

C-DEFORMATION in SUPERSYMMETRY and SUPERGRAVITY

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based on

hep-th/0502026, 0504191, 0506071, 0602066, 0602115

all published in NPB, PLB, CQG

with T. Hatanaka, Y. Kobayashi (TMU), S. Sasaki (TIT)

BASIC IDEAS

spacetime non-commutativity
eff. field theory ← from strings

Seiberg, Witten (1999)
Duguri, Vafa (2003)
Seiberg (2003)
Ferrara, Lledo (2000)
Klemm, Tamassia, Penati
(2001)... etc.

superspace non-ANTI-commutativity
susy field theory ← from (nonperturbative) superstrings
and branes

- gravi-photon $\langle C_{\mu\nu}^+ \rangle \neq 0$!
- partial susy breaking to $N = 1/2$

NON-COMMUTATIVITY

Motivation

- quantum spacetime
- controlled non-locality in quantum field theory
- signatures of underlying theory (strings and branes)
- abelian solitons
- singularity resolution

Technicalities

$$[x^\mu, x^\nu]_- = i B^{\mu\nu} = \text{const.} \neq 0$$

Moyal-Weyl star product:

$$f(x) \star g(x) = f(x) \exp\left(\frac{i}{2} \overleftarrow{\partial}_\mu B^{\mu\nu} \overrightarrow{\partial}_\nu\right) g(x)$$

Problems

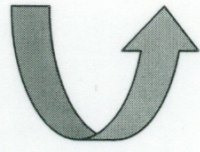
- violation of fundamental symmetries, e.g. Lorentz inv.
- UV/IR mixing \rightarrow questionable renormalization
- in finitely many derivatives \rightarrow unsolvable equations etc.

N=1 non-anticommutative superspace

Gravitphoton background \rightarrow N=1 non(anti)commutative superspace
(Ooguri and Vafa (2003), Seiberg (2003))

$$\left\{ \theta^\alpha, \theta^\beta \right\} = C^{\alpha\beta} \neq 0 \quad \text{Chiral part only}$$

$$\left\{ \bar{\theta}^{\dot{\alpha}}, \bar{\theta}^{\dot{\beta}} \right\} = \left\{ \bar{\theta}^{\dot{\alpha}}, \theta^\mu \right\} = 0$$



On four dimensional N=1 Euclidean superspace

$$f(\theta) \star g(\theta) = f(\theta) \exp \left[-\frac{C^{\alpha\beta} \overleftarrow{\partial} \overrightarrow{\partial}}{2 \partial\theta^\alpha \partial\theta^\beta} \right] g(\theta)$$

star product

$$= f(\theta) \left[1 - \frac{C^{\alpha\beta} \overleftarrow{\partial} \overrightarrow{\partial}}{2 \partial\theta^\alpha \partial\theta^\beta} + \det C \frac{\overleftarrow{\partial} \overrightarrow{\partial}}{\partial\theta^\alpha \partial\theta^\beta} \right] g(\theta)$$

Preserves only N=1/2 supersymmetry

NAC deformation—Lower star power example

NAC deformed Wess-Zumino model (Seiberg(2003))

$$\begin{aligned}\mathcal{L} &= \int d^4\theta \bar{\Phi} \star \Phi + \int d^2\theta \left(\frac{1}{2} m \Phi \star \Phi + \frac{1}{3} g \Phi \star \Phi \star \Phi \right) \\ &\quad + \int d^2\bar{\theta} \left(\frac{1}{2} \bar{m} \bar{\Phi} \star \bar{\Phi} + \frac{1}{3} \bar{g} \bar{\Phi} \star \bar{\Phi} \star \bar{\Phi} \right) \\ &= \mathcal{L}(c=0) - \frac{1}{3} g \det c F^3 + (\text{total derivative})\end{aligned}$$

Quantum property and renormalizability

R.Britto, B.Feng, S.J.Rey (2003), R.Britto, B.Feng (2003)

S.Terashima J.Yee (2003), A.Romagnoni (2003)

O.Lunin, S.J.Rey (2003), D.Berenstein, S.J.Rey (2003)

A.T.Banin, I.L.Buchbinder, N.G.Pletney (2004)

Etc...

NAC deformed gauge theory (Seiberg(2003))

Modified Wess-Zumino gauge

$$V_c(y, \theta, \bar{\theta}) = V(y, \theta, \bar{\theta}) - \frac{i}{4} \bar{\theta}^2 \theta^\alpha \varepsilon_{\alpha\beta} C^{\beta\gamma} \sigma_{\gamma\dot{\gamma}}^\mu \{ \bar{\lambda}^{\dot{\gamma}}, A_\mu \}$$

$$\int d^2\theta \operatorname{Tr} W^\alpha W_\alpha|_{\star} = \int d^2\theta \operatorname{Tr} W^\alpha W_\alpha (C=0) - i C^{\mu\nu} \operatorname{Tr} F_{\mu\nu} \bar{\lambda} \bar{\lambda} + \frac{|C|^2}{4} \operatorname{Tr} (\bar{\lambda} \bar{\lambda})^2$$
$$\int d^2\theta \operatorname{Tr} \bar{W}_{\dot{\alpha}} \bar{W}^{\dot{\alpha}}|_{\star} = \int d^2\theta \operatorname{Tr} W^\alpha W_\alpha (C=0) - i C^{\mu\nu} \operatorname{Tr} F_{\mu\nu} \bar{\lambda} \bar{\lambda} + \frac{|C|^2}{4} \operatorname{Tr} (\bar{\lambda} \bar{\lambda})^2$$

+ (total derivative)

Various works

- T.Araki, K.Ito, A.Ohtsuka (2003,2004,2005)
M.Alishahiha, A.Ghods, N.Sadooghi (2004)
I.Jack, D.R.T.Jones, L.A.Worthy (2004), A.Imaanpur (2003)
P.A.Grassi, R.Ricci, D.Robles-Llana (2003)
M.Billo, M.Frau, I.Pesando and A.Lerda (2004)
Etc...

4D single chiral model deformation

(T.Hatanaka, S.V.Ketov, Y.Kobayashi and S.S, (hep-th/0502026))

Superpotential part $\Phi = \phi + i\sqrt{2}\theta\psi + \theta^2 M$

$$\int d^2\theta W_*(\Phi) = \frac{1}{2c} [W(\phi + cM) - W(\phi - cM)]$$
$$c \equiv \sqrt{-\det C} - \frac{\chi^2}{4cM} [W_{,\phi}(\phi + cM) - W_{,\phi}(\phi - cM)]$$

Non(anti)commutativity induces splitting effect in the superpotential

Generally, to calculate Kahler part deformation is very complicated
→ Chiral reduced model is proposed



Result is

$$\begin{aligned}
 L_b = & \frac{1}{2} \partial^\mu \bar{\phi} \partial_\mu \bar{\phi} [K_{,\bar{\phi}\bar{\phi}}(\phi + cM, \bar{\phi}) + K_{,\bar{\phi}\bar{\phi}}(\phi - cM, \bar{\phi})] \\
 & + \frac{1}{2} \partial^2 \bar{\phi} [K_{,\bar{\phi}}(\phi + cM, \bar{\phi}) + K_{,\bar{\phi}}(\phi - cM, \bar{\phi})] \\
 & + \frac{\bar{M}}{2c} [K_{,\bar{\phi}}(\phi + cM, \bar{\phi}) - K_{,\bar{\phi}}(\phi - cM, \bar{\phi})] \\
 & + \frac{1}{2c} [W(\phi + cM) - W(\phi - cM)] + \bar{M} \frac{\partial \bar{W}}{\partial \bar{\phi}}
 \end{aligned}$$

$$\begin{aligned}
 L_f = & -\frac{1}{4c} \bar{\chi}^2 [K_{,\bar{\phi}\bar{\phi}}(\phi + cM, \bar{\phi}) - K_{,\bar{\phi}\bar{\phi}}(\phi - cM, \bar{\phi})] \\
 & - \frac{i}{2cM} (\chi \sigma^\mu \bar{\chi}) \partial_\mu \bar{\phi} [K_{,\bar{\phi}\bar{\phi}}(\phi + cM, \bar{\phi}) - K_{,\bar{\phi}\bar{\phi}}(\phi - cM, \bar{\phi})] \\
 & - \frac{i}{2cM} (\chi \sigma^\mu \partial_\mu \bar{\chi}) [K_{,\bar{\phi}}(\phi + cM, \bar{\phi}) - K_{,\bar{\phi}}(\phi - cM, \bar{\phi})] \\
 & - \frac{\bar{M}}{4cM} \chi^2 [K_{,\phi\bar{\phi}}(\phi + cM, \bar{\phi}) - K_{,\phi\bar{\phi}}(\phi - cM, \bar{\phi})] \\
 & \dots
 \end{aligned}$$

Multi chiral deformation

Two dimensional result (L.A-Gaume and M.A.V-Mozo (2005))

Four dimensional chiral reduced model (HKKS(hep-th/0506071))

Chiral reduced multi-chiral NAC deformed NLSM Lagrangian

$$\mathcal{L} = \frac{1}{2}M^i Y_{,i} + \frac{1}{2}\partial^\mu \bar{\phi}^{\bar{p}} \partial_\mu \bar{\phi}^{\bar{q}} Z_{,\bar{p}\bar{q}} + \frac{1}{2}\partial^2 \bar{\phi}^{\bar{p}} Z_{,\bar{p}} - \frac{1}{4}(\chi^i \chi^j) Y_{,ij} - \frac{1}{2}i(\chi^i \sigma^\mu \bar{\chi}^{\bar{p}}) \partial_\mu \bar{\phi}^{\bar{q}} Z_{,\bar{p}\bar{q}} - \frac{1}{2}i(\chi^i \sigma^\mu \partial_\mu \bar{\chi}^{\bar{p}}) Z_{,i\bar{p}}$$

$$Z(\phi, \bar{\phi}, M) = \int_{-1}^1 d\xi K^\xi = \int_{-1}^1 d\xi K(\phi^i + \xi c M^i, \bar{\phi}^{\bar{j}})$$

$$Y(\phi, \bar{\phi}, M, \bar{M}) = \bar{M}^{\bar{p}} Z_{,\bar{p}} - \frac{1}{2}(\bar{\chi}^{\bar{p}} \bar{\chi}^{\bar{q}}) Z_{,\bar{p}\bar{q}} + c \int_{-1}^1 d\xi \xi [\partial^\mu \bar{\phi}^{\bar{p}} \partial_\mu \bar{\phi}^{\bar{q}} K^\xi_{,\bar{p}\bar{q}} + \partial^2 \bar{\phi}^{\bar{p}} K^\xi_{,\bar{p}}]$$

$$\mathcal{L}_{\text{pot.}} = \frac{1}{2}M^i \tilde{W}_{,i} - \frac{1}{4}(\chi^i \chi^j) \tilde{W}_{,ij} + \bar{M}^{\bar{p}} \bar{W}_{,\bar{p}} - \frac{1}{2}(\bar{\chi}^{\bar{p}} \bar{\chi}^{\bar{q}}) \bar{W}_{,\bar{p}\bar{q}}$$

$$\tilde{W}(\phi, M) = \int_{-1}^1 d\xi W(\phi^i + \xi c M^i)$$

Deformed CP¹ metric with superpotential

$$M = \frac{\alpha - \sqrt{\alpha^2 + (2c\bar{\phi}(1 + \kappa^{-2}\phi\bar{\phi})\bar{W}_{,\bar{\phi}})^2}}{2c^2\kappa^{-2}\bar{\phi}^2\bar{W}_{,\bar{\phi}}}$$

On-shell auxiliary field

$$g_{\phi\bar{\phi}} = \frac{\alpha\kappa^{-2}c^2\bar{\phi}^2(\bar{W}_{,\bar{\phi}})^2}{\left(\alpha - \sqrt{\alpha^2 + (2c\bar{\phi}(1 + \kappa^{-2}\phi\bar{\phi})\bar{W}_{,\bar{\phi}})^2}\right)^2} \sqrt{\alpha^2 + (2c\bar{\phi}(1 + \kappa^{-2}\phi\bar{\phi})\bar{W}_{,\bar{\phi}})^2}$$

$$g_{\phi\phi} = 0$$

$$g_{\bar{\phi}\bar{\phi}} = \frac{-2\alpha^{-1}c^2(1 + \kappa^{-2}\phi\bar{\phi})\bar{W}_{,\bar{\phi}}}{\left(\alpha - \sqrt{\alpha^2 + (2c\bar{\phi}(1 + \kappa^{-2}\phi\bar{\phi})\bar{W}_{,\bar{\phi}})^2}\right)^2} \sqrt{\alpha^2 + (2c\bar{\phi}(1 + \kappa^{-2}\phi\bar{\phi})\bar{W}_{,\bar{\phi}})^2} \\ \times [4c^2\bar{\phi}^2(\bar{W}_{,\bar{\phi}})^3(1 + \kappa^{-2}\phi\bar{\phi}) \\ + \alpha(\alpha - \sqrt{\alpha^2 + (2c\bar{\phi}(1 + \kappa^{-2}\phi\bar{\phi})\bar{W}_{,\bar{\phi}})^2})(2\bar{W}_{,\bar{\phi}} + \bar{\phi}\bar{W}_{,\bar{\phi}\bar{\phi}})]$$



Metric is controlled by the superpotential

FURTHER APPLICATIONS

- \mathbb{C} -deformation of *effective potentials* in *susy gauge theories* (hep-th/0502026)

\mathbb{C} -deformed Veneziano-Yankielowicz superpotential:

- (ii) unbroken gauge invariance
- (iii) no further dynamical susy breaking
- ✓ (i) non-vanishing gluino condensate

- \mathbb{C} -deformation of $N=2$ extended *supergravity* (hep-th/0602115):

- (i) partial local susy breaking (to $N=1/2$)
- (ii) relation to *gravitino* condensate:

$$\langle C_{\mu\nu}^+ \rangle = \frac{\kappa}{2} \langle \bar{\Psi}_{[\mu} \chi_{\nu]} \rangle^+$$

and much more...

CONCLUSION

\mathbb{C} -deformation \sim simple window into some nonperturbative physics of *susy gauge theories* and *superstrings*.