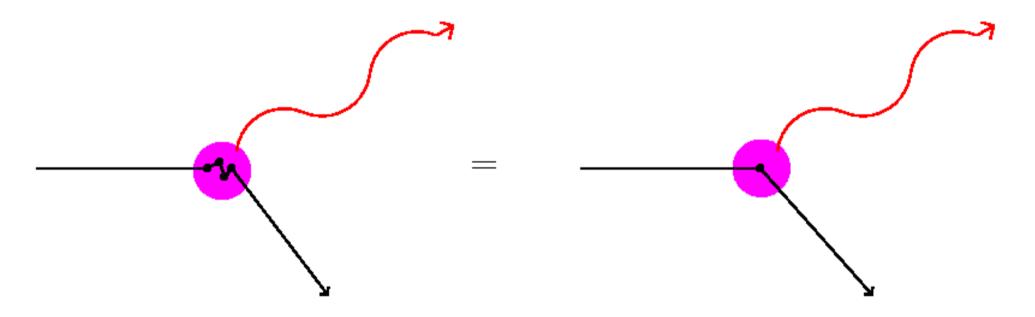
stopping distance $\propto -\ln E$

Actual result (weak coupling):



The LPM Effect (QED)

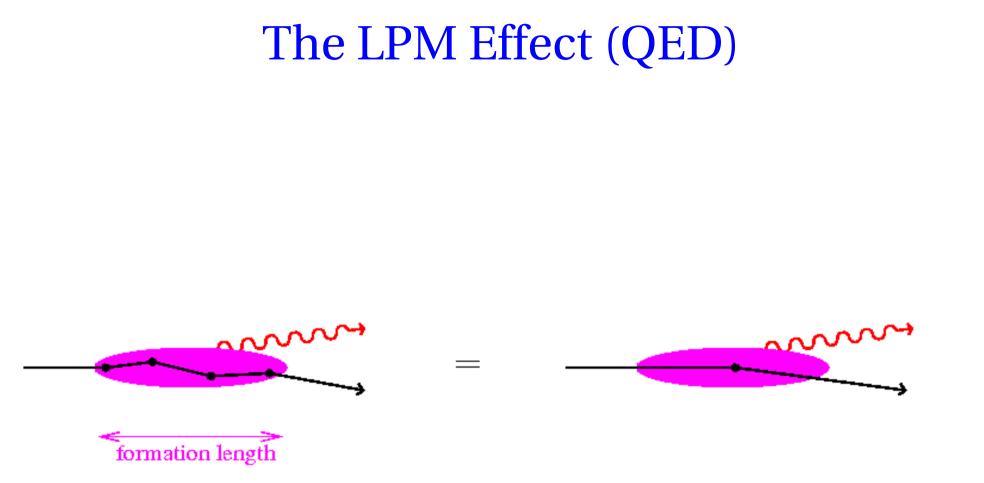
<u>*Warm-up*</u>: Recall that light cannot resolve details smaller than its wavelength.



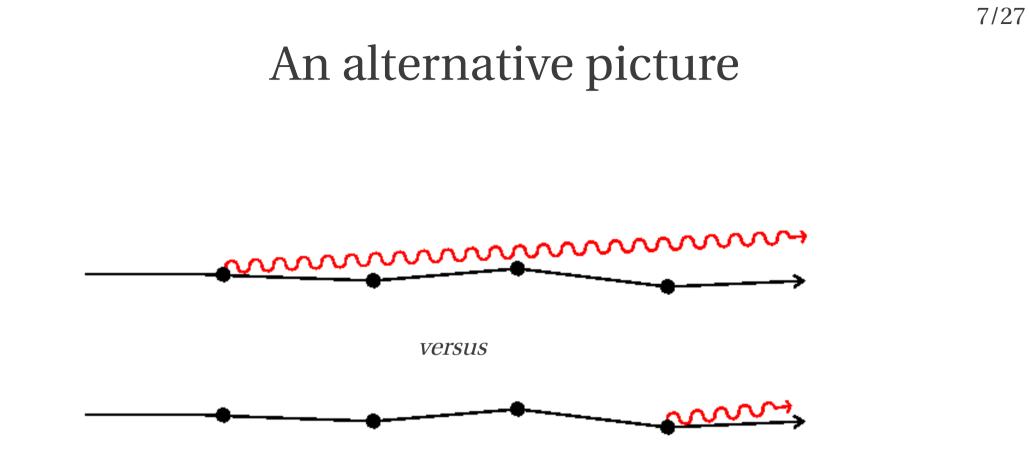
[Photon emission from different scatterings have same phase \rightarrow coherent.]

Now: Just Lorentz boost above picture by a lot!





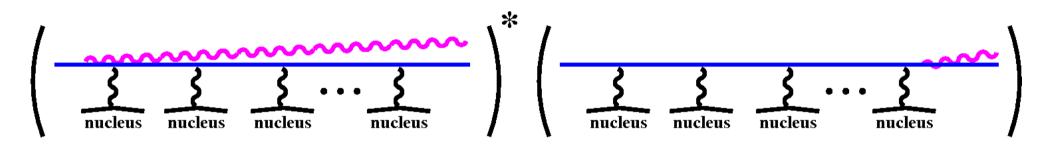
Note: (1) **bigger** *E* requires bigger boost \rightarrow more time dilation \rightarrow **longer formation length** (2) big boost \rightarrow this process is **very collinear**.



Are these two possibilities in phase? Or does the interference average to zero?

IN PHASE if (i) everything is nearly collinear	\checkmark
(ii) particle and photon have nearly same velocity	✓ (speed of light)

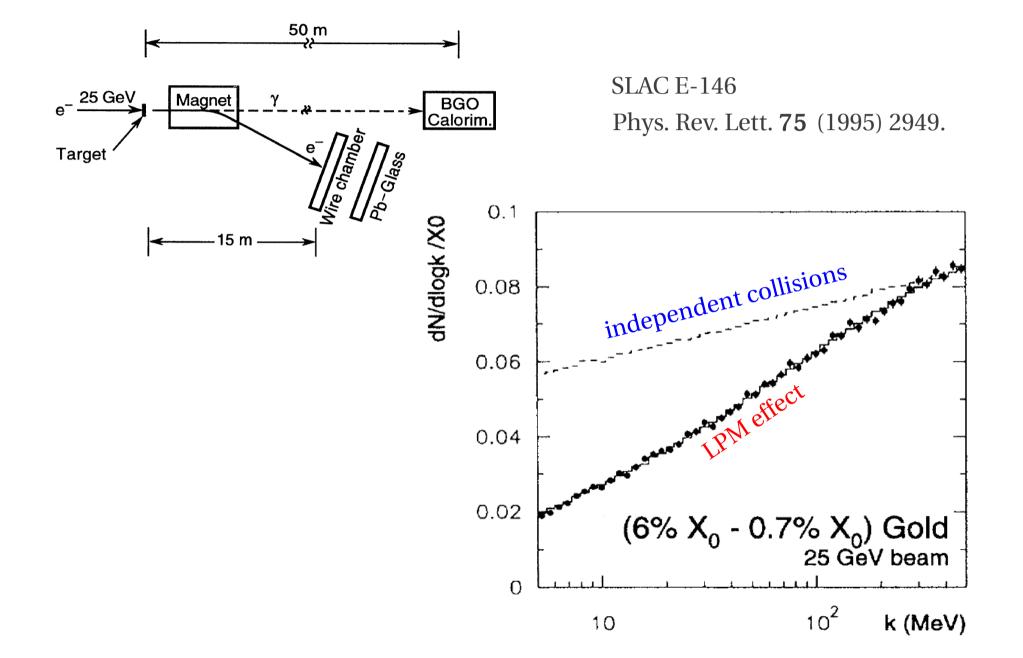
A real calculation involves...



soft photons: Can do a classical EM calculation hard photons: Much trickier!

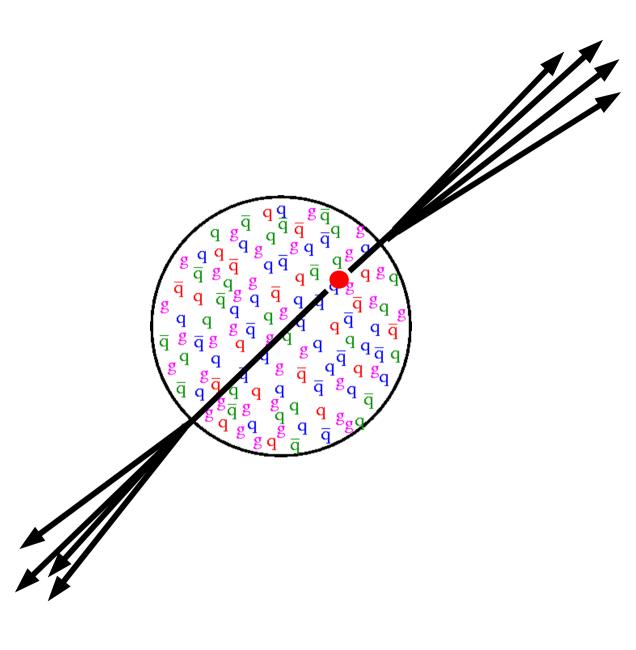
(Landau+Pomeranchuk 1953) (Migdal 1955)

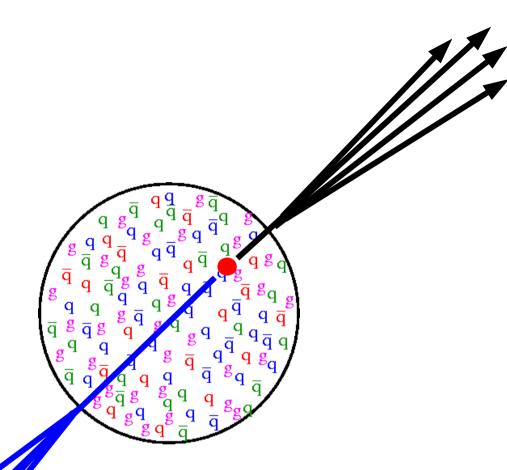
Experimental Measurement of LPM (QED)



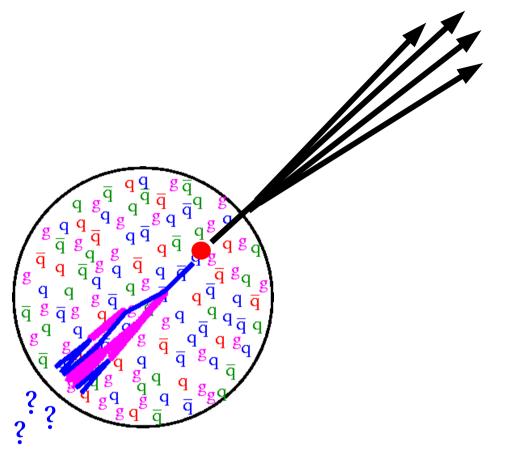
<u>Jet Quenching in</u> <u>Quark-Gluon Plasmas</u>



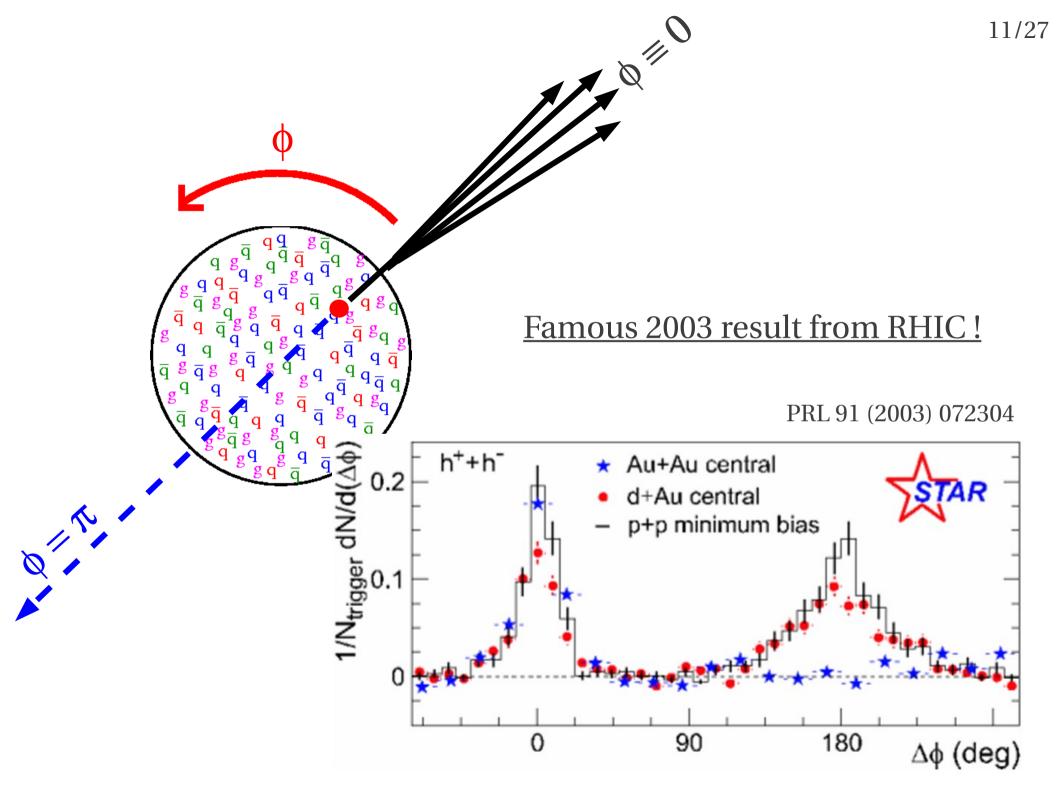


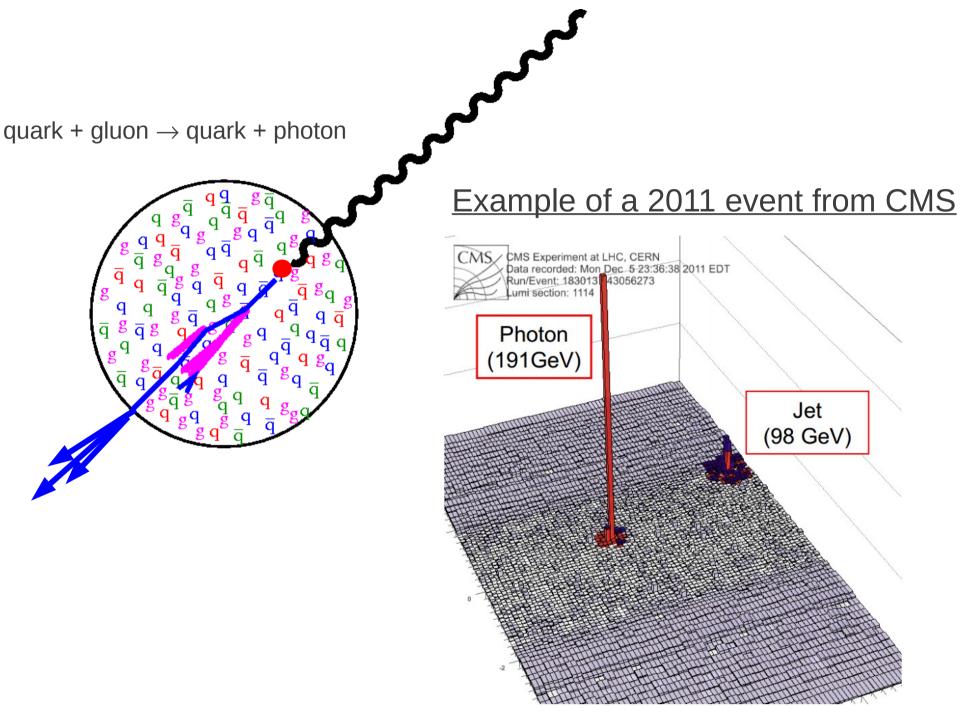


This one had to go a long way! It might shower and perhaps even stop <u>inside</u> the medium.



This one had to go a long way! It might shower and perhaps even stop <u>inside</u> the medium.

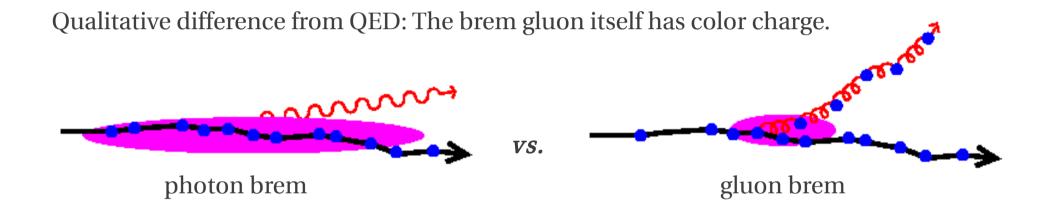




[from CMS highlights talk at Quark Matter 2012]

The LPM Effect (QCD)

Basic formalism worked out in the 1990s by BDMPS-Z (with ongoing caveats about how to apply it to realistic experimental situations).

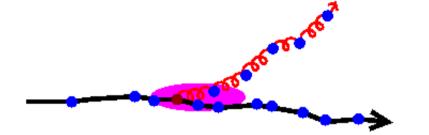


Parametrically, still gives rise to a rate suppression that grows like $E^{1/2}$.

stopping distance $\propto E^{1/2}$

One of the Caveats

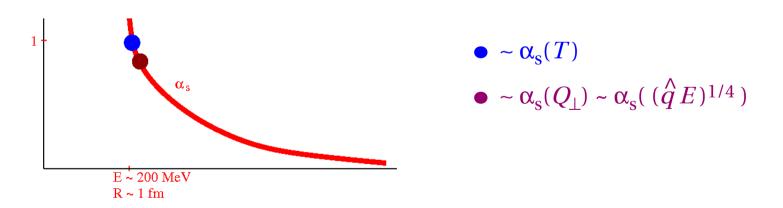




implicitly assumes that it makes sense to talk about individual quarks and gluons.

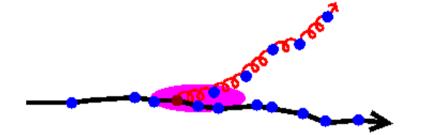
Talking about individual quanta makes sense for <u>weakly</u>-coupled systems, but not so much for strongly-coupled systems.

<u>Complication</u>: QCD interaction strength depends on scale.



One of the Caveats

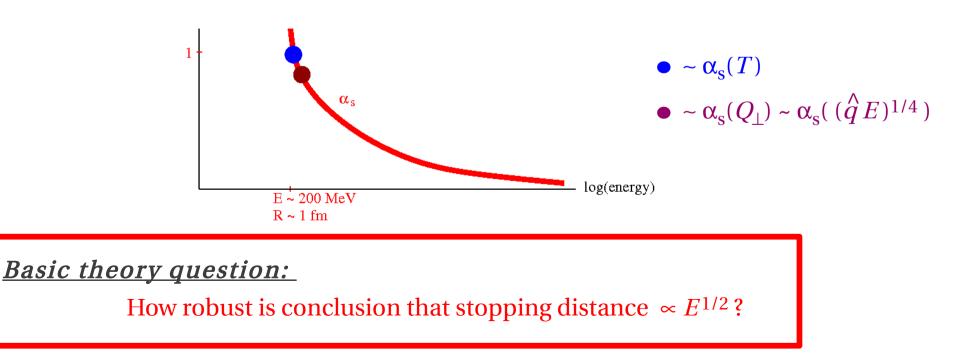




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Applied String Theory

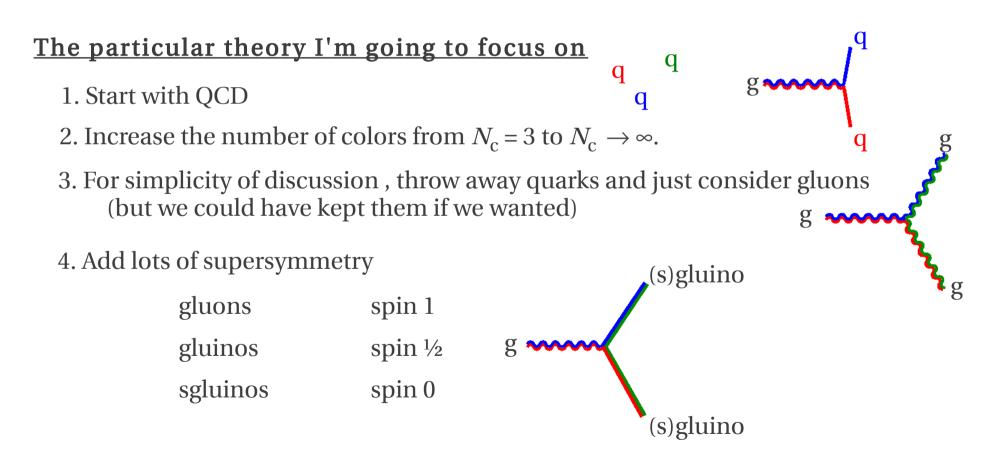
also referred to as, variously,

Applied Holography Gauge-String Duality Gauge-Gravity Duality the AdS/CFT correspondence Basic theory question:

How robust is conclusion that stopping distance $\propto E^{1/2}$?

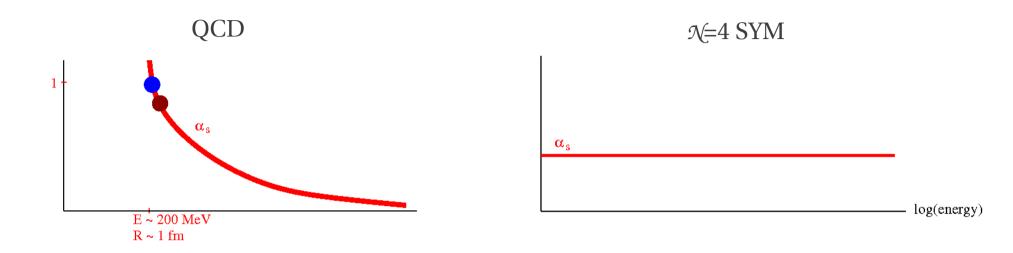
So far, too hard to study theoretically in QCD iteslf.

But we can study robustness by looking at theories similar to QCD that *can* be solved.



"large- N_c $\mathcal{N}=4$ Super-Yang-Mills (SYM) theory"

At zero temperature, this theory does not behave much like QCD



confining!

scale invariant!

no scales \rightarrow no confinement

 $[Note: value of \alpha_s can be set to whatever you want to study.]$

But a QCD plasma is very similar to a SYM plasma (at weak coupling)

Units: $\hbar = c = 1$

Debye screening length
$$\sim rac{\#}{lpha_{
m s}^{1/2}T}$$

Plasma frequency $\sim \# \alpha_{
m s}^{1/2} T$

Correlation length of color magnetic fields $\sim \frac{\#}{\alpha_{\rm s}T}$

Shear viscosity
$$\sim \frac{\#T^3}{\alpha_{\rm s}^2 \ln \frac{1}{\alpha_{\rm s}}}$$

Stopping distance $\sim \frac{\#E^{1/2}}{\alpha_{\rm s}^2 T^{3/2} \ln^{1/2}(E/T)}$
etc.

SYM is a good theorist playground for studying *lots* of things, (including the robustness of "stopping distance $\propto E^{1/2}$ ") because

(large- N_c N=4) SYM can <u>also</u> be solved at strong coupling!

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Gauge-Gravity Duality (the AdS/CFT correspondence)

Quantum Field Theory

3+1 dimensional space+time

no gravity (flat space)

𝔧=4 SYM a conformal field theory (CFT)

large $N_{\rm c}$

strong coupling

a theory of classical waves

4+1 dimensional space+time

a curved spacetime (General Relativity)

the spacetime is Anti-deSitter space (AdS)

can ignore quantum effects

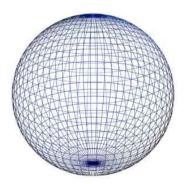
General Relativity instead of non-perturbative string theory

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A flavor of a thought #1

If the theories are the same, the symmetries of the 3+1 field theory should be the same as those of the 4+1 curved spacetime...

<u>isometries</u>



A flavor of a thought #1

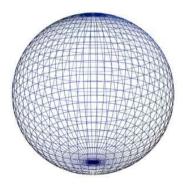
If the theories are the same, the symmetries of the 3+1 field theory should be the same as those of the 4+1 curved spacetime...

<u>symmetries</u>

translations rotations

Lorentz boosts

<u>isometries</u>



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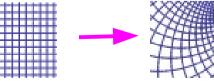
translations

rotations

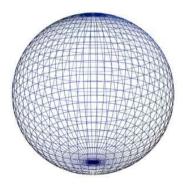
Lorentz boosts

scaling transformation

conformal transformations



<u>isometries</u>



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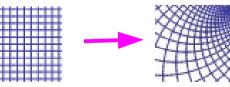
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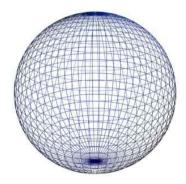
Lorentz boosts

scaling transformation

conformal transformations

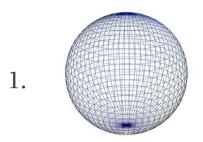


<u>isometries</u>



symmetry group of $3+1 \dim CFT = \text{isometry group of } 4+1 \dim AdS$

What is AdS?



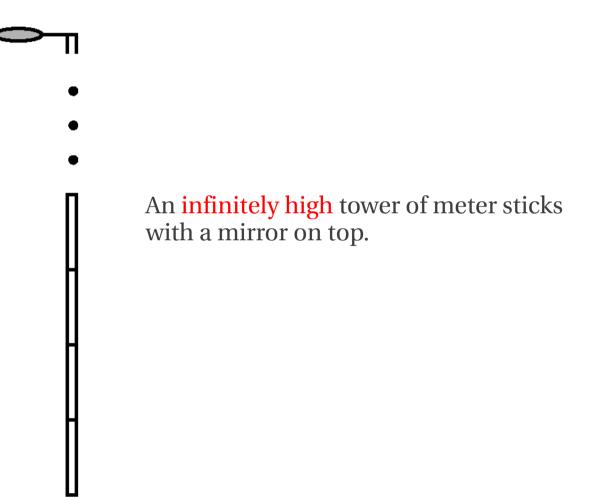
= a picture of a 2-dimensional surface called the "2-sphere"

- 2. Now instead think of a 5-sphere (if the goal is to work our way to 4+1 dim AdS)
- 3. Now imagine the surface instead has negative curvature everywhere (like a saddle except that every point has to look the same!) Unlike the sphere, this space has infinite size.
- 4. Finally, trade one of the five spatial dimensions for a time dimension.

Irrelevant connection to cosmology:

If our universe had nothing in it but the cosmological constant, it would be a 3+1 dim. deSitter spacetime. If the cosmological case were negative instead of positive, it would be a 3+1 dim. anti-deSitter spacetime,

<u>A flavor of a thought #2</u>

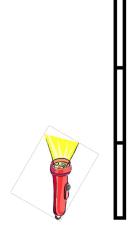


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<u>A flavor of a thought #2</u>

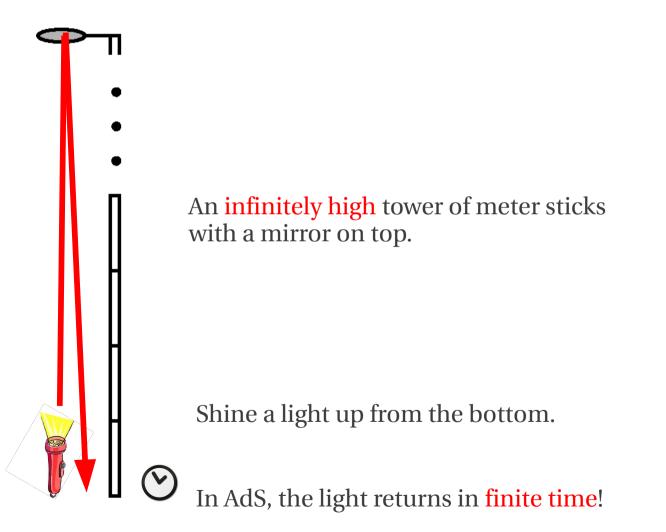
An **infinitely high** tower of meter sticks with a mirror on top.

Shine a light up from the bottom.

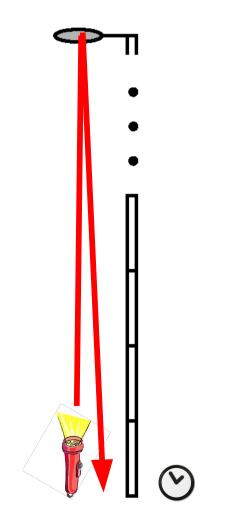


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<u>A flavor of a thought #2</u>



<u>A flavor of a thought #2</u>



How??

Think Pound and Rebka in 1959

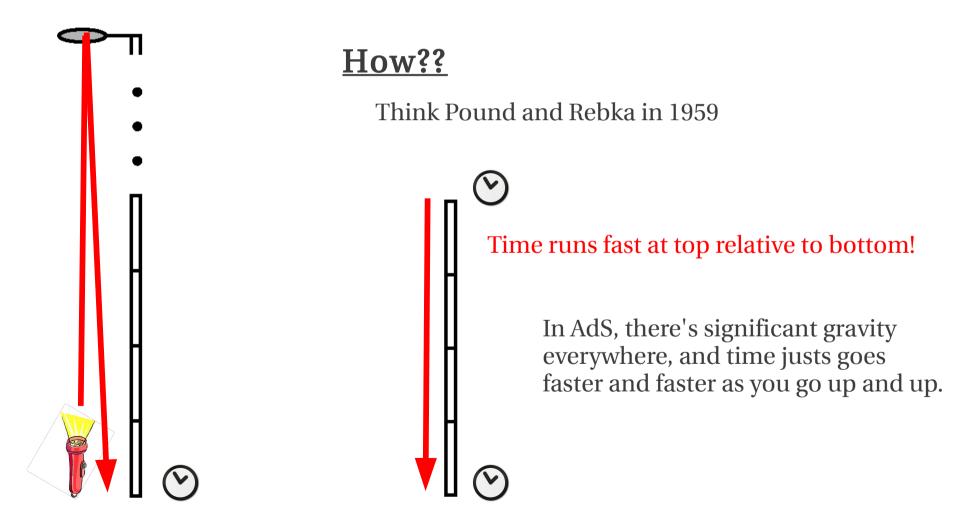
Time runs fast at top relative to bottom!

In AdS, there's significant gravity everywhere, and time justs goes faster and faster as you go up and up.

Why would a 3+1 dim. theory be equivalent to a 4+1 dim. theory?

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<u>A flavor of a thought #2</u>

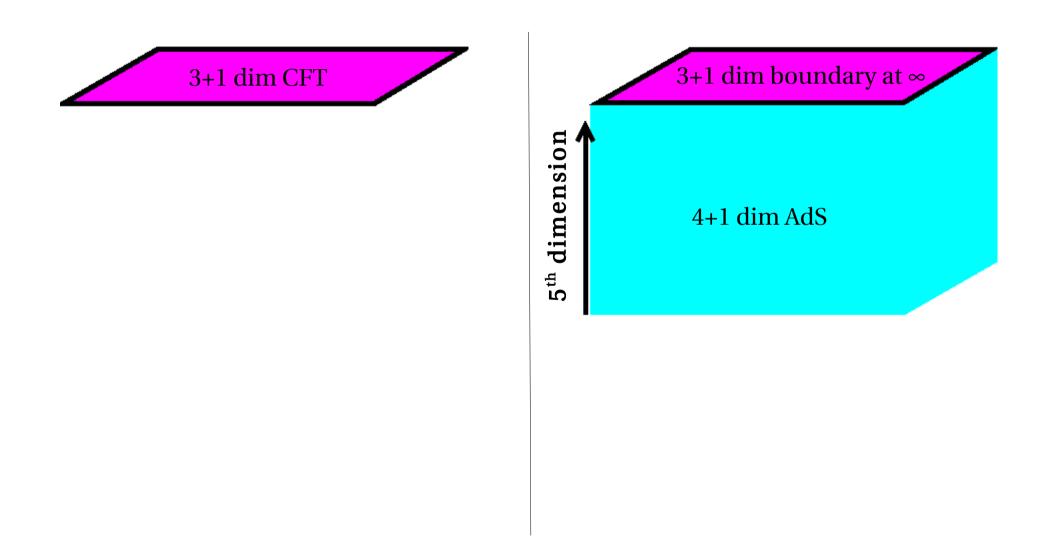


Upshot: What's happening at ∞ in AdS is not decoupled.

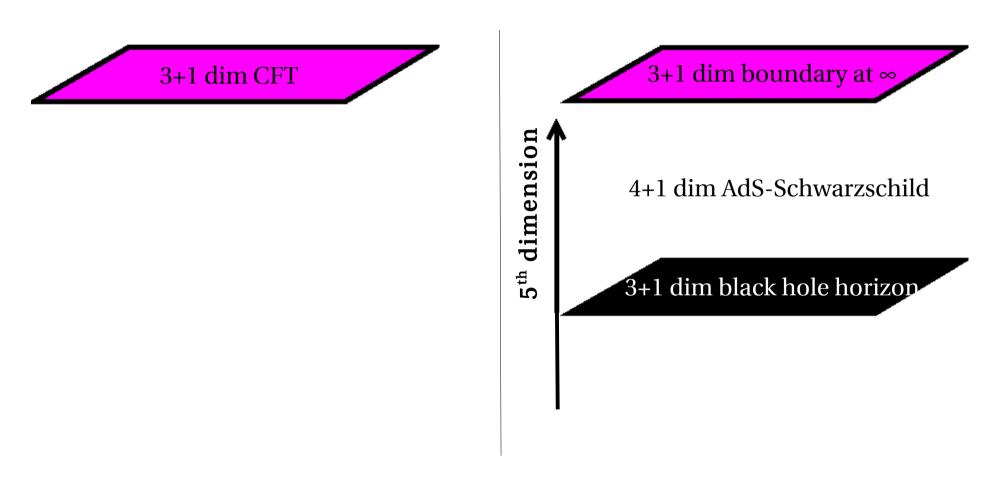
Why would a 3+1 dim. theory be equivalent to a 4+1 dim. theory?

22/27

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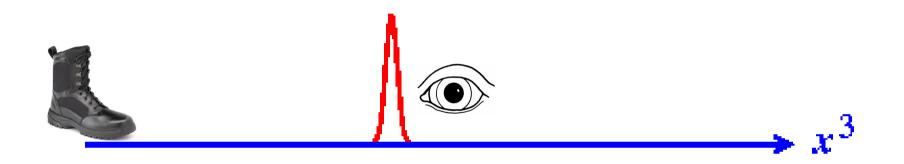
Q: How to study the theory at finite temperature? ^{23/27} A: Put a <u>black hole</u> in on the gravity side.

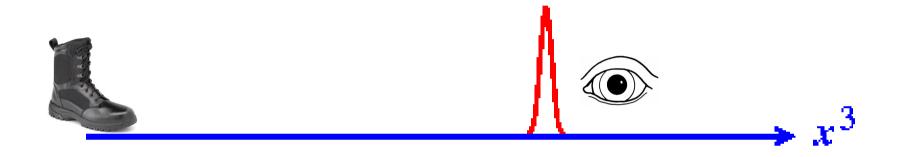


temperature for QFT = Hawking temp. of black hole



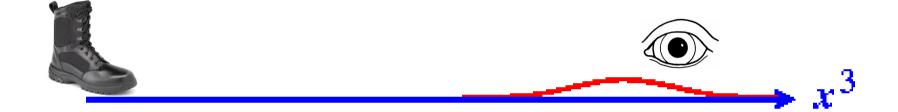


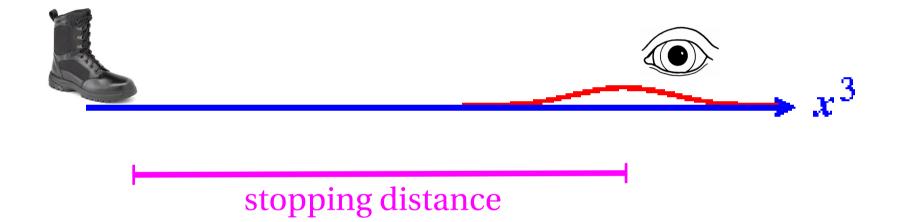


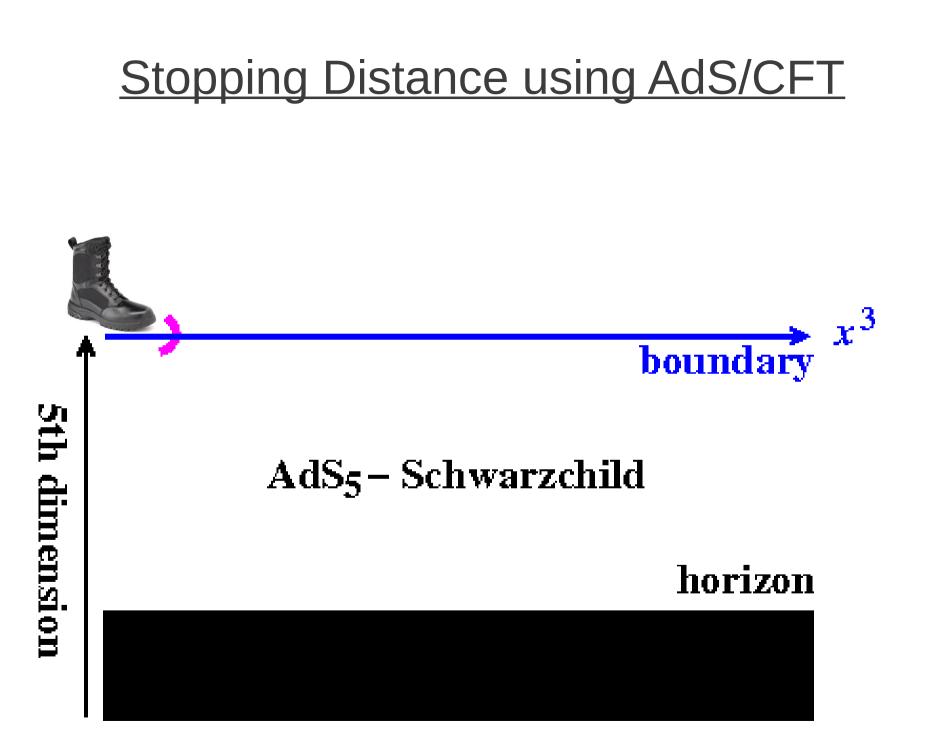








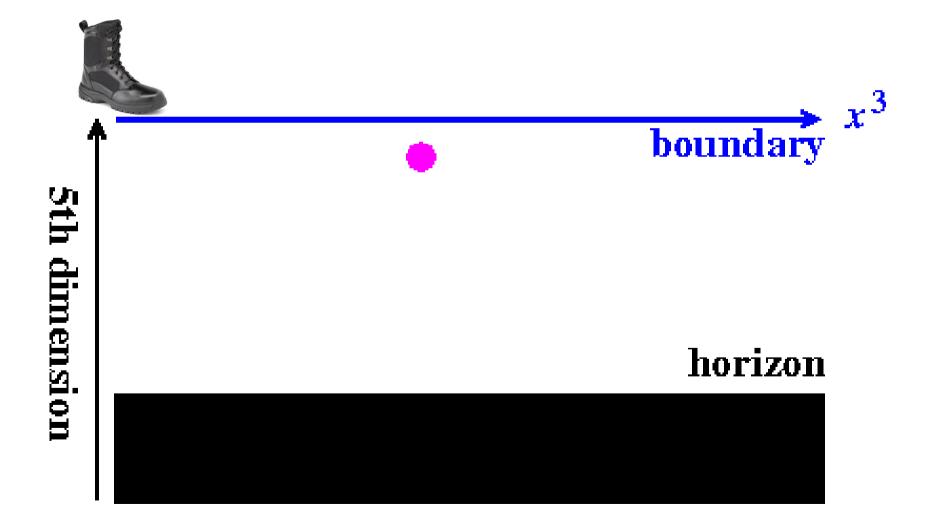


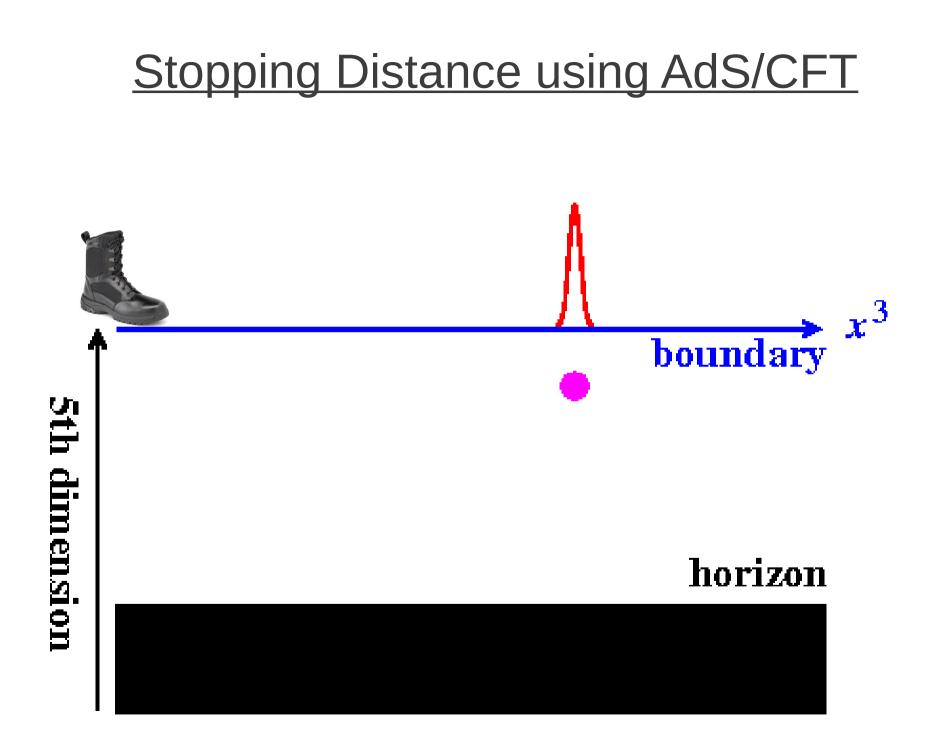


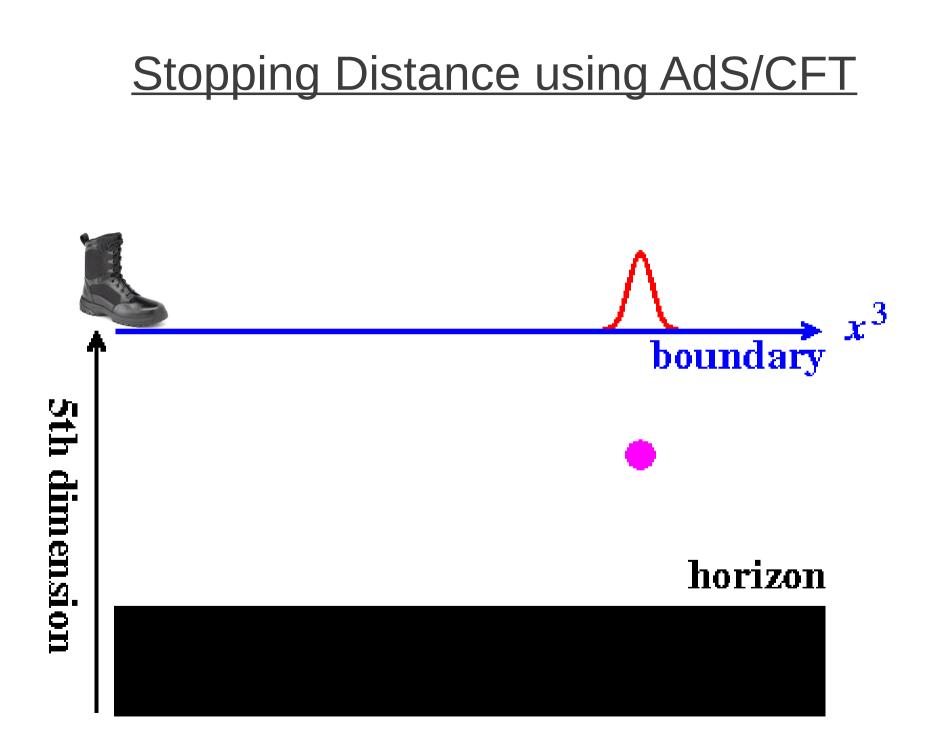
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Stopping Distance using AdS/CFT x^3 boundary 5th dimension horizon

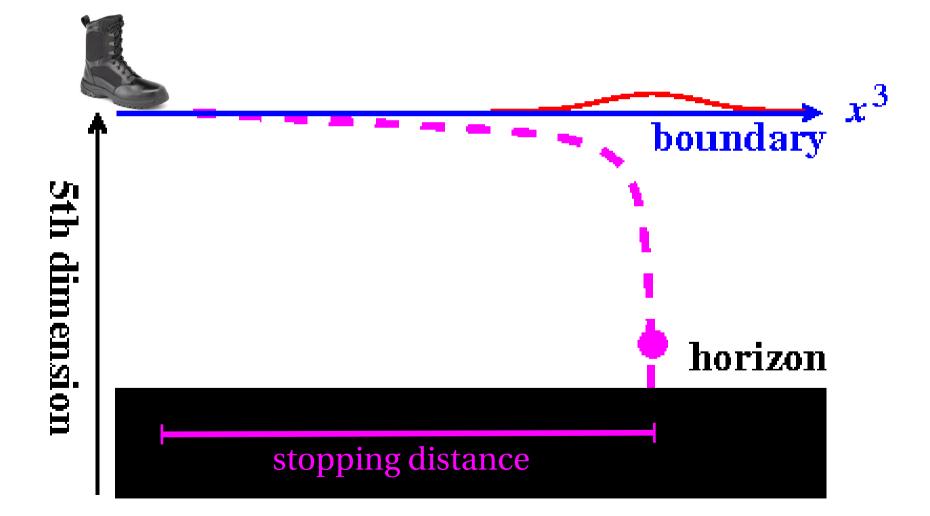
Stopping Distance using AdS/CFT







Stopping Distance using AdS/CFT



<u>Results</u>

weak coupling: stopping distance $\propto E^{1/2}$

strong coupling: stopping distance $\propto E^{1/3}$

[Gubser, Gollota, Pufu, Rocha; Hatta, Iancu, Mueller; Chesler, Jensen, Karch, Yaffe (2008)]

What did we learn?

1) <u>A point specific to the field</u>

One should look for corrections to the LPM energy dependence

stopping distance $\propto E^{\frac{1}{2} + \dots}$

in real QCD, even if the relevant coupling for high-energy jets is not BIG with a capital "B"

2) <u>A general point about the nature of theoretical research</u>

One of many cases where basic, curiosity-driven research into theory (e.g. string theory, supergravity in weird curved spacetimes, ...) yielded insights in unanticipated applications.

What I personally have been doing lately (w/ Diana Vaman, Phil Szepietwoski, Gabriel Wong)

When I said

strong coupling: stopping distance $\propto E^{1/3}$

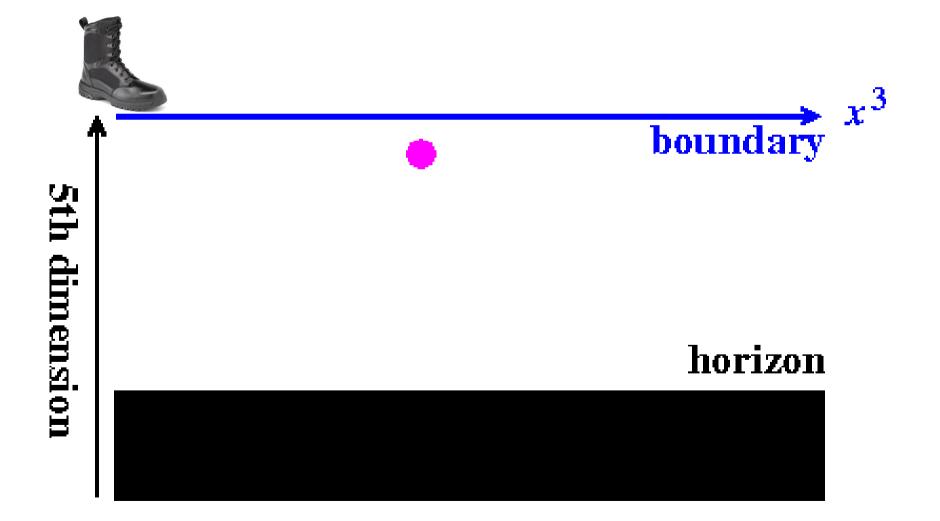
strong really meant coupling $\rightarrow \infty$.

But is this really the high energy behavior if, instead, coupling = large but finite ? What happens if, for example,



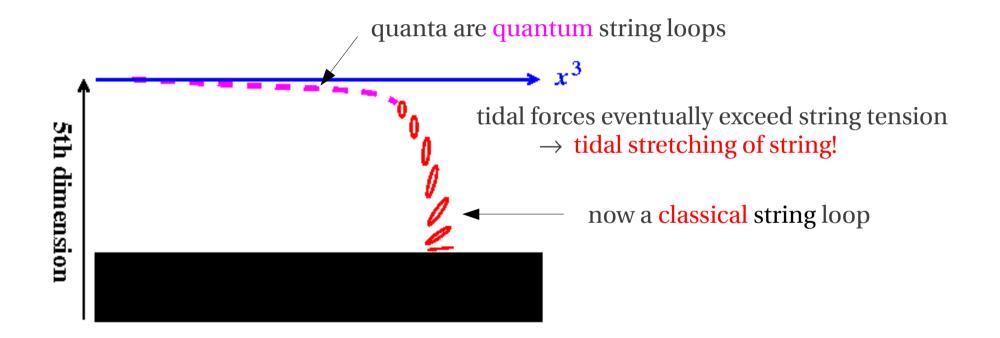
which you really can't explore in any calculation that takes coupling = ∞ .

<u>Reminder</u>



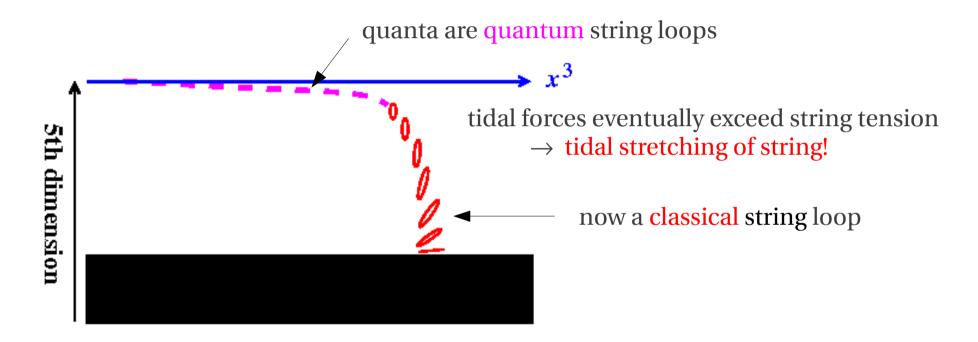
But really

= w/ internal degrees of freedom in ground state
 proper size ~ (string tension)^{-1/2}



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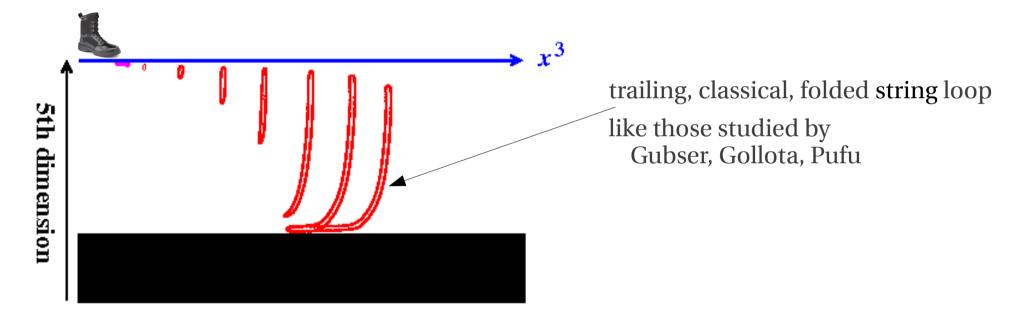
<u>Result</u>

Interesting stuff, but does not change the following:

strong but finite coupling: maximum stopping distance $\propto E^{1/3}$

A quick aside for those who follow AdS/CFT jet stuff

- Q: What happens if you adjust the "kick" (virtuality) that creates the jet so that it stops earlier than the maximum stopping distance?
- A: It's possible to get the following:



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