



The gluon-fusion cross section at N³LO in the soft limit

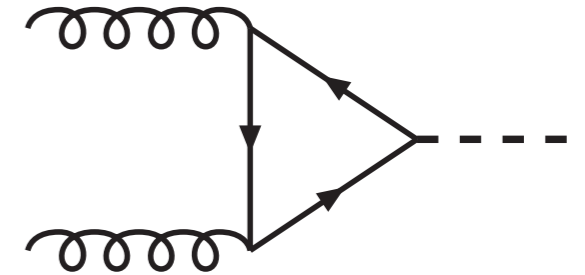
Claude Duhr

in collaboration with C. Anastasiou, F. Dulat, E. Furlan,
T. Gehrmann, F. Herzog, B. Mistlberger

Amplitudes 2014
Saclay, 11/06/2014

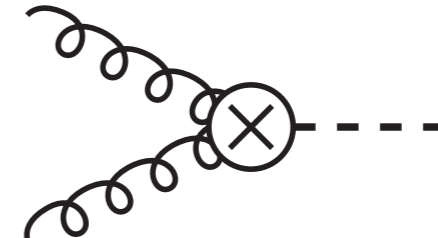
The gluon fusion cross section

- The dominant Higgs production mechanism at the LHC is gluon fusion.
➔ Loop-induced process.



- For a light Higgs boson, the dimension five operator describing a tree-level coupling of the gluons to the Higgs boson

$$\mathcal{L} = \mathcal{L}_{QCD,5} - \frac{1}{4v} C_1 H G_{\mu\nu}^a G_a^{\mu\nu}$$



- Top-mass corrections known at NNLO.

[Harlander, Ozeren; Pak, Rogal, Steinhauser; Ball, Del Duca, Marzani, Forte, Vicini; Harlander, Mantler, Marzani, Ozeren]

- In the rest of the talk, I will only concentrate on the effective theory.

The gluon fusion cross section

- The gluon fusion cross section is given in perturbation theory by

$$\sigma(pp \rightarrow H + X) = \tau \sum_{ij} \int_{\tau}^1 dz \mathcal{L}_{ij}(z) \hat{\sigma}_{ij}(\tau/z)$$

- The (partonic) cross section depends up to an overall scale only on the ratio

$$\tau = \frac{m^2}{s} \qquad z = \frac{m^2}{\hat{s}}$$

- The partonic cross section known at NLO and NNLO.

[Dawson; Djouadi, Spira, Zerwas;
Harlander, Kilgore; Anastasiou, Melnikov;
Ravindran, Smith, van Neerven]

	σ [8 TeV]	$\delta\sigma$ [%]
LO	9.6 pb	$\sim 25\%$
NLO	16.7 pb	$\sim 20\%$
NNLO	19.6 pb	$\sim 7 - 9\%$
N3LO	???	$\sim 4 - 8\%$

[Fixed order only]

The gluon fusion cross section

- So far no complete computation is available.
 - ➔ Scale variation at N³LO.
 - ➔ Approximate N³LO results exist.

[Moch, Vogt; Ball, Bonvini, Forte, Marzani, Ridolfi; Bühler, Lazopulos

The gluon fusion cross section

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- Can we push the state of the art one order higher?

- Challenge:

performed so far...

- ➔ Uncharted territory!

- ➔ New conceptual challenges.

HopfAlgebras
ExpansionByRegions
MellinBarnes
MultiloopAmplitudes
ZetaValues
NestedSums
SoftCurrent
UV-Singularities
HypergeometricFunctions
Renormalisation
Symbols
Polylogarithms
IR-Singularities
DimensionalShiftIdentities
DifferentialEquations
IBPs
PhaseSpace
SectorDecomposition
Unitarity
MassFactorization

Prime example of how new developments from the amplitude community can have impact on phenomenology.

Outline

- The inclusive gluon-fusion cross section.
- Ingredients entering the cross section at threshold:
 - ➔ Soft triple-real emission.
 - ➔ Soft double-virtual-real emission.
 - ➔ Soft virtual-double-real emission.
- The gluon-fusion cross section in the soft limit.
- Conclusion & outlook.

The inclusive gluon-
fusion cross section

The gluon fusion cross section

- A cross section computation requires two ingredients:

$$\hat{\sigma} = \int d\Phi |\mathcal{M}|^2$$

The gluon fusion cross section

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Matrix
element

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Phase space
integration

Matrix
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The gluon fusion cross section

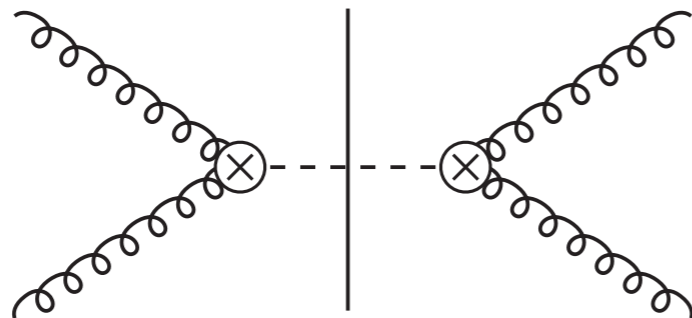
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- Example:



The gluon fusion cross section

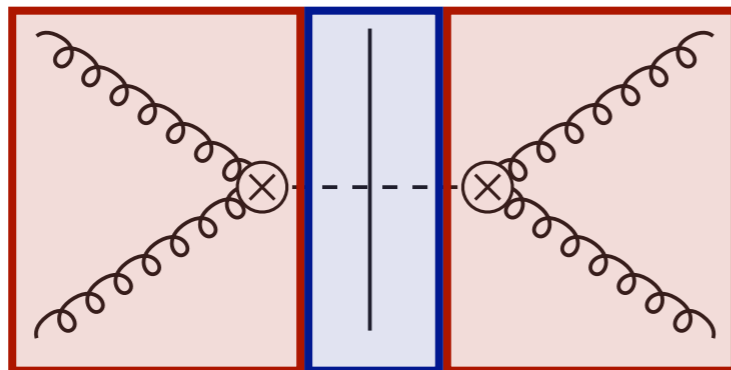
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Phase space
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- Example:

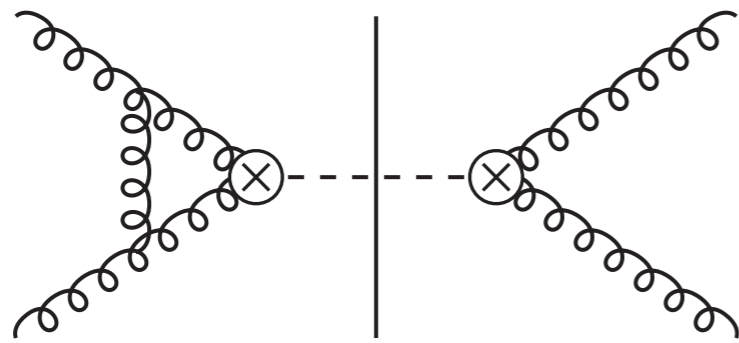


$$\int d\Phi_1 \mathcal{M}^{(0)} \mathcal{M}^{(0)*}$$

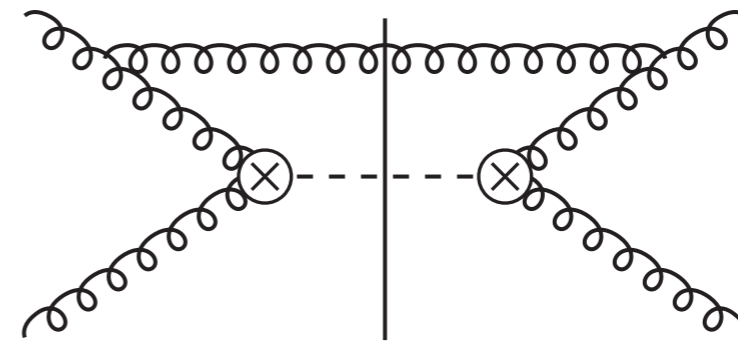
The gluon fusion cross section

- A_t

[Dawson; Djouadi, Spira, Zerwas]



Virtual corrections ('loops')

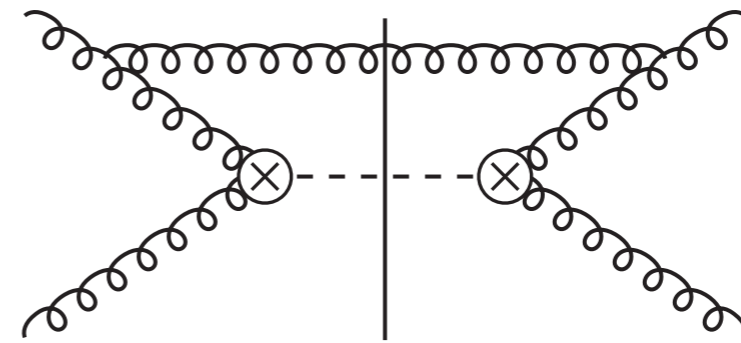
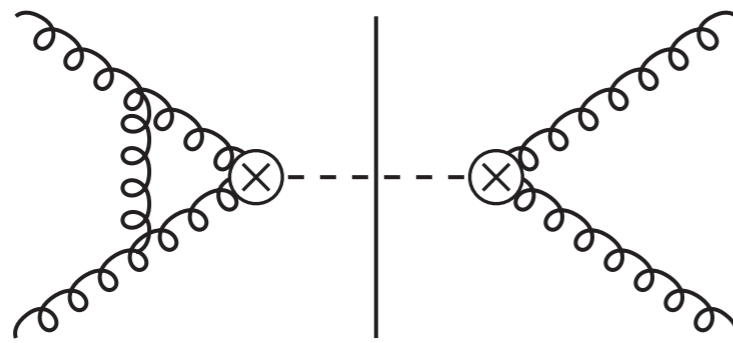


Real emission

The gluon fusion cross section

- At

[Dawson; Djouadi, Spira, Zerwas]



Virtual corrections ('loops')

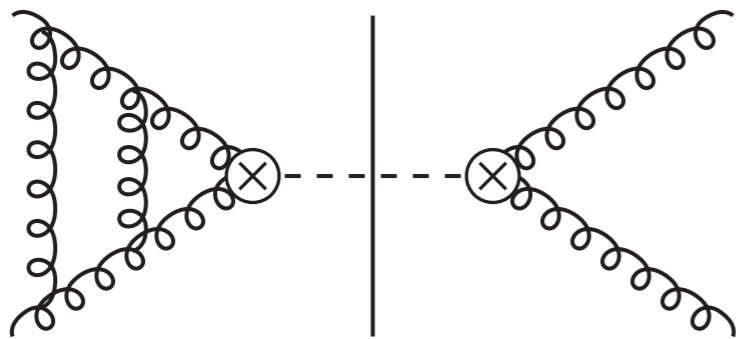
Real emission

- Both contributions are individually divergent:
 - ➔ UV divergences are handled by renormalization.
 - ➔ IR divergences cancelled by PDF counterterms.

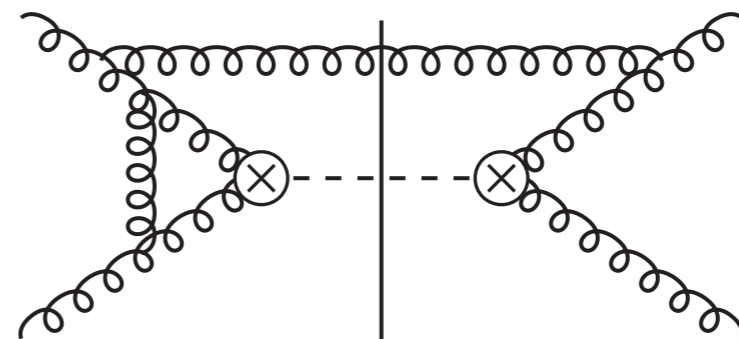
The gluon fusion cross section

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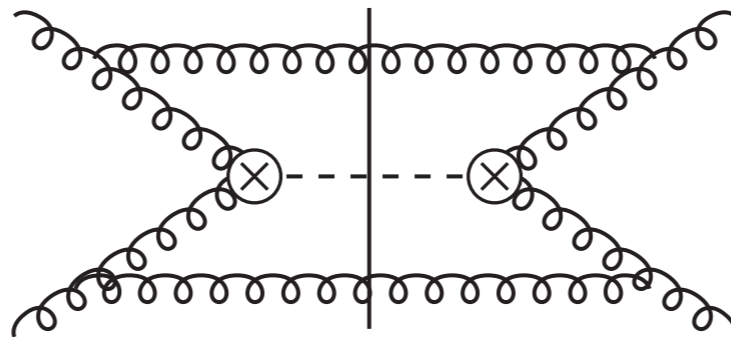
[Harlander, Kilgore; Anastasiou, Melnikov; Ravindran, Smith, van Neerven]



Double virtual



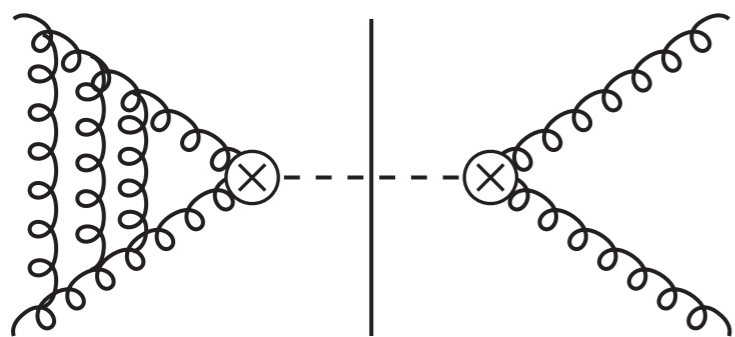
Real-virtual



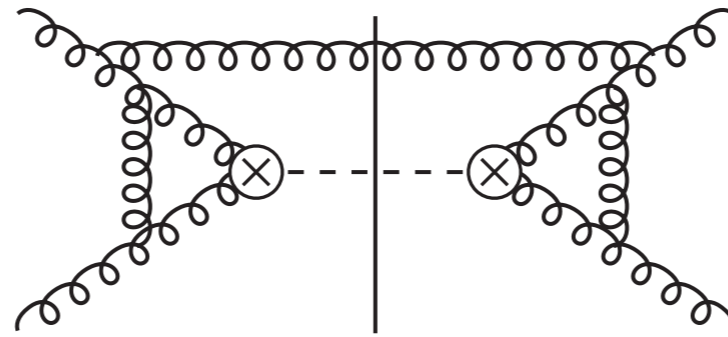
Double real

The gluon fusion cross section

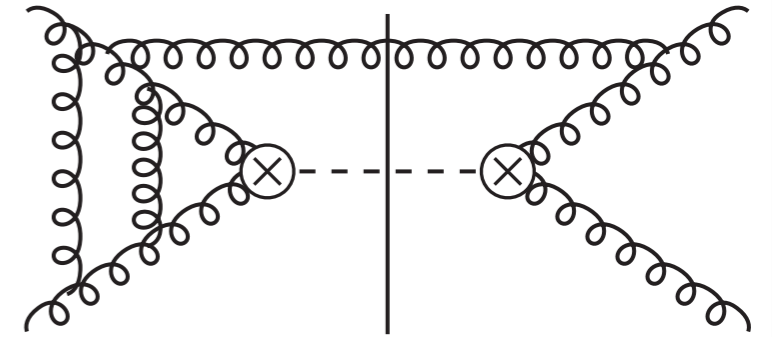
● At



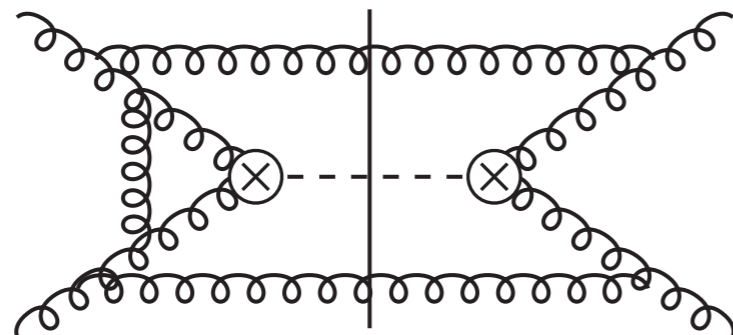
Triple virtual



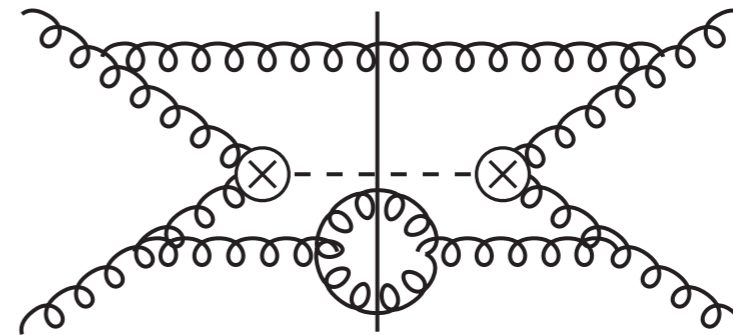
Real-virtual squared



Double virtual real



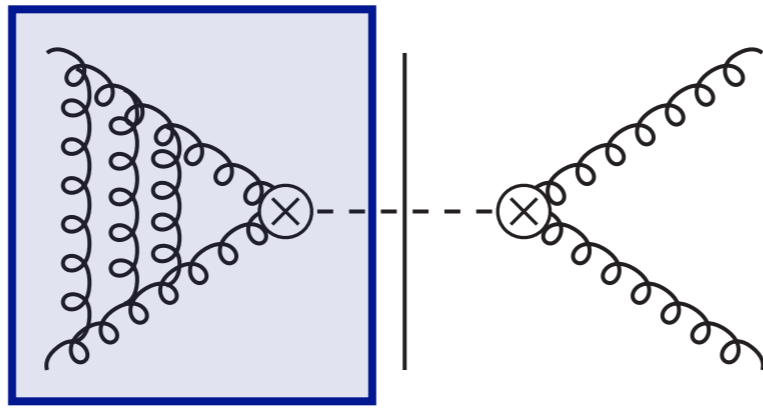
Double real virtual



Triple real

Triple virtual corrections

- The triple virtual corrections are directly related to the QCD form factor



- The QCD form factor is known

- ➔ at one loop.

- ➔ at two loops.

- ➔ at three loops.

[Gonsalves; Kramer, Lampe;
Gehrmann, Huber, Maître]

[Baikov, Chetyrkin, Smirnov, Smirnov, Steinhauser;
Gehrmann, Glover, Huber, Ikizlerli, Studerus]

- It is not the loops that are the problem!

Unitarity

- Optical theorem:

$$\text{Im} \text{ (loop diagram) } = \int d\Phi \text{ (cut diagrams)}$$

- ➔ Discontinuities of loop amplitudes are phase space integrals.
- Discontinuities of loop integrals are given by rule

$$\frac{1}{p^2 - m^2 + i\epsilon} \rightarrow \delta_+(p^2 - m^2) = \delta(p^2 - m^2) \theta(p^0)$$

- ➔ See Ruth Britto's talk.

Reverse-unitarity

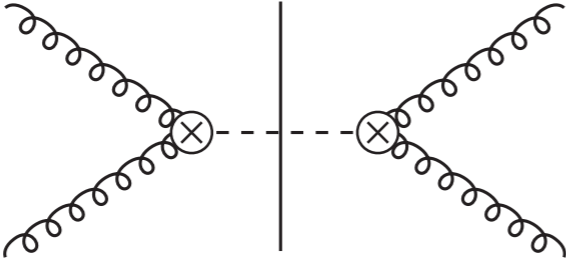
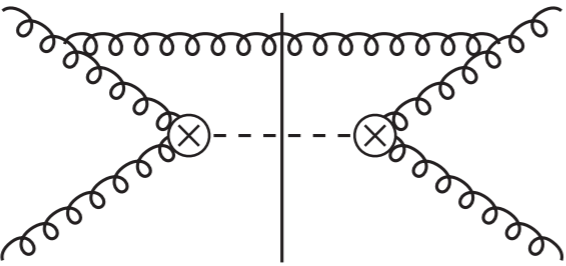
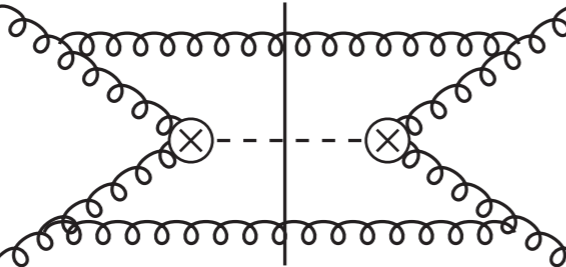
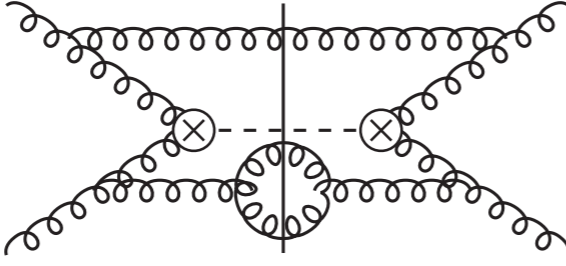
- Optical theorem:

$$\text{Im} \text{ (circle with 4 arrows)} = \int d\Phi \text{ (two ovals with 4 arrows and a dashed line)}$$

- We can read the optical theorem ‘backwards’ and write inclusive phase space integrals as unitarity cuts of loop integrals. [Anastasiou, Melnikov; Anastasiou, Dixon, Melnikov, Petriello]
 - ➔ Rather than computing phase-space integrals, we can compute loop integrals with cuts!
 - ➔ Makes inclusive phase space integrals accessible to all the technology developed for multi-loop computations!
 - ▶ Integration-by-parts & differential equations.

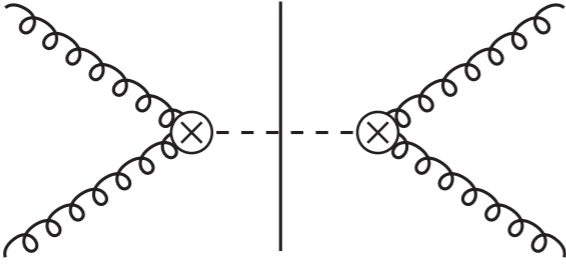
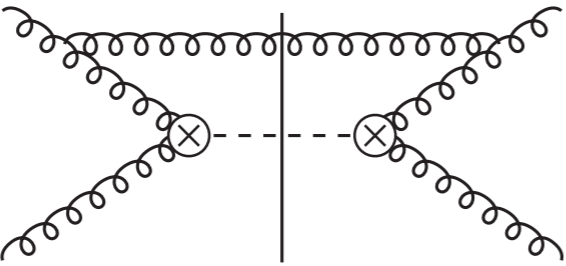
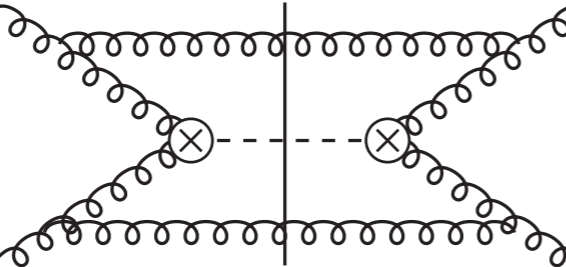
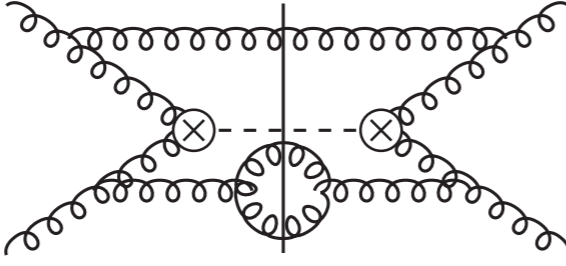
Reverse-unitarity @ N3LO

Growth in complexity for real emission

LO		1 diagram	1 integral
NLO			
NNLO			
N3LO			

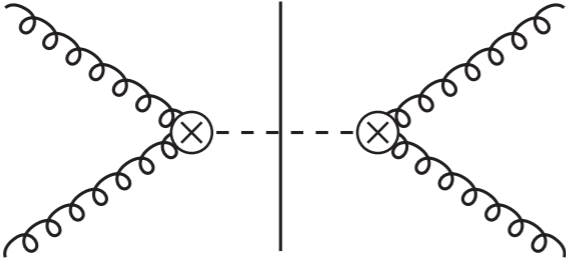
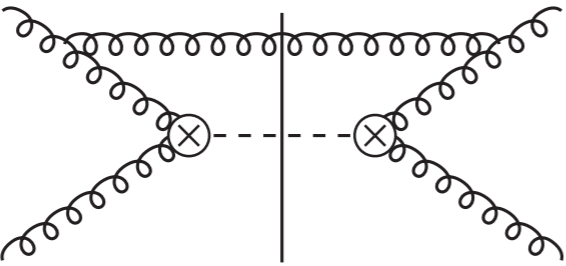
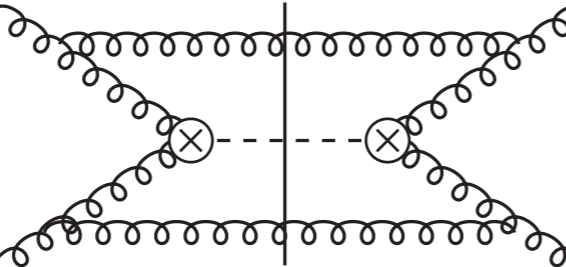
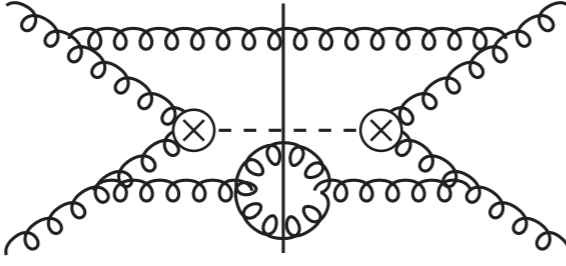
Reverse-unitarity @ N3LO

Growth in complexity for real emission

LO		1 diagram	1 integral
NLO		10 diagrams	1 integral
NNLO			
N3LO			

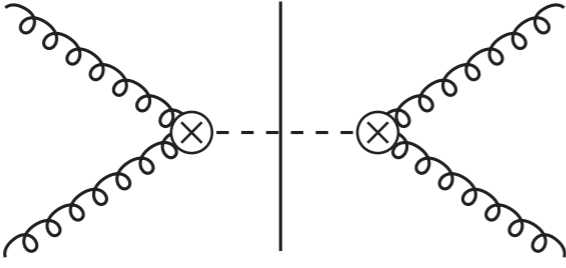
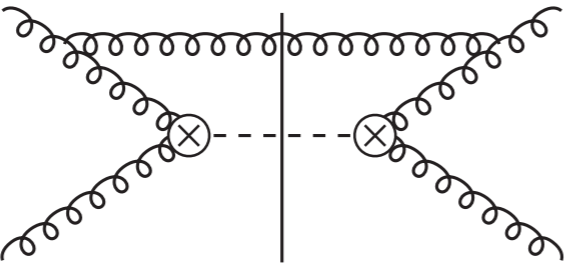
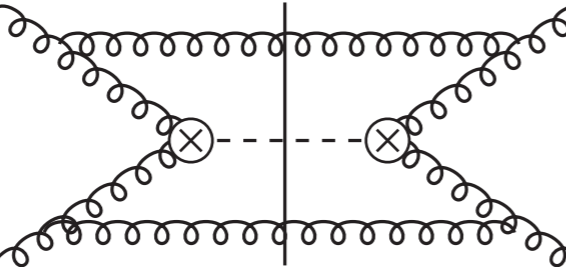
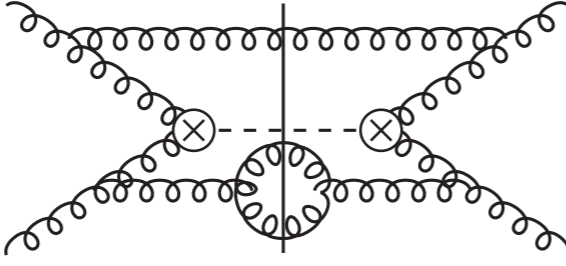
Reverse-unitarity @ N3LO

Growth in complexity for real emission

LO		1 diagram	1 integral
NLO		10 diagrams	1 integral
NNLO		381 diagrams	18 integrals
N3LO			

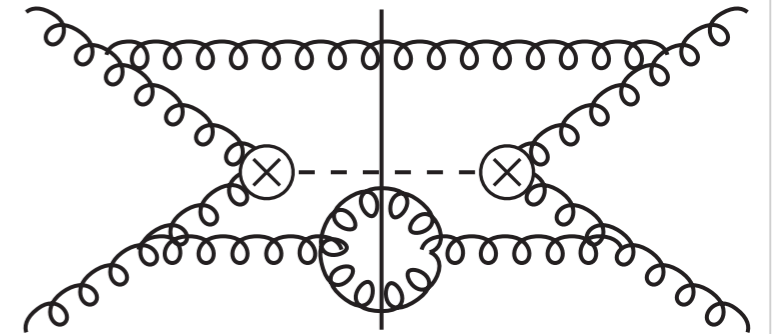
Reverse-unitarity @ N3LO

Growth in complexity for real emission

LO		1 diagram	1 integral
NLO		10 diagrams	1 integral
NNLO		381 diagrams	18 integrals
N3LO		26565 diagrams	~500 integrals

The threshold expansion

- ~ 500 master integrals only for triple real double real NNLO).
 - ➔ Tough nut to crack!
- Concentrate on some approximation first!



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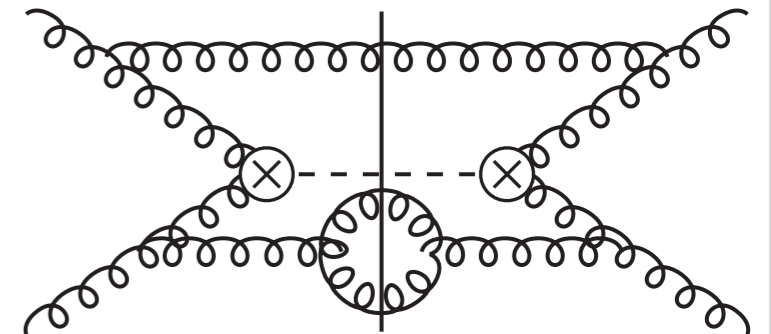
- The gluon fusion cross section depends on one single parameter:

$$z = \frac{m^2}{s} \quad \bar{z} = 1 - z$$

- Close to threshold ($z \sim 1$), we can approximate the triple real cross section by a power series:

$$\hat{\sigma}(z) = \sigma_{-1} + \sigma_0 + (1 - z) \sigma_1 + \mathcal{O}(1 - z)^2$$

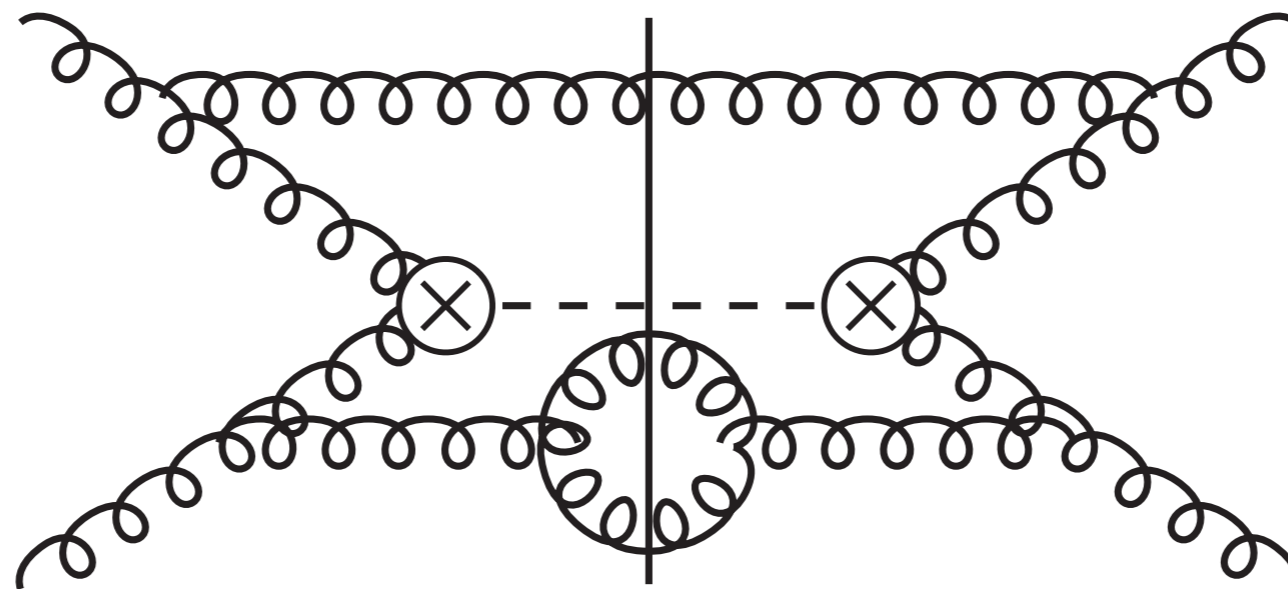
- Goal:



The threshold expansion

- Formally, this expansion corresponds to ‘expansion by regions’. [Beneke, Smirnov]
 - ➔ In the limit:
momenta are, e.g., ‘hard’, ‘soft’ or ‘collinear’.
 - ➔ Extend this to inclusive phase space.
 - ➔ Advantage:
themselves.
- Higgs production at threshold (soft-virtual):
 - ➔ Every real gluon is soft.
 - ➔ Every virtual gluon is either hard or soft.
- N.B.:
virtual and/or real gluon!
 - ➔ Universality of soft emissions!

Soft t emissions



Triple real emission

- ~500 master integrals.

- Subprocesses:

$$g g \rightarrow H g g g$$

$$g q \rightarrow H g g q$$

$$q \bar{q} \rightarrow H g g g$$

$$q q \rightarrow H g q q$$

$$g g \rightarrow H g q \bar{q}$$

$$g q \rightarrow H q \bar{q} q$$

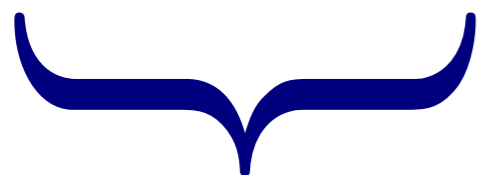
$$q \bar{q} \rightarrow H g q \bar{q}$$

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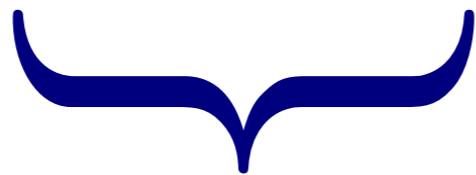
$$g q \rightarrow H Q \bar{Q} q$$

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$$q Q \rightarrow H g q Q$$



Soft-virtual



Next-to-soft-virtual

- If we concentrate on the first two terms in the expansion, all ~500 master integrals can be reduced to only 10 integrals!

NNLO example

- NNLO integral:

$$\int d\Phi_3 = \bar{z}^{3-4\epsilon} \Phi_3^S(\epsilon) \sum_{n=0}^{\infty} \frac{(1-\epsilon)_n (2-2\epsilon)_n}{(4-4\epsilon)_n} \bar{z}^n$$

$$= \bar{z}^{3-4\epsilon} \Phi_3^S(\epsilon) \left[1 + \frac{1-\epsilon}{2} \bar{z} + \frac{(1-\epsilon)(2-\epsilon)(3-2\epsilon)}{4(5-4\epsilon)} \bar{z}^2 + \mathcal{O}(\bar{z}^3) \right]$$

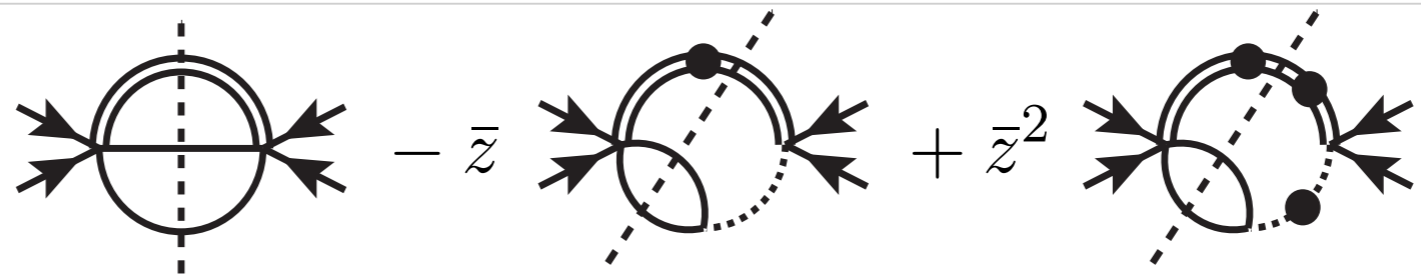
$$\Phi_3^S(\epsilon) = \frac{1}{2(4\pi)^{3-2\epsilon}} \frac{\Gamma(1-\epsilon)^2}{\Gamma(4-4\epsilon)}$$

- Diagrammatic expansion:

$$\int d\Phi_3 = \bar{z}^{3-4\epsilon} \left[\text{Diagram 1} - \bar{z} \text{Diagram 2} + \bar{z}^2 \text{Diagram 3} + \mathcal{O}(\bar{z}^3) \right]$$

➔ The coefficients themselves have a loop interpretation.

NNLO example

$$\int d\Phi_3 = \bar{z}^{3-4\epsilon} \left[\text{tree} - \bar{z} \text{NLO} + \bar{z}^2 \text{NNLO} + \mathcal{O}(\bar{z}^3) \right]$$


NNLO example

$$\int d\Phi_3 = \bar{z}^{3-4\epsilon} \left[\text{Diagram 1} - \bar{z} \text{Diagram 2} + \bar{z}^2 \text{Diagram 3} + \mathcal{O}(\bar{z}^3) \right]$$

- Using IBP identities:

$$\text{Diagram 2} = -\frac{1-\epsilon}{2} \text{Diagram 1},$$

$$\text{Diagram 3} = \frac{(1-\epsilon)(2-\epsilon)(3-2\epsilon)}{4(5-4\epsilon)} \text{Diagram 1}$$

NNLO example

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- To be compared with exact result:

$$\bar{z}^{3-4\epsilon} \Phi_3^S(\epsilon) \left[1 + \frac{1-\epsilon}{2} \bar{z} + \frac{(1-\epsilon)(2-\epsilon)(3-2\epsilon)}{4(5-4\epsilon)} \bar{z}^2 + \mathcal{O}(\bar{z}^3) \right]$$

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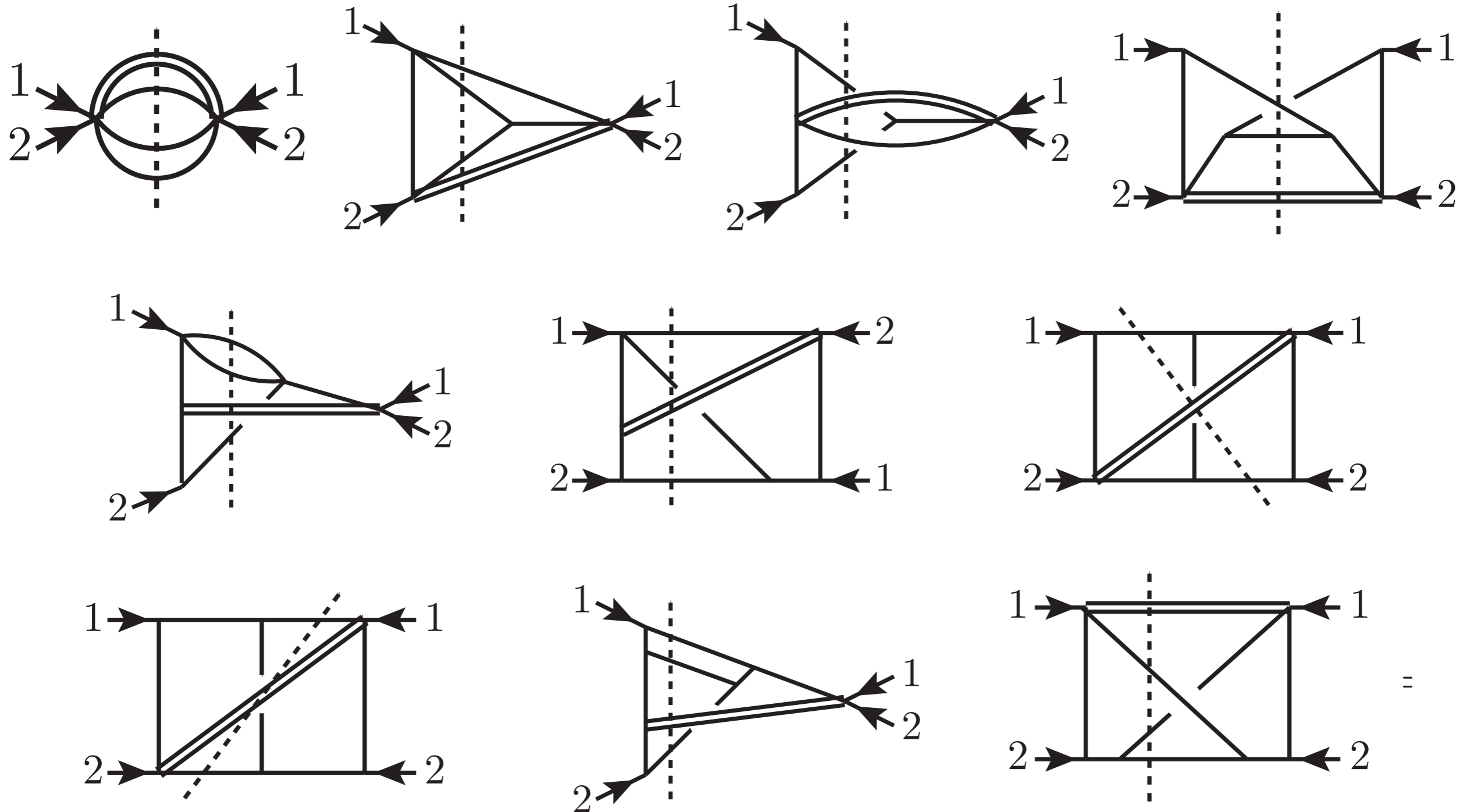
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Soft triple real emissions



[Anastasiou, Dulat, CD, Mistlberger]

Soft triple real emissions

- We were able to compute all the master integrals analytically.
- General strategy:
 - ➔ There is a canonical way to turn each of these integrals into a Mellin-Barnes integral.

$$\mathcal{F}_6(\epsilon) = \frac{\Gamma(6 - 6\epsilon)}{\epsilon \Gamma(1 - 6\epsilon) \Gamma(1 - \epsilon)^2} \int_{-i\infty}^{+i\infty} \frac{dz_1 dz_2}{(2\pi i)^2} \Gamma(-z_1) \Gamma(z_1 + 1) \Gamma(-z_2) \Gamma(z_2 + 1) \\ \times \frac{\Gamma(-\epsilon + z_1 - z_2) \Gamma(z_2 - \epsilon) \Gamma(-2\epsilon - z_1 + z_2) \Gamma(-\epsilon - z_1 + z_2)}{\Gamma(-\epsilon + z_2 + 1) \Gamma(-2\epsilon - z_1 + z_2 + 1)}.$$

- ➔ All of the integrals can be computed as a Laurent series in dimensional regularization.
- ➔ One of the integrals required use of symbols and coproducts

Soft triple real emissions

$$\begin{aligned}\mathcal{F}_6(\epsilon) = & \frac{10}{\epsilon^5} - \frac{137}{\epsilon^4} + \frac{1}{\epsilon^3} \left(40 \zeta_2 + 675 \right) + \frac{1}{\epsilon^2} \left(320 \zeta_3 - 548 \zeta_2 - 1530 \right) \\ & + \frac{1}{\epsilon} \left(1500 \zeta_4 - 4384 \zeta_3 + 2700 \zeta_2 + 1620 \right) + 5160 \zeta_5 + 320 \zeta_2 \zeta_3 - 20550 \zeta_4 \\ & + 21600 \zeta_3 - 6120 \zeta_2 - 648 + \epsilon \left(18340 \zeta_6 + 1280 \zeta_3^2 - 70692 \zeta_5 - 4384 \zeta_2 \zeta_3 \right. \\ & \left. + 101250 \zeta_4 - 48960 \zeta_3 + 6480 \zeta_2 \right) + \mathcal{O}(\epsilon^2).\end{aligned}$$

- Intriguing observation:

→ All
all

- How can we be sure that we have obtained the correct results..?

Soft triple real emissions

- We can compute the Mellin-Barnes integrals numerically and compare to our analytic results.

Soft triple real emissions

- We can compute the Mellin-Barnes integrals numerically and compare to our analytic results.
- The integrals in four dimensions are related to the integrals in six dimensions:

$$\mathcal{F}_6(D+2)\mathcal{R} = \frac{(4256 - 6684D + 4224D^2 - 1345D^3 + 216D^4 - 14D^5)}{3(D-4)^2(D-3)(D-2)^2} + \frac{(D-4)(3D-10)}{9(D-2)^2(3D-7)}\mathcal{F}_2(D) - \frac{(D-4)^3}{24(D-2)(3D-11)(3D-7)}\mathcal{F}_6(D),$$

- ➔ Similar to dimensional shift identities for loops. [Tarasov]
- The integrals are finite in six dimensions.
 - ➔ Strong constraint!

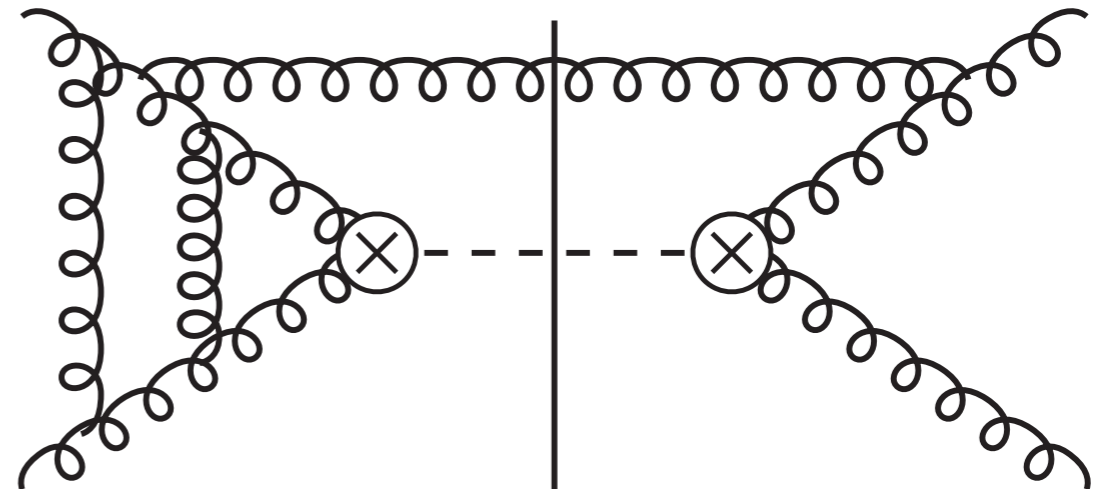
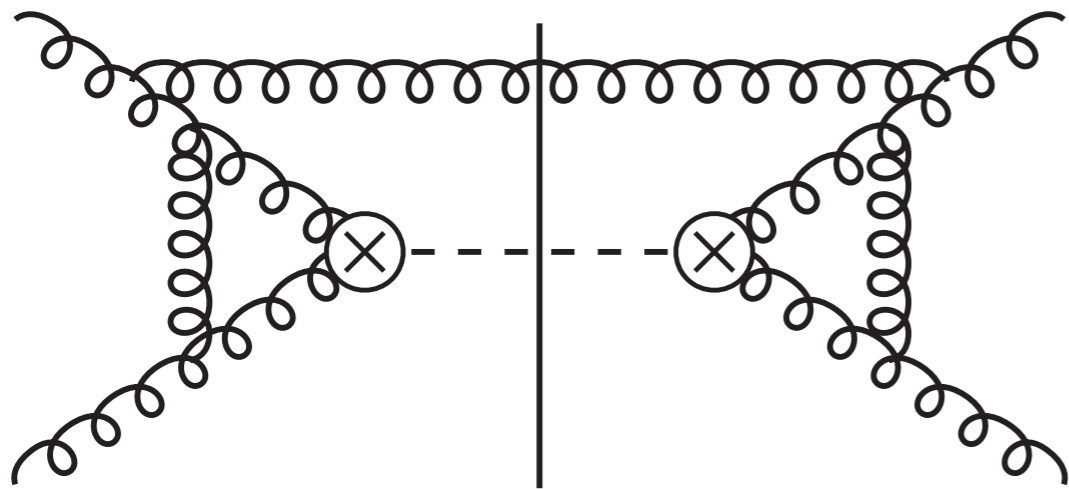
Soft triple real cross sections

- The integrals immediately allow us to write down the first two terms in the soft expansion of the cross section, e.g.,

$$\begin{aligned}
 \sigma_{gg \rightarrow H+gq\bar{q}}^{S(0)} &= \frac{2^5}{3^7} \frac{1}{8(N_c^2 - 1)^2} (4\pi\alpha_S)^3 \Phi_4^S(\epsilon) C_A C_F c_H^2 N_f \\
 &\times \left\{ \frac{153090}{\epsilon^4} - \frac{1604043}{\epsilon^3} + \frac{1}{\epsilon^2} (-29160\zeta_2 + 4903902) \right. \\
 &+ \frac{1}{\epsilon} (-204120\zeta_3 + 321732\zeta_2 - 4833675) - 874800\zeta_4 + 2252124\zeta_3 - 911088\zeta_2 \\
 &+ 203535 + \epsilon(-2711880\zeta_5 - 233280\zeta_2\zeta_3 + 9651960\zeta_4 - 6290136\zeta_3 - 492210\zeta_2 \\
 &+ 1667109) + \epsilon^2(-9360360\zeta_6 - 816480\zeta_3^2 + 29921076\zeta_5 + 2573856\zeta_2\zeta_3 \\
 &- 26589060\zeta_4 - 4323186\zeta_3 + 4693212\zeta_2 + 1294731) \\
 &+ 2C_A C_F \left[\frac{167670}{\epsilon^4} - \frac{1743039}{\epsilon^3} + \frac{1}{\epsilon^2} (-29160\zeta_2 + 5267592) + \frac{1}{\epsilon} (-204120\zeta_3 \right. \\
 &+ 321732\zeta_2 - 5183163) - 874800\zeta_4 + 2252124\zeta_3 - 911088\zeta_2 + 337959 \\
 &+ \epsilon(-2711880\zeta_5 - 233280\zeta_2\zeta_3 + 9651960\zeta_4 - 6290136\zeta_3 - 492210\zeta_2 + 1651749) \\
 &+ \epsilon^2(-9360360\zeta_6 - 816480\zeta_3^2 + 29921076\zeta_5 + 2573856\zeta_2\zeta_3 - 26589060\zeta_4 \\
 &\left. - 4323186\zeta_3 + 4693212\zeta_2 + 1284491) \right] + \mathcal{O}(\epsilon^3) \left. \right\}.
 \end{aligned}$$

[Anastasiou, Dulat,
CD, Mistlberger]

Soft double-virtual-real emissions



Double-virtual-real emission

- Subprocesses:

$$g g \rightarrow H g$$


Soft-virtual

$$g q \rightarrow H q$$


Next-to-soft-virtual

$$q \bar{q} \rightarrow H g$$

- The phase space is trivial (2-body phase space).

$$|\mathcal{M}^{(L)}(g g \rightarrow H g)\rangle = \sum_{k=0}^L \varepsilon^\mu J_\mu^{a(k)} |\mathcal{M}^{(L-k)}(g g \rightarrow H)\rangle$$

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- If the final-state gluon is soft into the soft current:

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- The soft current had previously been computed

➔ at one loop.

[Catani, Grazzini]

➔ at two loops, up to finite terms.

[Badger, Glover]

- At N3LO, we need higher-order terms at two loops.

The two-loop soft current

- Two parallel computations of these higher-order terms (for the interference with the Born soft current):
 - ➔ Two-loop Wilson line computation up to weight 6. [Li, Zhu]
 - ➔ Extraction from the two-loop matrix element for $\gamma^* \rightarrow q \bar{q} g$ to all orders in epsilon. [CD, Gehrmann]

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- **Method:**

integrals as an expansion in the soft limit while keeping the coefficients exact in epsilon:

$$F_i(y, z; \epsilon) = \sum_{m,n=0}^2 y^{-m\epsilon} z^{-n\epsilon} f_{i,mn}(y, z; \epsilon),$$

$$f_{i,mn}(y, z; \epsilon) = \sum_{k=r_y}^{\infty} \sum_{l=r_z}^{\infty} c_{i,mn}^{kl}(\epsilon) y^k z^l$$

The two-loop soft current

- The two-loop soft current to all orders in dimensional regularization

$$r_{soft}^{(2)} = N N_f R_1(\epsilon) + N^2 R_2(\epsilon),$$

$$R_1(\epsilon) = \frac{2\Gamma(-2\epsilon)}{(1+\epsilon)\Gamma(4-2\epsilon)} \frac{\Gamma(1-2\epsilon)^2 \Gamma(1+2\epsilon)^2}{\Gamma(1-\epsilon)^2 \Gamma(1+\epsilon)^2} \left[3 \frac{\Gamma(1-\epsilon)\Gamma(1-2\epsilon)}{\Gamma(1-3\epsilon)} - \frac{(1+\epsilon^3)}{\epsilon^2(1+\epsilon)} \frac{\Gamma(1-2\epsilon)^2}{\Gamma(1-4\epsilon)} \right],$$

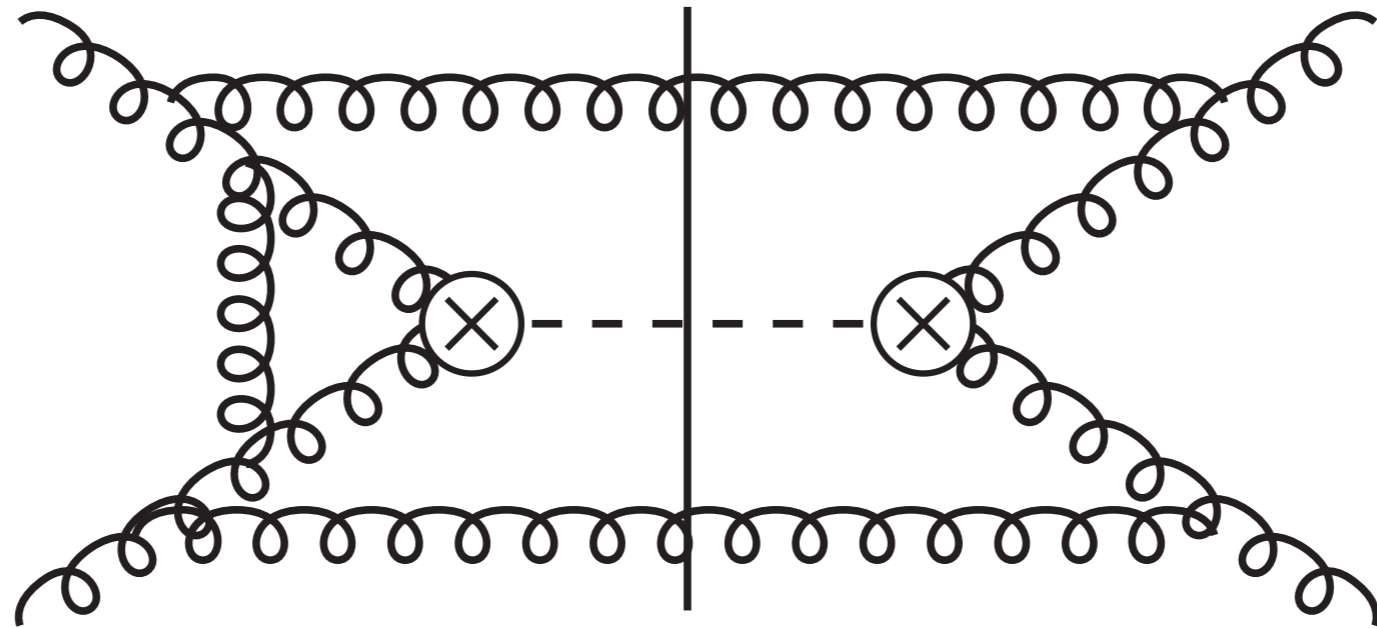
$$R_2(\epsilon) = \frac{\Gamma(1-2\epsilon)^3 \Gamma(1+2\epsilon)^2}{6\epsilon^4 \Gamma(1-\epsilon)\Gamma(1+\epsilon)^2 \Gamma(1-3\epsilon)} \left\{ (1+4\epsilon) {}_4F_3(1, 1, 1-\epsilon, -4\epsilon; 2, 1-3\epsilon, 1-2\epsilon; 1) \right. \\ \left. - 6\epsilon [\psi(1-3\epsilon) + \psi(1-2\epsilon) - \psi(1-\epsilon) - \psi(1+\epsilon)] + \frac{(14\epsilon^3 + 4\epsilon^2 + 5\epsilon - 3)}{2(1+\epsilon)(3-2\epsilon)(1-2\epsilon)} \right\} \\ + \frac{(1+4\epsilon)}{3\epsilon^4(1+2\epsilon)} \frac{\Gamma(1-2\epsilon)^4 \Gamma(1+2\epsilon)^2}{\Gamma(1-\epsilon)^2 \Gamma(1+\epsilon)^2 \Gamma(1-4\epsilon)} \left\{ 2 {}_3F_2(1, -2\epsilon, 2\epsilon+1; 1-\epsilon, 2\epsilon+2; 1) \right. \\ \left. - \frac{\Gamma(1+\epsilon)\Gamma(1-2\epsilon)}{\Gamma(1-\epsilon)} {}_3F_2(-2\epsilon, \epsilon+1, 2\epsilon+1; 1-\epsilon, 2\epsilon+2; 1) + \frac{(1+2\epsilon)(6\epsilon^4 + 13\epsilon^3 - 16\epsilon^2 - 38\epsilon + 3)}{4(1+4\epsilon)(1+\epsilon)(3-2\epsilon)(1-2\epsilon)} \right\},$$

- The soft-virtual RVV contribution to Higgs@N3LO is easily obtained from this by multiplying by the phase space.

Real-virtual squared

- Contribution from one-loop-squared can easily be computed exactly.
- We did the computation in four different ways:
 - ➔ Threshold expansion by expanding hypergeometric functions.
 - ➔ Threshold expansion from expansion by regions.
 - ➔ Reverse-unitarity and differential equations.
 - ➔ Direct integration of the matrix element over phase space.
[Anastasiou, CD, Dulat, Herzog, Mistlberger]
- Confirmed by independent computation. [Kilgore]
- Full two-loop matrix element is also known. [Glover, Gehrmann, Jaquier, Koukoutsakis]
 - ➔ Can be done in the same way, but need two-loop collinear counterterms.

Soft virtual-double-real emissions



Virtual-double-real emission

- Subprocesses:

$$g g \rightarrow H g g$$

$$g q \rightarrow H g q$$

$$q \bar{q} \rightarrow H g g$$

$$q q \rightarrow H q q$$

$$g g \rightarrow H q \bar{q}$$

$$q \bar{q} \rightarrow H q \bar{q}$$

$$q Q \rightarrow H q Q$$

$$q \bar{q} \rightarrow H Q \bar{Q}$$

$$q \bar{Q} \rightarrow H q \bar{Q}$$

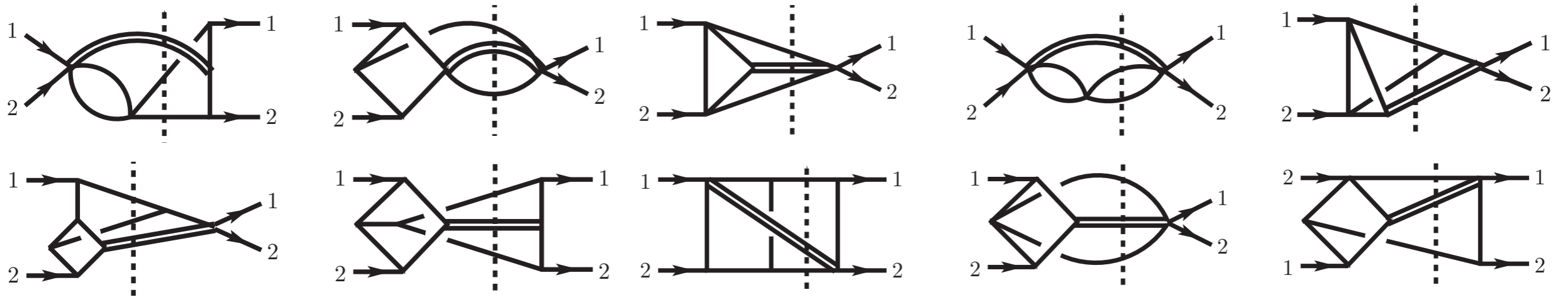
Virtual-double-real emission

- Subprocesses:

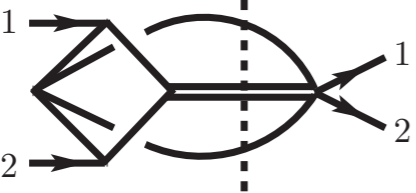
$$\begin{array}{llll} g g \rightarrow H g g & g q \rightarrow H g q & q \bar{q} \rightarrow H g g & q q \rightarrow H q q \\ g g \rightarrow H q \bar{q} & & q \bar{q} \rightarrow H q \bar{q} & q Q \rightarrow H q Q \\ & & q \bar{q} \rightarrow H Q \bar{Q} & q \bar{Q} \rightarrow H q \bar{Q} \end{array}$$

- The soft-virtual term receives contributions from two regions:
 - ➔ The virtual gluon is hard.
 - ➔ The virtual gluon is soft.
- The hard region is trivial (tree-level emission of two soft gluons).
- The soft region can be dealt with in a way similar to the soft triple real emission.
 - ➔ IBP reduction in soft limit and soft master integrals.

VRR soft master integrals



- All integrals can be computed by combining the soft expansion for the virtuals with the phase space techniques developed for the triple real emission. [Anastasiou, CD, Dulat, Furlan, Herzog, Mistlberger

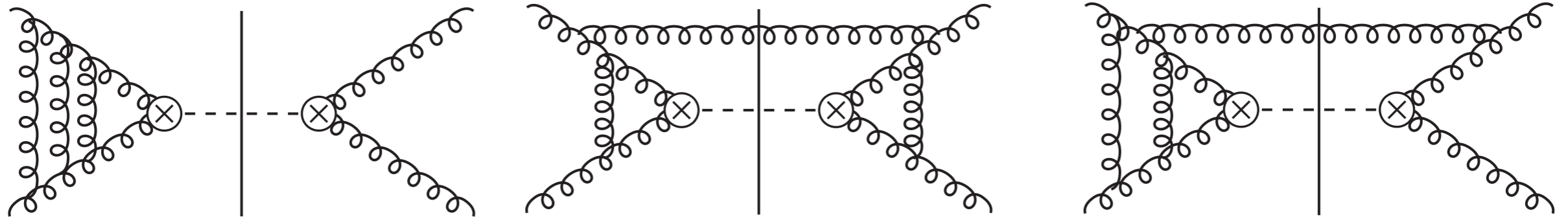


$$\begin{aligned}
 &= -\frac{4\Gamma(4-4\epsilon)\Gamma(1-3\epsilon)}{\epsilon(1+\epsilon)(1-2\epsilon)\Gamma(3-6\epsilon)\Gamma(1-\epsilon)} {}_3F_2(1, 1, 1-\epsilon; 2-3\epsilon, 2+\epsilon; 1) \\
 &= -\frac{12\zeta_2}{\epsilon} - 8\zeta_2 - 36\zeta_3 + (-112\zeta_2 - 24\zeta_3 + 33\zeta_4)\epsilon + (720\zeta_3\zeta_2 - 672\zeta_2 \\
 &\quad - 336\zeta_3 + 22\zeta_4 - 450\zeta_5)\epsilon^2 + (1512\zeta_3^2 + 480\zeta_2\zeta_3 - 2016\zeta_3 - 4032\zeta_2 + 308\zeta_4 \\
 &\quad - 300\zeta_5 + \frac{16881}{4}\zeta_6)\epsilon^3 + \mathcal{O}(\epsilon^4).
 \end{aligned}$$

- Results recently confirmed by independent computation using Wilson lines. [Li, von Manteuffel, Schabinger, Zhu

The gluon-fusion cross
section in the soft limit

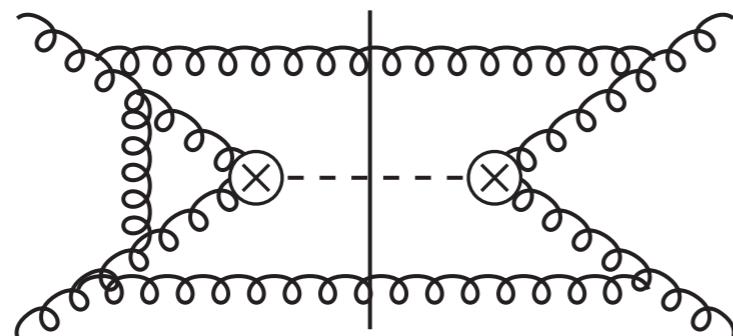
N3LO status: soft-virtual



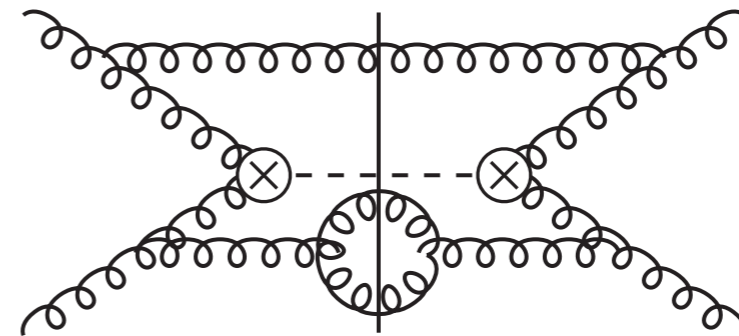
✓ Triple virtual

✓ Real-virtual squared

✓ Double virtual real



✓ Double real virtual



✓ Triple real

✓ +

The soft-virtual approximation

- The

$$\hat{\sigma}(z) = \boxed{\sigma_{-1}} + \sigma_0 + (1 - z) \sigma_1 + \mathcal{O}(1 - z)^2$$

- The soft-virtual term receives contributions from a 'pole' at $z \sim 1$:

$$(1 - z)^{-1+n\epsilon} = \frac{\delta(1 - z)}{n\epsilon} + \left[\frac{1}{1 - z} \right]_+ + n\epsilon \left[\frac{\log(1 - z)}{1 - z} \right]_+ + \mathcal{O}(\epsilon^2)$$

- Plus-distribution terms already known. [Moch, Vogt
- Complete three-loop corrections are contained the delta function term.
 - ➔ The soft-virtual term contains the complete three-loop corrections plus the correction from the emission of up to three soft gluons.

The soft-virtual approximation

- At NLO and NNLO, the soft-virtual term reads ($\mu_R = \mu_F = m_H$)

$$\hat{\sigma}_{gg}^{SV}(z) = \frac{\pi C(\mu^2)^2}{v^2 (N^2 - 1)^2} \sum_{k=0}^{\infty} \left(\frac{\alpha_s}{\pi}\right)^k \hat{\eta}^{(k)}(z)$$

$$\hat{\eta}^{(0)}(z) = \delta(1 - z) \qquad \hat{\eta}^{(1)}(z) = 2 C_A \zeta_2 \delta(1 - z) + 4 C_A \left[\frac{\log(1 - z)}{1 - z} \right]_+$$

$$\begin{aligned} \hat{\eta}^{(2)}(z) = & \delta(1 - z) \left\{ C_A^2 \left(\frac{67}{18} \zeta_2 - \frac{55}{12} \zeta_3 - \frac{1}{8} \zeta_4 + \frac{93}{16} \right) + N_F \left[C_F \left(\zeta_3 - \frac{67}{48} \right) - C_A \left(\frac{5}{9} \zeta_2 + \frac{1}{6} \zeta_3 + \frac{5}{3} \right) \right] \right\} \\ & + \left[\frac{1}{1 - z} \right]_+ \left[C_A^2 \left(\frac{11}{3} \zeta_2 + \frac{39}{2} \zeta_3 - \frac{101}{27} \right) + N_F C_A \left(\frac{14}{27} - \frac{2}{3} \zeta_2 \right) \right] \\ & + \left[\frac{\log(1 - z)}{1 - z} \right]_+ \left[C_A^2 \left(\frac{67}{9} - 10 \zeta_2 \right) - \frac{10}{9} C_A N_F \right] \\ & + \left[\frac{\log^2(1 - z)}{1 - z} \right]_+ \left(\frac{2}{3} C_A N_F - \frac{11}{3} C_A^2 \right) + \left[\frac{\log^3(1 - z)}{1 - z} \right]_+ 8 C_A^2. \end{aligned}$$

The soft-virtual approximation

- All the integrals can be computed analytically!
 - ➔ 22 three-loop. [Baikov, Chetyrkin, Smirnov, Smirnov, Steinhauser; Gehrmann, Glover, Huber, Ikizlerli, Studerus]
 - ➔ 3 double-virtual-real. [CD, Gehrmann; Li, Zhu]
 - ➔ 7 real-virtual-squared. [Anastasiou, CD, Dulat, Herzog, Mistlberger; Kilgore]
 - ➔ 10 virtual-double-real. [Anastasiou, CD, Dulat, Furlan, Herzog, Mistlberger; Li, von Manteuffel, Schabinger, Zhu]
 - ➔ 8 triple real. [Anastasiou, CD, Dulat, Mistlberger]
- In addition, one needs:
 - ➔ three-loop splitting functions. [Moch, Vogt, Vermaseren]
 - ➔ three-loop beta function. [Tarasov, Vladimirov, Zharkov; Larin, Vermaseren]
 - ➔ three-loop Wilson coefficient. [Chetyrkin, Kniehl, Steinhauser; Schroder, Steinhauser; Chetyrkin, Kuhn, Sturm]

Higgs soft-virtual @ N3LO

$$\begin{aligned}
\hat{\eta}^{(3)}(z) = & \delta(1-z) \left\{ C_A^3 \left(-\frac{2003}{48} \zeta_6 + \frac{413}{6} \zeta_3^2 - \frac{7579}{144} \zeta_5 + \frac{979}{24} \zeta_2 \zeta_3 - \frac{15257}{864} \zeta_4 - \frac{819}{16} \zeta_3 + \frac{16151}{1296} \zeta_2 + \frac{215131}{5184} \right) \right. \\
& + N_F \left[C_A^2 \left(\frac{869}{72} \zeta_5 - \frac{125}{12} \zeta_3 \zeta_2 + \frac{2629}{432} \zeta_4 + \frac{1231}{216} \zeta_3 - \frac{70}{81} \zeta_2 - \frac{98059}{5184} \right) \right. \\
& \quad \left. + C_A C_F \left(\frac{5}{2} \zeta_5 + 3 \zeta_3 \zeta_2 + \frac{11}{72} \zeta_4 + \frac{13}{2} \zeta_3 - \frac{71}{36} \zeta_2 - \frac{63991}{5184} \right) + C_F^2 \left(-5 \zeta_5 + \frac{37}{12} \zeta_3 + \frac{19}{18} \right) \right] \\
& + N_F^2 \left[C_A \left(-\frac{19}{36} \zeta_4 + \frac{43}{108} \zeta_3 - \frac{133}{324} \zeta_2 + \frac{2515}{1728} \right) + C_F \left(-\frac{1}{36} \zeta_4 - \frac{7}{6} \zeta_3 - \frac{23}{72} \zeta_2 + \frac{4481}{2592} \right) \right] \left. \right\} \\
& + \left[\frac{1}{1-z} \right]_+ \left\{ C_A^3 \left(186 \zeta_5 - \frac{725}{6} \zeta_3 \zeta_2 + \frac{253}{24} \zeta_4 + \frac{8941}{108} \zeta_3 + \frac{8563}{324} \zeta_2 - \frac{297029}{23328} \right) + N_F^2 C_A \left(\frac{5}{27} \zeta_3 + \frac{10}{27} \zeta_2 - \frac{58}{729} \right) \right. \\
& \quad \left. + N_F \left[C_A^2 \left(-\frac{17}{12} \zeta_4 - \frac{475}{36} \zeta_3 - \frac{2173}{324} \zeta_2 + \frac{31313}{11664} \right) + C_A C_F \left(-\frac{1}{2} \zeta_4 - \frac{19}{18} \zeta_3 - \frac{1}{2} \zeta_2 + \frac{1711}{864} \right) \right] \right\} \\
& + \left[\frac{\log(1-z)}{1-z} \right]_+ \left\{ C_A^3 \left(-77 \zeta_4 - \frac{352}{3} \zeta_3 - \frac{152}{3} \zeta_2 + \frac{30569}{648} \right) + N_F^2 C_A \left(-\frac{4}{9} \zeta_2 + \frac{25}{81} \right) \right. \\
& \quad \left. + N_F \left[C_A^2 \left(\frac{46}{3} \zeta_3 + \frac{94}{9} \zeta_2 - \frac{4211}{324} \right) + C_A C_F \left(6 \zeta_3 - \frac{63}{8} \right) \right] \right\} \\
& + \left[\frac{\log^2(1-z)}{1-z} \right]_+ \left\{ C_A^3 \left(181 \zeta_3 + \frac{187}{3} \zeta_2 - \frac{1051}{27} \right) + N_F \left[C_A^2 \left(-\frac{34}{3} \zeta_2 + \frac{457}{54} \right) + \frac{1}{2} C_A C_F \right] - \frac{10}{27} N_F^2 C_A \right\} \\
& + \left[\frac{\log^3(1-z)}{1-z} \right]_+ \left\{ C_A^3 \left(-56 \zeta_2 + \frac{925}{27} \right) - \frac{164}{27} N_F C_A^2 + \frac{4}{27} N_F^2 C_A \right\} \\
& + \left[\frac{\log^4(1-z)}{1-z} \right]_+ \left(\frac{20}{9} N_F C_A^2 - \frac{110}{9} C_A^3 \right) + \left[\frac{\log^5(1-z)}{1-z} \right]_+ 8 C_A^3.
\end{aligned}$$

[Anastasiou, CD, Dulat, Furlan, Gehrmann, Herzog, Mistlberger]

Conclusion & Outlook

- We have completed the first computation of the Higgs boson cross section at N³LO in the soft-virtual approximation.
- ‘Amplitude-technology’ played a crucial role in the computation of the integrals
- More terms in the expansion/full computation in progress.
- Soft term already allows to extract interesting results:
 - ➔ Extension to Drell-Yan @ N³LO in the soft limit.
[Ahmed, Mahakhud, Mathews, Rana, Ravindran]
 - ➔ Extension to rapidity distribution @ N³LO in soft limit.
[Ahmed, Mandal, Rana, Ravindran]
 - ➔ Soft-gluon resummation at N³LL.
[Bonvini, Marzani; Catani, Cieri, de Florian, Ferrara, Grazzini]

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Back up

Higgs soft-virtual @ N³LO

- How can we be sure that we got it right?
 - ➔ Plus-distribution terms agree with Moch & Vogt.
 - ➔ All master integrals were computed analytically and cross checked numerically.
 - ➔ Independent computations of matrix elements and integrals.
 - All but the triple-real contribution have been confirmed by other groups.

Higgs soft-virtual @ N3LO

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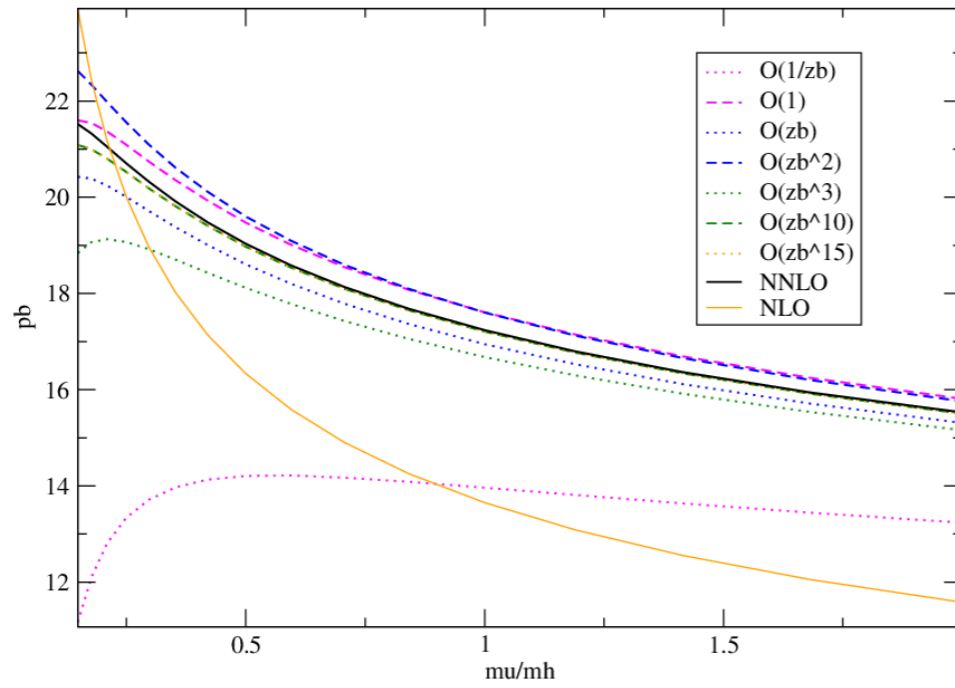
- **Caveat emptor!**

$$\int dx_1 dx_2 [f_i(x_1) f_j(x_2) z g(z)] \left[\frac{\hat{\sigma}_{ij}(s, z)}{z g(z)} \right]_{\text{threshold}} \quad \lim_{z \rightarrow 1} g(z) = 1.$$

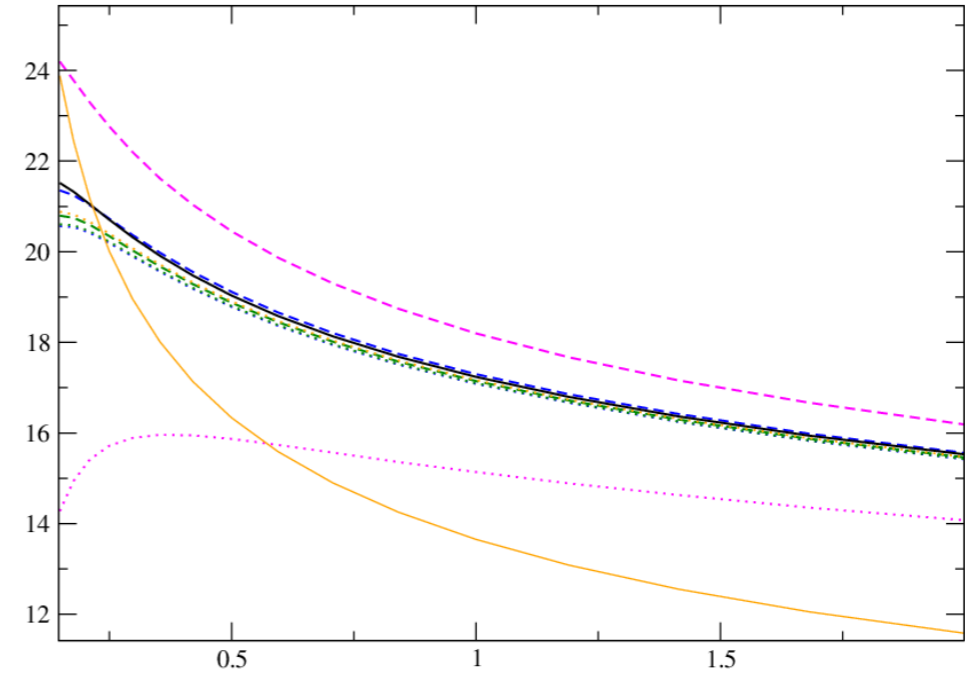
- ➔ Formally all these choices are equivalent!

Higgs soft-virtual @ N3LO

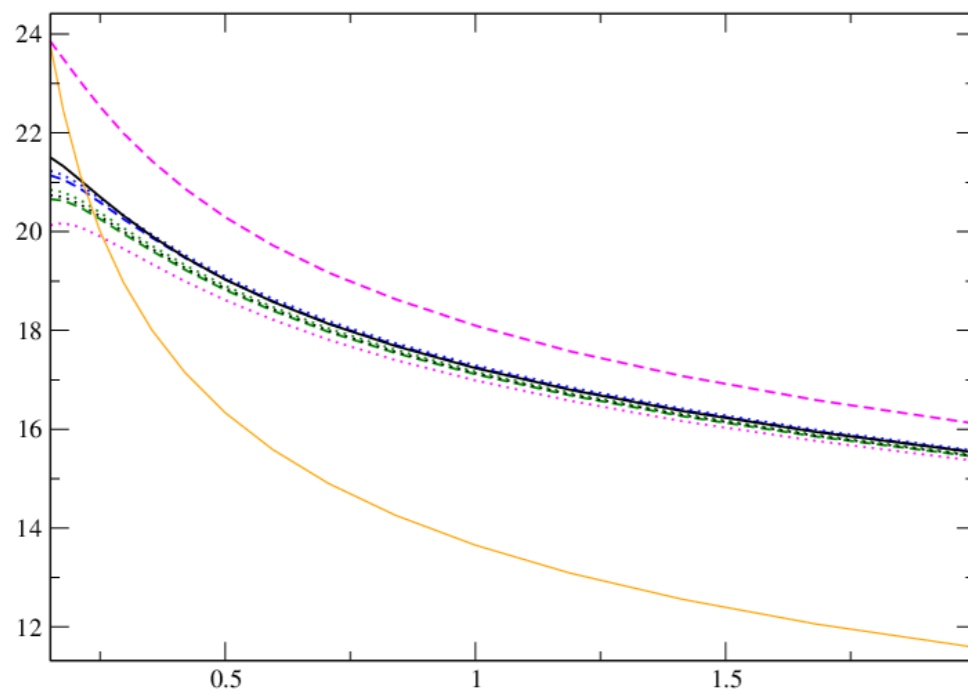
NNLO Soft Expansion $g(z)=1/z$



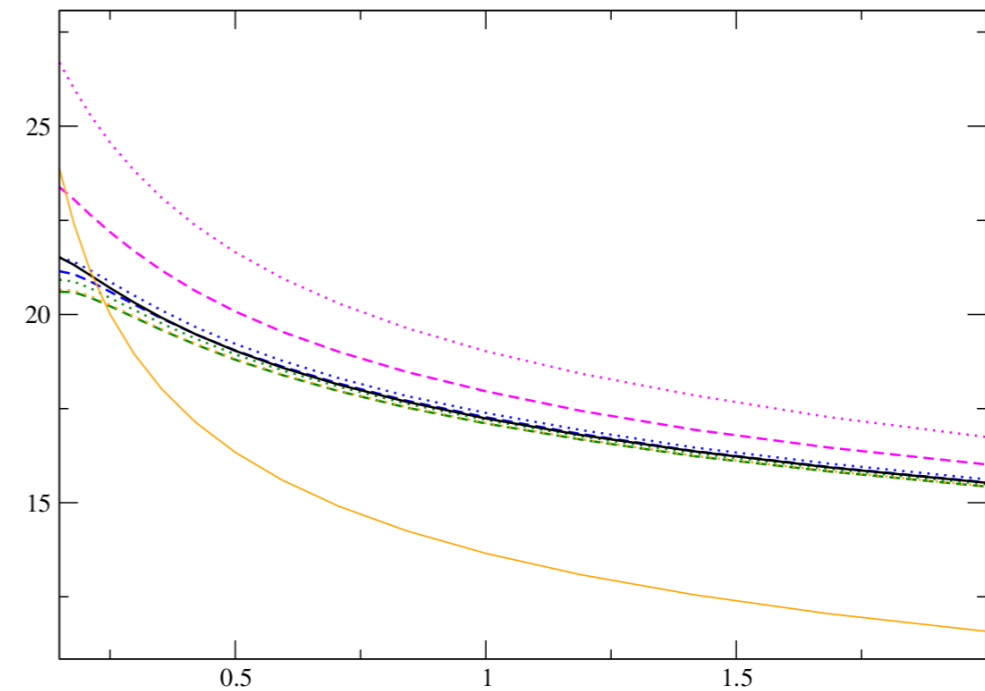
Soft Expansion NNLO $g(z)=1$



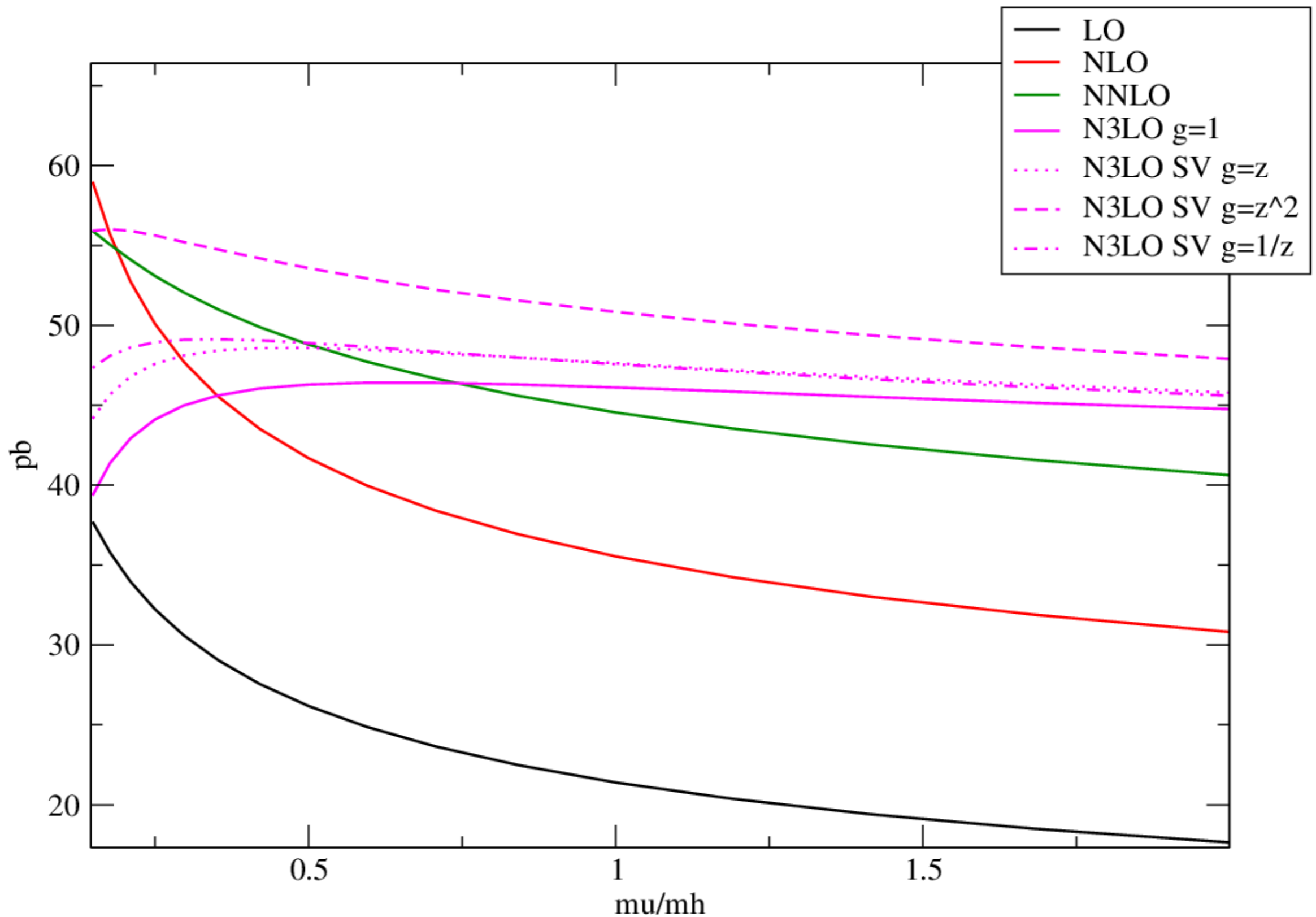
Soft Expansion NNLO $g(z)=z$



NNLO $g(z)=z^2$



Higgs soft-virtual @ N3LO



Higgs soft-virtual @ N3LO

--- 7 TeV — 8 TeV - - - 13 TeV — 14 TeV - - - 100 TeV

