

# MODULAR GRAPH FUNCTIONS

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ABSTRACT. We consider properties of *modular graph functions*, which are non-holomorphic modular functions associated with the Feynman graphs for a conformal scalar field theory on a two-dimensional torus. Such functions arise, for example, in the low energy expansion of genus-one Type II superstring amplitudes. We demonstrate that these functions are sums, with rational coefficients, of special values of *single-valued elliptic multiple polylogarithms*, which will be introduced in this paper. This insight suggests the many interrelations between these modular graph functions (a few of which were motivated in an earlier paper) may be obtained as a consequence of identities involving elliptic polylogarithms.

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## 1. INTRODUCTION

Superstring perturbation theory may be merely an approximation to a complete non-perturbative formulation of string theory, but it already exhibits a remarkably rich mathematical structure. The perturbative series is given by a topological expansion for two-dimensional surfaces which represent string world-sheets. For the closed superstring theories the perturbative series is given by the sum over all genera  $g$ , with

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*Date:* December 22, 2015.

DAMTP-2015-86; IPhT-t15/202, IHES/P/15/29.

$g \geq 0$ , of functional integrals over orientable (super)Riemann surfaces. In the case of four-graviton amplitudes, which will be the prototype and point of departure of this paper, the explicit expressions for tree-level ( $g = 0$ ) and genus-one ( $g = 1$ ) were obtained in [1], for genus-two ( $g = 2$ ) in [2], and a certain amount is known about the leading low energy behavior of the genus-three ( $g = 3$ ) case in [3]. For the open string theory, a summation over boundaries and cross-caps must also be included.

The low energy expansion of string theory corresponds to an expansion valid when the energies and momenta are small in units of the inverse of the string length scale  $\ell_s$ , a parameter which is related to Newton's gravitational constant. The lowest order contribution corresponds to Einstein's theory, while higher order corrections become important for strong gravitational fields. The structure of these higher order corrections is of considerable mathematical interest. In particular, their coefficients in the low energy expansion of tree-level  $N$ -particle amplitudes in open superstring theory are multiple zeta values (MZVs), which are special values of multiple polylogarithms (MPLs). The analogous coefficients of tree-level  $N$ -particle closed superstring amplitudes are single-valued multiple zeta values (svMZVs) [4,5], which in turn are special values of single-valued multiple polylogarithms (svMPLs), following the terminology introduced in [6–8].

Much less is understood about the low energy expansion of string amplitudes with higher genus. In the genus-one closed superstring case the coefficients in the low energy expansion are given by integrals of non-holomorphic modular functions over the complex structure modulus  $\tau$  of the torus that is defined by the string world-sheet. These modular functions, which can be expressed as multiple sums, are generalizations of non-holomorphic Eisenstein series. Specific classes of such functions have been shown to satisfy a number of very intriguing relationships [9,10]. These are reminiscent of the algebraic relationships between multiple zeta values but, in the present case, they are relationships between functions defined on an elliptic curve. In the genus-two closed superstring case, a connection has been uncovered with the Zhang-Kawazumi invariant [11], which satisfies equally intriguing relations [12], but whose study remains incomplete.

In the present paper, we specialize to the case of a conformal scalar field theory on a two-dimensional torus. To every Feynman graph with vanishing external momenta on each vertex,  $L$  loops, and  $w$  scalar Green functions on the edges of the graph, we associate a non-holomorphic modular function of depth  $L$  and weight  $w$ , which we shall refer to as a *modular graph function*. Certain classes of graphs will produce vanishing modular functions, such as any one-edge reducible graph, and any graph in which at least one vertex supports only a single Green function edge. In this paper, we shall consider graphs with non-derivative couplings only, but this restriction can be easily lifted if needed.

We shall also introduce *single-valued elliptic multiple polylogarithms* via Feynman graphs in which all vertices but two have vanishing external momenta. These functions are of the elliptic type. One purpose of this paper is to demonstrate that a modular graph function may be expressed as the value of a single-valued elliptic multiple polylogarithm (sveMPL)<sup>1</sup> when its argument is set equal to 1. This is an elliptic analogue of the familiar statement that the single-valued multiple-zeta values discussed in [6–8] are the values of single-valued multiple polylogarithms with their arguments set equal to 1. This connection between modular graph functions and single-valued elliptic multiple polylogarithms suggests a compelling origin of the many interrelations between modular graph functions (a few of which were motivated in [9]) as a consequence of identities involving elliptic polylogarithms.

**1.1. Outline of paper.** Section 2 will give a brief overview of some of the relevant features of polylogarithms, multiple polylogarithms, multiple zeta values, their single-valued projections, and their elliptic generalisations that will enter into the subsequent ideas in the paper. Section 3 will discuss the modular graph functions and single-valued elliptic multiple polylogarithms that arise in the low energy expansion of the perturbative amplitudes in superstring theory and which are expressed in terms of Feynman graphs for a conformal scalar field theory on a two-dimensional torus. In particular, we will show that every modular graph function is given by a single valued elliptic multiple polylogarithm evaluated at a special point. Section 4 illustrates this feature by considering some of the infinite classes of graphs studied in [9, 10, 16, 17] by other methods. Section 5 develops the conjecture which states that single-valued elliptic multiple polylogarithms are linear combinations of elliptic polylogarithms with rational coefficients, offers a proof for the infinite class of star graphs, and outlines some of the arguments for general graphs. In section 6, the validity of the conjecture is shown to lead to a corollary stating that the non-leading coefficients of the Laurent expansion of the constant Fourier mode of modular graph functions are single-valued multiple zeta values. A summary and further thoughts on the basis of modular graph functions, and their further generalisations, is relayed to section 7.

Appendix A presents some relations between Eichler integrals, elliptic polylogarithms, and holomorphic Eisenstein series in preparation of appendix B where the algebraic properties of multiple polylogarithms will be used to evaluate the modular graph function associated with the simplest two-loop graph,  $C_{1,1,1}(q)$ . This calculation makes use of various reduction identities for multiple sums that are determined in appendix C.

**1.2. A note on acronyms.** This subject suffers from a plethora of acronyms, which we will use in places to avoid repeating long names. These are summarised as follows.

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<sup>1</sup>The elliptic multiple polylogarithms considered in this paper are not holomorphic and differ from those introduced in [13], which are of relevance to the open string annulus amplitude, as discussed in [14, 15]. Although there is a clear relationship between the open string and the closed string, this relationship is not a subject studied in this paper.

MPL	Multiple polylogarithm.
svMPL	Single-valued multiple polylogarithm.
eMPL	Elliptic multiple polylogarithm.
sveMPL	Single-valued elliptic multiple polylogarithm.
MZV	Multiple zeta-value.
svMZV	Single-valued multiple zeta-value.

## 2. SOME BASIC FEATURES OF MULTIPLE POLYLOGARITHMS

The emphasis in this paper is the analogy between elliptic functions that arise as coefficients in the low energy expansion in genus-one closed superstring amplitudes with single-valued multi-zeta values (svMZVs) that arise in the expansion of the tree-level amplitudes [4, 5]. For clarity we will here give a brief (and incomplete) review of the relationship of these quantities to multiple polylogarithms and to single-valued multiple polylogarithms. There are many detailed reviews of this large subject in the literature and we note in particular the elementary introduction in [18].

**2.1. Polylogarithms.** The polylogarithm  $\text{Li}_a(z)$  is defined for any value of  $a \in \mathbb{C}$  by the power series expansion

$$\text{Li}_a(z) = \sum_{k=1}^{\infty} \frac{z^k}{k^a}, \quad (2.1)$$

which is absolutely convergent for  $|z| < 1$ . The polylogarithm is a natural generalisation of the logarithm since we have  $\text{Li}_1(z) = -\log(1-z)$ . Alternatively, the function  $\text{Li}_a(z)$  may be defined by the integral representation

$$\text{Li}_a(z) = \frac{z}{\Gamma(a)} \int_1^{\infty} \frac{dt}{t} \frac{(\log t)^{a-1}}{t-z}, \quad (2.2)$$

which coincides with (2.1) for  $|z| < 1$ , but may be analytically continued to all  $z \in \mathbb{C} \setminus [1, +\infty[$ . The resulting  $\text{Li}_a(z)$  manifestly has a branch points at  $z = 1$ , but it also has a branch point at  $z = 0$  on its higher Riemann sheets. Therefore, the function  $\text{Li}_a(z)$  for generic value  $a \in \mathbb{C}$  is multiple-valued and has interesting monodromies, again generalising the properties of  $\text{Li}_1(z) = -\ln(1-z)$ . Of particular interest is the relation of  $\text{Li}_a(z)$  to the Riemann zeta function  $\zeta(a)$  via

$$\text{Li}_a(1) = \zeta(a), \quad (2.3)$$

and to the Bernoulli polynomials  $B_n(x)$  via

$$\text{Li}_n(e^{2i\pi x}) + (-1)^n \text{Li}_n(e^{-2i\pi x}) = -\frac{(2i\pi)^n}{n!} B_n(x), \quad (2.4)$$

for  $n \in \mathbb{N}$  and  $x \in [0, 1]$ . Recall that  $B_n(x)$  is the  $n$ -th Bernoulli polynomial which is defined by the expansion in powers of  $t$  of the following generating function

$$\frac{t e^{xt}}{e^t - 1} = \sum_{n=0}^{\infty} B_n(x) \frac{t^n}{n!}. \quad (2.5)$$

Moreover, the function  $\text{Li}_a(z)$  satisfies the differential relation

$$z \frac{\partial \text{Li}_a(z)}{\partial z} = \text{Li}_{a-1}(z), \quad (2.6)$$

and the converse integral relation

$$\text{Li}_{a+1}(z) = \int_0^z \frac{dy}{y} \text{Li}_a(y). \quad (2.7)$$

**2.2. Single-valued polylogarithms.** At the cost of giving up holomorphicity in  $z$ , it is possible to construct a single-valued polylogarithm associated with  $\text{Li}_a(z)$ . The Bloch–Wigner dilogarithm is the prototype for such single-valued polylogarithms, and is defined as follows

$$D(z) = \Im \left( \text{Li}_2(z) + \log(1-z) \log|z| \right), \quad z \in \mathbb{C} \setminus \{0, 1\}. \quad (2.8)$$

In this paper we will make use of the generalisations of the Bloch–Wigner dilogarithm introduced by Zagier in [19]. They are single-valued (real analytic) polylogarithms  $D_{a,b}(z)$  for  $z \in \mathbb{C} \setminus [1, \infty[$  and  $a, b \in \mathbb{N}$ , and can be expressed in terms of sums of ordinary polylogarithms  $\text{Li}_k(z)$  for integer index  $k$  by

$$\begin{aligned} D_{a,b}(z) = & (-1)^{a-1} \sum_{k=a}^{a+b-1} \binom{k-1}{a-1} \frac{(-2 \log|z|)^{a+b-1-k}}{(a+b-1-k)!} \text{Li}_k(z) \\ & + (-1)^{b-1} \sum_{k=b}^{a+b-1} \binom{k-1}{b-1} \frac{(-2 \log|z|)^{a+b-1-k}}{(a+b-1-k)!} (\text{Li}_k(z))^*. \end{aligned} \quad (2.9)$$

These functions satisfy the complex conjugation relation  $D_{a,b}(z)^* = D_{b,a}(z)$ , and we define their *weight* to be  $a + b - 1$ . In particular, the function  $D_{a,a}(z)$  with  $a \in \mathbb{N}$  has weight  $2a - 1$  and is real and single-valued on  $\mathbb{C} \setminus \{0, 1\}$ . It is given by a finite linear combination of the  $\text{Li}_k(z)$ -functions

$$D_{a,a}(z) = 2\Re \left( (-1)^{a-1} \sum_{k=0}^{a-1} \binom{k+a-1}{a-1} \frac{(-2 \log|z|)^{a-1-k}}{(a-1-k)!} \text{Li}_{a+k}(z) \right), \quad (2.10)$$

and will play an important role in the sequel of this paper. The following are examples at low weights. At weight one, we have,

$$D_{1,1}(z) = -2 \log|1-z|^2, \quad (2.11)$$

which is manifestly single-valued. At weight 2 we have

$$\begin{aligned} D_{1,2}(z) &= 2 \log |z| \log(1-z) + 2i \Im(\text{Li}_2(z)), \\ D_{2,1}(z) &= D_{1,2}(z) - 4iD(z), \end{aligned} \quad (2.12)$$

where  $D(z)$  is the Bloch-Wigner dilogarithm of (2.8) which is single-valued on  $\mathbb{C} \setminus \{0, 1\}$ . At weight 3 we have

$$\begin{aligned} D_{1,3}(z) &= 2(\log |z|)^2 \text{Li}_1(z) - 2 \log |z| \text{Li}_2(z) + 2\Re(\text{Li}_3(z)), \\ D_{2,2}(z) &= 4 \log |z| \Re(\text{Li}_2(z)) - 4\Re(\text{Li}_3(z)). \end{aligned} \quad (2.13)$$

It is not difficult to check that these functions are algebraically independent.

**2.3. Multiple polylogarithms.** The multi-variable polylogarithm  $\text{Li}_{a_1, \dots, a_r}(z_1, \dots, z_r)$  is referred to as a multiple polylogarithm and was defined by [20, 21]

$$\text{Li}_{a_1, \dots, a_r}(z_1, \dots, z_r) = \sum_{1 < m_1 < \dots < m_r} \prod_{i=1}^r \frac{z_i^{m_i}}{m_i^{a_i}}, \quad (2.14)$$

with  $a_i \in \mathbb{N}$  for  $1 \leq i \leq r$ , and with  $|z_i| \leq 1$  ( $1 \leq i \leq r-1$ ) and  $|z_r| < 1$ . When  $a_i \in \mathbb{N} \geq 2$  for all  $1 \leq i \leq r$  this function is defined for  $|z_i| \leq 1$  for  $1 \leq i \leq r$ .

The *weight* of a multiple polylogarithm is the sum of its indices  $\sum_{i=1}^r a_i$  and its *depth* is the number of indices,  $r$ . The classical polylogarithm functions  $\text{Li}_a(z)$  are special cases with depth  $r = 1$ .

The multiple polylogarithms in (2.14) satisfy two kinds of differential relations. One of these reduces the depth by one but leaves the weight unchanged

$$z_i \frac{\partial}{\partial z_i} \text{Li}_{a_1, \dots, a_r}(z_1, \dots, z_r) = \text{Li}_{a_1, \dots, a_{i-1}, \dots, a_r}(z_1, \dots, z_r) \quad 1 \leq i \leq r, \quad (2.15)$$

while the other reduces both the weight and the depth by one

$$(1 - z_r) \frac{\partial}{\partial z_r} \text{Li}_{a_1, \dots, a_{r-1}, 1}(z_1, \dots, z_r) = \text{Li}_{a_1, \dots, a_{r-1}}(z_1, \dots, z_{r-1} z_r). \quad (2.16)$$

Together with the initial conditions  $\text{Li}_{a_1, \dots, a_r}(0, \dots, 0) = 0$  these differential equations determine the multiple polylogarithms by multiple integration. This leads to Chen iterated integrals [22] which endow the space of multiple polylogarithms with a *shuffle algebra*. Since we will not make use of this construction we refer to [18] for details and references. For additional reference, we note that these properties have been implemented in the algebraic program `Hyperint` by Erik Panzer [23] and in `MPL` by Christian Bogner in [24].

From the series representation one derives the *stuffle relations* as described in [25], and implemented in [26, 27]. The stuffle relation of two multiple polylogarithms is given

by

$$\text{Li}_{a_1, \dots, a_r}(x_1, \dots, x_r) \text{Li}_{b_1, \dots, b_s}(y_1, \dots, y_s) = \sum_{k=\max(r,s)}^{r+s} \text{Li}_{c_1, \dots, c_k}(z_1, \dots, z_k), \quad (2.17)$$

where the set  $c(k) = (c_1, \dots, c_k)$  is obtained by adding zeroes in all possible ways to the set  $(a_1, \dots, a_r)$  and  $(b_1, \dots, b_s)$  to get length  $k \leq r + s$  sets denoted  $a(k)$  and  $b(k)$ , and then adding the two sets  $c(k) = a(k) + b(k)$ . The argument  $z_i$  associated with the index  $c(k)_i = a(k)_i + b(k)_i$  is obtained as follows:

- if  $a(k)_i \neq 0$  and  $b(k)_i = 0$  then  $z_i = x_i$ ;
- if  $a(k)_i = 0$  and  $b(k)_i \neq 0$  then  $z_i = y_i$ ;
- if  $a(k)_i \neq 0$  and  $b(k)_i \neq 0$  then  $z_i = x_i y_i$ .

Note that since we have assumed that  $a_i \geq 1$  for all  $i = 1, \dots, r$  and  $b_j \geq 1$  for all  $j = 1, \dots, s$ , the case where both  $a(k)_i = 0$  and  $b(k)_i = 0$  cannot arise.

We illustrate this with two examples that will be useful later. The first example is the product of two depth-1 polylogarithms

$$\text{Li}_a(x) \text{Li}_b(y) = \text{Li}_{a+b}(xy) + \text{Li}_{a,b}(x, y) + \text{Li}_{b,a}(y, x). \quad (2.18)$$

The second example is the product of a depth-2 multiple polylogarithm by a depth-1 polylogarithm

$$\begin{aligned} \text{Li}_{a,b}(x, y) \text{Li}_c(z) &= \text{Li}_{a+c,b}(xz, y) + \text{Li}_{a,b+c}(x, yz) \\ &\quad + \text{Li}_{a,b,c}(x, y, z) + \text{Li}_{a,c,b}(x, z, y) + \text{Li}_{c,a,b}(z, x, y). \end{aligned} \quad (2.19)$$

**2.4. Single-valued multiple polylogarithms.** The special case of single-variable multiple polylogarithms, which have the form

$$\text{Li}_{a_1, \dots, a_r}(z) = \text{Li}_{a_1, \dots, a_r}(1, \dots, 1, z), \quad (2.20)$$

have monodromies around  $z = 0$  and  $z = 1$ . However there is an important subclass of such functions that are single valued on  $\mathbb{C} \setminus \{0, 1\}$  and were constructed by Francis Brown in [6, 8]. They are obtained by appropriate linear combinations of single-variable multiple polylogarithms of the same weight and define functions on  $\mathbb{C} \setminus \{0, 1\}$  without monodromies. These combinations are examples of *single-valued multiple polylogarithms* (svMPLs). Generalising this discussion to define single-valued multiple polylogarithms with more than one variable argument appears to be very complicated and will not be needed for our purposes.

**2.5. Multiple zeta-values and single-valued multiple zeta-values.** Multiple zeta values (MZVs) [28] provide a natural generalization of Riemann zeta values. The multiple zeta function may be defined by the following multiple sums for  $\Re(a_i) > 1$ ,

$$\zeta(a_1, \dots, a_r) = \sum_{0 < n_1 < \dots < n_r} \prod_{i=1}^r \frac{1}{n_i^{a_i}}, \quad (2.21)$$

and analytically continued to  $a_i \in \mathbb{C}$ . For  $a_i \in \mathbb{N}$ , and  $a_r > 1$ , the quantities  $\zeta(a_1, \dots, a_r)$  give *multiple zeta values*. The multiple polylogarithm evaluated with all its arguments equal to 1 is a multiple zeta value given by

$$\text{Li}_{a_1, \dots, a_r}(1, \dots, 1) = \zeta(a_1, \dots, a_r), \quad (2.22)$$

when  $a_i \in \mathbb{N}$  with  $a_r > 1$ .<sup>2</sup>

The basis of MZVs has been understood in terms of periods on the Riemann sphere with marked points [29]. This interpretation also arises naturally in the analysis of the low energy expansion of string theory tree amplitudes [30, 31]. The stuffle and shuffle relations for multiple polylogarithms lead to a wealth of algebraic relations between MZVs.

The single-valued multiple zeta values (svMZVs) are obtained by setting the arguments of single-valued multiple polylogarithms to 1. For example,

$$\begin{aligned} D_{1,2}(1) &= D_{2,1}(1) = 0, \\ D_{1,3}(1) &= 2\zeta(3) \\ D_{2,2}(1) &= -4\zeta(3). \end{aligned} \quad (2.23)$$

It is not difficult to check that these functions have the properties

$$\begin{aligned} D_{a,b}(1) &= 0 & a+b \in 2\mathbb{N}-1 \\ D_{a,b}(1) &\in \zeta(a+b-1) \times \mathbb{Z} & a+b \in 2\mathbb{N}. \end{aligned} \quad (2.24)$$

This has led Brown to define *single-valued zeta values*  $\zeta_{sv}$  that are given by [8]

$$\begin{aligned} \zeta_{sv}(2n) &= 0 & n \in \mathbb{N} \\ \zeta_{sv}(2n+1) &= 2\zeta(2n+1). \end{aligned} \quad (2.25)$$

It is important to appreciate that the definition of a single-valued polylogarithm on  $\mathbb{C} \setminus \{0, 1\}$  is not unique for weight greater than two.<sup>3</sup> However, when evaluated at special values of its arguments, any definition of a single-valued multiple polylogarithm (svMPL) gives rise to the same, unique, svMZV. The functions  $D_{a,b}(z)$  in (2.9) give a particular definition of svMPLs.

**2.6. The basis of svMZVs.** The dimension  $d_w$  of the subspace of multiple zeta values of weight  $w$  in the ring over  $\mathbb{Q}$  of multiple zeta values is given by the expansion of a rational function [28, 32, 33]

$$\sum_{w=0}^{\infty} d_w x^w = \frac{1}{1-x^2-x^3}. \quad (2.26)$$

<sup>2</sup>When  $a_r = 1$  the expression diverges and has to be regularised, but this case will not concern us in the sequel.

<sup>3</sup>A discussion of the general properties of such functions and the relations between various equivalent definitions may be found, for example, in [6].



The dimension  $d_w^{sv}$  of the subspace of weight  $w$  in the ring over  $\mathbb{Q}$  of single-valued multiple-zeta values is smaller than  $d_w$  [8]. For example, at weight 11 the basis of ordinary MZVs has dimension 9 and is composed of

$$\begin{aligned} &\zeta(3, 5, 3), \zeta(3, 5)\zeta(3), \zeta(3)^2\zeta(5), \zeta(11), \\ &\zeta(2)\zeta(3)^3, \zeta(2)^4\zeta(3), \zeta(2)^3\zeta(5), \zeta(2)^2\zeta(7), \zeta(2)\zeta(9). \end{aligned} \quad (2.27)$$

Since by definition we have  $\zeta_{sv}(2) = 0$  all the basis elements for MZVs on the second line vanish are mapped to zero as elements of svMZVs. At weight 11 a further reduction happens since according to [8] we have the relation

$$\zeta_{sv}(3, 5, 3) = 2\zeta(3, 5, 3) - 2\zeta(3)\zeta(3, 5) - 10\zeta(3)^2\zeta(5), \quad (2.28)$$

and since  $\zeta_{sv}(3, 5) = -10\zeta(3)\zeta(5)$ , the basis of svMZVs at weight 11 has dimension 3 and is composed of

$$\zeta_{sv}(3, 5, 3), \zeta_{sv}(3)^2\zeta_{sv}(5), \zeta_{sv}(11). \quad (2.29)$$

Corollary 7.4 of [8] gives the dimensions  $d_w^{sv}$  of the subspace of svMZVs of weight  $w$

$$\sum_{w=0}^{\infty} d_w^{sv} x^w = \prod_{n=0}^{\infty} (1 - x^{2n+1})^{-\ell_{2n+1}}, \quad (2.30)$$

where the positive integers  $\ell_n$  are defined by the following product relation,

$$\prod_{n=1}^{\infty} (1 - x^n)^{\ell_n} = 1 - x^2 - x^3. \quad (2.31)$$

Note that, up to weight 12, the non-zero values for  $\ell$  are  $\ell_2 = \ell_3 = \ell_5 = \ell_7 = \ell_8 = \ell_9 = \ell_{10} = 1$  and  $\ell_{11} = \ell_{12} = 2$ , so that the non-zero values for  $d_w^{sv}$  are  $d_3^{sv} = d_5^{sv} = d_6^{sv} = d_7^{sv} = d_8^{sv} = 1$ , while  $d_9^{sv} = d_{10}^{sv} = 2$  and  $d_{11}^{sv} = 3$  in agreement with our earlier calculation. Single-valued MZVs are very efficiently obtained using Schnetz' Maple routines [34].

We will give an example in section 4.3 of how such svMZVs enter into the Laurent expansions of the modular graph functions. They have already appeared in Feynman diagrams in quantum field theory [7] and in the low energy expansion of closed string three level amplitudes [4, 5].

**2.7. Single-valued elliptic polylogarithms.** In the main part of this paper we will consider generalisations of polylogarithms and multiple polylogarithms that are defined on a general elliptic curve, that enters, for example, into the integrand of the genus one closed string amplitude. An elliptic curve  $\mathcal{E}$  can be described either as  $\mathcal{E} \simeq \mathbb{C}/(\mathbb{Z} + \tau\mathbb{Z})$  where  $\tau$  is the period ratio or as  $\mathcal{E} \simeq \mathbb{C}^\times/q^{\mathbb{Z}}$  where  $q = \exp(2i\pi\tau)$ , and  $\mathbb{C}^\times$  is the multiplicative group of non-zero complex numbers. A point  $P$  on the elliptic curve is represented as  $\zeta(P) = \exp(2i\pi(u\tau + v)) \in \mathbb{C}^\times/q^{\mathbb{Z}}$  where  $u$  and  $v$  are cartesian coordinate in  $[0, 1]$ .

We will make extensive use of the elliptic polylogarithms that were constructed by Zagier [19]. These can be expressed as infinite sums over positive integers  $n$  of images of the polylogarithms  $\text{Li}_k(z)$  under the translations  $z \rightarrow z + n\tau$ . To be explicit, these functions are defined by

$$D_{a,b}(q; \zeta) = \sum_{n \geq 0} D_{a,b}(q^n \zeta) + (-1)^{a+b} \sum_{n \geq 1} D_{a,b}(q^n / \zeta) + \frac{(4\pi\tau_2)^{a+b-1}}{(a+b)!} B_{a+b}(u), \quad (2.32)$$

where the function  $D_{a,b}(x)$  was defined in (2.9),  $B_n$  is the  $n$ -th Bernoulli polynomial, and  $\zeta = e^{2\pi i(u\tau+v)}$  corresponds to a point on the torus with modulus  $q = e^{2\pi i\tau}$  for  $u, v \in [0, 1]$ . Zagier showed that if  $a$  and  $b$  are two positive integers, this function is equal to the double sum

$$D_{a,b}(q; \zeta) = \frac{(2i\tau_2)^{a+b-1}}{2i\pi} \sum_{(m,n) \neq (0,0)} \frac{e^{2i\pi(nu-mv)}}{(m\tau+n)^a (m\bar{\tau}+n)^b}. \quad (2.33)$$

The function  $D_{a,b}(q; \zeta)$  is invariant under  $\zeta \rightarrow q\zeta$  since the expression in (2.33) is invariant under the shifts  $u \rightarrow u+1$  and  $v \rightarrow v+1$ , therefore this is a function on the elliptic curve. Furthermore, under the action of  $SL(2, \mathbb{Z})$

$$\begin{aligned} \tau &\rightarrow \frac{\alpha\tau + \beta}{\gamma\tau + \delta} & \alpha\delta - \beta\gamma &= 1, & \alpha, \beta, \gamma, \delta &\in \mathbb{Z} \\ (u, v) &\rightarrow (\delta u - \gamma v, -\beta u + \alpha v) \end{aligned} \quad (2.34)$$

$D_{a,b}(q; \zeta)$  transforms as a modular form of weight  $(1-b, 1-a)$ , namely,

$$D_{a,b}(q; \zeta) \rightarrow (\gamma\tau + \delta)^{b-1} (\gamma\bar{\tau} + \delta)^{a-1} D_{a,b}(q; \zeta). \quad (2.35)$$

The special case  $a = b$  will be of particular importance in this paper and is given by

$$D_{a,a}(q; \zeta) = \sum_{n \geq 0} D_{a,a}(q^n \zeta) + \sum_{n \geq 1} D_{a,a}(q^n / \zeta) + \frac{(4\pi\tau_2)^{2a-1}}{(2a)!} B_{2a}(u), \quad (2.36)$$

where  $D_{a,a}(x)$  is given by (2.10). The elliptic polylogarithms defined in this manner are manifestly single-valued, as follows from their form in (2.33).

It will be established later in this paper that we may define *single-valued elliptic multiple polylogarithms* (sveMPLs) as integrals of products of  $D_{1,1}(q; \zeta)$  factors over the torus, with one argument left unintegrated. Modular graph functions may then be described as the values of sveMPLs with the unintegrated argument set equal to 1. The terminology *single-valued* arises because sveMPLs have no monodromy as functions of their unintegrated arguments. This mimics the terminology used for single-valued MZVs, which are the values of single-valued multiple polylogarithms with arguments set equal to 1.

These functions should be distinguished from the elliptic multiple polylogarithms that arise in [13], which are holomorphic since they are only functions of  $q$  or  $\bar{q}$ . These

are of relevance to loop amplitudes in open string theory [14] as well as arising in dimensionally regularised Feynman graphs in quantum field theory at two loops [35–37] and three loops [38]. References [35, 38] give a motivic approach to these objects.

### 3. FEYNMAN GRAPHS ASSOCIATED WITH A TORUS

Our motivation for considering modular graph functions originated from the study of the low energy expansion of the genus-one four-graviton amplitude in Type II superstring theory. The graphs generated in this expansion have up to four vertices joined by the scalar Green function on the torus, as will be reviewed below. However, the structure of such graphs generalises in an obvious fashion to graphs with an arbitrary number of vertices, and it is these general graphs that will be considered in this section and sections 4, 5, and 6 below. Actually, this obvious generalisation will not quite encompass all the contributions that can arise in the  $N$ -string amplitudes at genus-one for  $N > 4$ , as was discussed to a limited extent in [39]. We will comment on the structure of such graphs in section 7.

The four-graviton amplitude at genus one is given by an integral over the moduli space  $\mathcal{M}_1$  of a torus  $\Sigma$  of a partial amplitude  $\mathcal{B}_4$  evaluated at fixed modulus.<sup>4</sup> We define the following partial amplitude  $\mathcal{B}_N(s_{ij}|\tau)$  for arbitrary  $N$

$$\mathcal{B}_N(s_{ij}|\tau) = \prod_{n=1}^N \int_{\Sigma} \frac{d^2 z_n}{\tau_2} \exp \left( \sum_{1 \leq i < j \leq N} s_{ij} G(z_i - z_j|\tau) \right). \quad (3.1)$$

The scalar Green function  $G(z_i - z_j|\tau)$  on the torus will be discussed in the subsection below. In superstring theory, the parameters  $s_{ij}$  are related to the momenta  $k_i$  of  $N$  gravitons, for  $i = 1, \dots, N$ . Each momentum  $k_i$  is a null-vector in 10-dimensional Minkowski space-time  $\mathbb{R}^{10}$  so that  $k_i \cdot k_i = 0$ . The relation between the  $s_{ij}$  and  $k_i$  is given by  $s_{ij} = -\ell_s^2 k_i \cdot k_j / 2$ , where  $\ell_s$  is the string length. The null-vector condition on  $k_i$  and overall momentum conservation  $\sum_{i=1}^N k_i = 0$  impose the following relations between the parameters  $s_{ij}$

$$s_{ii} = 0, \quad \sum_{j=1}^N s_{ij} = 0, \quad i = 1, \dots, N. \quad (3.2)$$

Actually, any partition of the sum over momenta into two disjoint sets will give rise to extra linear *partition relations* between the variables  $s_{ij}$ . As a result, for  $N \leq 3$ ,

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<sup>4</sup>A torus  $\Sigma$  with modulus  $\tau$  may be represented in the complex plane by  $\mathbb{C}/(\mathbb{Z}\tau + \mathbb{Z})$  and parametrized by a complex coordinate  $z = u\tau + v$  or equivalently by two real coordinates  $u, v \in \mathbb{R}/\mathbb{Z}$ . The volume form on  $\Sigma$  is normalized to  $d^2 z = idz \wedge d\bar{z}/2 = -\tau_2 du \wedge dv$ . The moduli space  $\mathcal{M}_1$  of orientable genus-one Riemann surfaces (or tori) may be represented by a fundamental domain for the action of  $PSL(2, \mathbb{Z})$  on the Poincaré upper half plane, which may be parametrized by  $\tau = \tau_1 + i\tau_2$  with  $\tau_1, \tau_2 \in \mathbb{R}$  and  $\mathcal{M}_1 = \{\tau, 0 < \tau_2, |\tau| \geq 1, |\tau_1| < \frac{1}{2}\}$ .

we have  $s_{ij} = 0$  for all  $1 \leq i, j \leq 3$  and thus  $\mathcal{B}_1 = \mathcal{B}_2 = \mathcal{B}_3 = 1$ . The first non-trivial case is the four-graviton amplitude with  $N = 4$ , and extra *partition relations*  $s_{12} = s_{34}$ ,  $s_{23} = s_{14}$ ,  $s_{13} = s_{24}$  resulting from partitioning the momenta into pairs.

The partial amplitude  $\mathcal{B}_N(s_{ij}|\tau)$ , with the parameters  $s_{ij}$  subject to the relations (3.2), is well-defined for any value of  $N$ , and will be taken here as the generating function for modular graph functions, although it is only when  $N = 4$  that  $\mathcal{B}_N(s_{ij}|\tau)$  corresponds precisely to the partial closed superstring amplitude.<sup>5</sup> The restrictions imposed by the extra *partition relations* between the parameters  $s_{ij}$  will be left unenforced, as they may be imposed easily, if needed, when constructing the partial amplitudes for actual string processes.

**3.1. The scalar Green function.** A scalar Green function  $G(z|\tau)$  is defined to be the inverse of the Laplace operator  $\Delta = 4\partial_z\partial_{\bar{z}}$  on the torus,  $\Sigma$ . Since  $\mathcal{B}_N$  is invariant under a  $z$ -independent shift of the Green function  $G$ , we may restrict the Green function to have vanishing integral on  $\Sigma$ . Therefore,  $G(z|\tau)$  obeys the following equations

$$\Delta G(z|\tau) = -4\pi\delta^{(2)}(z) + \frac{4\pi}{\tau_2} \quad \int_{\Sigma} d^2z G(z|\tau) = 0. \quad (3.3)$$

This determines the modular invariant formula for  $G$  in terms of the Jacobi  $\vartheta_1$ -function and the Dedekind  $\eta$ -function

$$G(z|\tau) = -\ln \left| \frac{\vartheta_1(z|\tau)}{\eta(\tau)} \right|^2 - \frac{\pi}{2\tau_2} (z - \bar{z})^2. \quad (3.4)$$

Equivalently, the Green function may be expressed as a Fourier sum in terms of the real variables  $u, v \in \mathbb{R}/\mathbb{Z}$  defined by  $z = u\tau + v$

$$G(z|\tau) = \sum_{(m,n) \neq (0,0)} \frac{\tau_2}{\pi|m\tau + n|^2} e^{2i\pi(mv - nu)}. \quad (3.5)$$

The integers  $m, n$  parametrise the discrete momenta  $p = m\tau + n \in \mathbb{Z}\tau + \mathbb{Z}$  of the Fourier modes on the torus.

To make contact with polylogarithms, we change variables for both the modulus  $\tau$  and the coordinate  $z$  on the torus, to new variables defined in terms of  $\tau$  and  $z$  by

$$q = e^{2\pi i\tau} \quad \zeta = e^{2\pi iz} = e^{2\pi i(u\tau + v)} \quad (3.6)$$

The condition  $0 < \tau_2$  guarantees that  $|q| < 1$ , while periodicity of  $\zeta$  in  $v$  shows that  $\Sigma$  is parametrized by  $\zeta \in \mathbb{C}^\times/q^{\mathbb{Z}}$ . Comparing the expression for  $D_{a,b}(q; \zeta)$  of (2.33) for  $a = b = 1$  with the Fourier sum in (3.5), we are led to the following identification [16, 19]

$$D_{1,1}(q; \zeta) = G(z|\tau). \quad (3.7)$$

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<sup>5</sup>See the discussion in section 7 for some comments concerning the  $N > 4$  amplitude.

Similarly, the expression  $D_{1,1}(q; \zeta_i/\zeta_j)$  will equal the Green function  $G(z_i - z_j|\tau)$  on the torus between the points  $i$  and  $j$ , with  $\zeta_i/\zeta_j = e^{2i\pi(z_i - z_j)}$ . In view of this identity, we shall often refer to  $D_{1,1}(q; \zeta)$  as the Green function as well.

Importantly for the subsequent discussions, the Green function  $D_{1,1}(q; \zeta)$  admits a simple expression in terms of polylogarithms  $2\Re\text{Li}_1(\cdot)$ , which are single-valued on  $\mathbb{C} \setminus \{0, 1\}$ , evaluated on an elliptic curve

$$D_{1,1}(q; \zeta) = 2\Re \left( \sum_{n \geq 0} \text{Li}_1(q^n \zeta) + \sum_{n \geq 1} \text{Li}_1(q^n / \zeta) + \pi \tau_2 B_2(u) \right), \quad (3.8)$$

where  $\zeta = q^u e^{2i\pi v}$  with  $u, v \in [0, 1]$ , and  $B_2(u)$  is the second Bernoulli polynomial.

**3.2. Modular graph functions.** The Fourier representation of  $G(z|\tau)$  makes it transparent that the Green function  $G(z|\tau)$  is invariant under the modular transformations of (2.34). As a result, the function  $\mathcal{B}_N(s_{ij}|\tau)$  is also modular invariant when the variables  $s_{ij}$  are left unchanged under the modular transformation. The series expansion of  $\mathcal{B}_N(s_{ij}|\tau)$  in powers of  $s_{ij}$  has a finite radius of convergence, and may be organized graphically in what physicists refer to as Feynman graphs, in this case for a conformal scalar field on the torus. Each term in this expansion corresponds to a specific graph, and will evaluate to a specific non-holomorphic modular function, whence the terminology *modular graph function*. Thus, we may view  $\mathcal{B}_N(s_{ij}|\tau)$  as a generating function for all such modular graph functions.

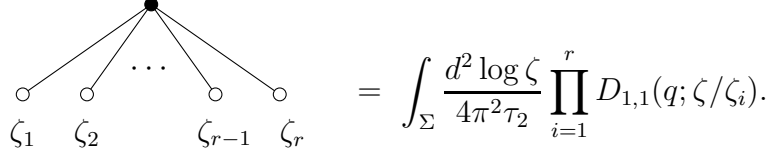
The building block of the modular graphs is the scalar Green function  $G(z_i - z_j|\tau)$ , which we associate with an edge in the graph between vertices  $z_i$  and  $z_j$ . Alternatively, and more conveniently for application to polylogarithm, the scalar Green function may be set to the function  $D_{1,1}(q; \zeta_i/\zeta_j)$ . We have the following graphical representation

$$\begin{array}{ccc} \zeta_i & & \zeta_j \\ \circ & \text{-----} & \circ \\ & & = D_{1,1}(q; \zeta_i/\zeta_j) = G(z_i - z_j|\tau). \end{array}$$

The Green function between two vertices  $\zeta_i$  and  $\zeta_j$  may enter into a given graph raised to an integer power  $n$ , for which we shall use the following graphical representation,

$$\begin{array}{ccc} \zeta_i & & \zeta_j \\ \circ & \text{-----} \textcircled{n} \text{-----} & \circ \\ & & = D_{1,1}(q; \zeta_i/\zeta_j)^n = G(z_i - z_j|\tau)^n. \end{array}$$

Finally, the integration at one vertex over a product of Green functions ending at that vertex will be denoted by a filled black dot, in contrast with an unintegrated vertex which will be represented by an unfilled dot. The basic ingredient in the graphical notation is depicted in the graph below



$$= \int_{\Sigma} \frac{d^2 \log \zeta}{4\pi^2 \tau_2} \prod_{i=1}^r D_{1,1}(q; \zeta/\zeta_i).$$

The low energy expansion of  $\mathcal{B}_N$  is obtained simply by expanding the exponential in (3.1) in powers of  $s_{ij}$ , or equivalently by expanding in powers of  $\ell_s^2$ , while keeping the momenta  $k_i$  fixed. The coefficient of the term which is homogeneous of degree  $w$  in the  $s_{ij}$  variables is given by a sum over all possible graphs with  $w$  Green functions joining at most  $N$  vertices, as studied for the special case of  $N = 4$  in [9, 16, 17]. Each weight  $w$  graph  $\Gamma_{w,N}$  is evaluated by the integration of the product of Green functions over the position of the vertices  $\zeta_i$  on the torus  $\Sigma$ . When no confusion is expected to arise, we shall abbreviate  $\Gamma_{w,N}$  by simply  $\Gamma$ . The resulting function will be denoted  $I_{\Gamma}(q)$  and is given by

$$I_{\Gamma}(q) = \prod_{k=1}^N \int_{\Sigma} \frac{d^2 \log \zeta_k}{4\pi^2 \tau_2} \prod_{1 \leq i < j \leq N} D_{1,1}(q; \zeta_i/\zeta_j)^{n_{ij}}, \quad (3.9)$$

where  $N$  is the number of vertices and  $n_{ij}$  is the number of Green functions joining the vertices  $i$  and  $j$ . The numbers  $n_{ij}$  are the entries of the adjacency matrix of the graph  $\Gamma = \Gamma_{w,N}$ . By construction, the integral  $I_{\Gamma}(q)$  is modular invariant, as it arises from the expansion in powers of  $s_{ij}$  of the modular invariant generating function  $\mathcal{B}_N(s_{ij}|\tau)$ , and therefore associates with a graph  $\Gamma$  a non-holomorphic modular function  $I_{\Gamma}(q)$ .

An alternative evaluation of the modular graph function  $I_{\Gamma}(q)$  is obtained by using the Fourier representation

$$D_{1,1}(q; \zeta) = \sum_{(m,n) \neq (0,0)} \frac{\tau_2}{\pi |m\tau + n|^2} e^{2i\pi(mv - nu)}, \quad (3.10)$$

where  $\zeta = e^{2\pi i(u\tau + v)}$ . The integration over  $\Sigma$  of the position  $\zeta_i$  of the vertex  $i$  enforces momentum conservation on all the Green functions that enter the vertex  $i$ . Carrying out the integrations over  $\Sigma$  for all the vertex positions  $\zeta_i$  for  $i = 1, \dots, N$  gives the constrained multiple sum representation for the graphs, as was studied already earlier for the case  $N = 4$  in [9, 10, 16, 17]. The general form of the sum may be schematically represented as follows

$$I_{\Gamma}(q) = \sum'_{p_1, \dots, p_w \in \mathbb{Z}\tau + \mathbb{Z}} \prod_{\alpha=1}^w \frac{\tau_2}{\pi |p_{\alpha}|^2} \prod_{i=1}^N \delta \left( \sum_{\alpha=1}^w C_{i\alpha} p_{\alpha} \right). \quad (3.11)$$

Here, the prime above the summation symbol indicates that the sums over  $p$  exclude the value 0; the Kronecker  $\delta$  symbol takes the value 1 when its argument vanishes and

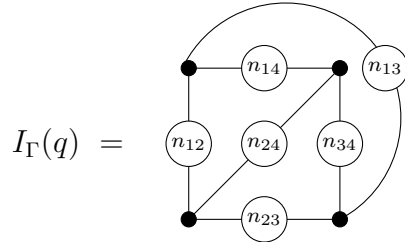
zero otherwise; the coefficients  $C_{i\alpha}$  are given as follows

$$C_{i\alpha} = \begin{cases} \pm 1 & \text{if edge } \alpha \text{ ends on vertex } i \\ 0 & \text{otherwise} \end{cases} \quad (3.12)$$

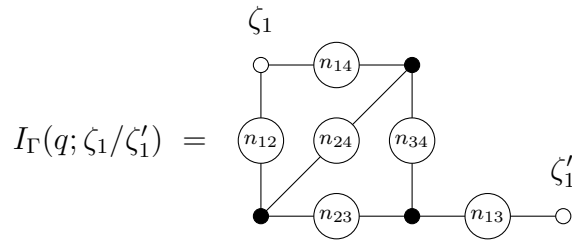
the sign being determined by the orientation of the momenta. Note that, given the weight  $w$ , all the information on the graph  $\Gamma$  is encoded in the coefficients  $C_{i\alpha}$ , the other parts of the sum in (3.11) being the same for all  $w$ . This momentum representation of the modular graph function clearly confirms its modular invariance.

**3.3. Single-valued nature of the modular graph functions.** In order to understand the statement that the modular graph functions are *single-valued*, we now show that they are naturally expressed in terms of elliptic analogues of single-valued multiple polylogarithms. We will call these functions *single-valued elliptic multiple polylogarithms*. We will demonstrate that the function  $I_\Gamma(q)$  is a special value of a more general function,  $I_\Gamma(q; \zeta)$ , associated with the graph  $\Gamma = \Gamma_{w,N}$  when all but one of the vertices, say vertex  $j$ , are integrated and the lines meeting the unintegrated vertex  $j$  are made to terminate at two separate points,  $\zeta_j$  and  $\zeta'_j$ . By translation invariance on the torus, we may set  $\zeta'_j = 1$  without loss of generality. The function  $I_\Gamma(q, \zeta)$  has no monodromies in  $\zeta$  and is therefore single-valued on the torus.

To be as concrete as possible, we shall concentrate here on the case of  $N = 4$ , but our considerations may be easily generalised to other values of  $N$ . An arbitrary graph  $\Gamma = \Gamma_{w,4}$  of weight  $w$  and with four vertices may be labeled as follows



The resulting elliptic function  $I_\Gamma(q; \zeta_1)$  is illustrated by the graph below with five vertices, of which two are unintegrated. The unintegrated vertex corresponds here to the vertex  $j = 1$  in the figure



The black dots represent integrated vertices, while the white dots continue to represent unintegrated vertices. To such a graph we associate an elliptic function that results from the integration over the remaining three vertices, which here are  $k = 2, 3, 4$

$$I_\Gamma(q; \zeta) = \prod_{k=2}^4 \int_\Sigma \frac{d^2 \log \zeta_k}{4\pi^2 \tau_2} \prod_{1 \leq i < j \leq 4} D_{1,1}(q; \zeta_j / \zeta_i)^{n_{ij}} \left( \frac{D_{1,1}(q; \zeta'_1 / \zeta_3)}{D_{1,1}(q; \zeta_1 / \zeta_3)} \right)^{n_{13}}. \quad (3.13)$$

The function  $I_\Gamma(q; \zeta)$  depends only on the ratio  $\zeta = \zeta_1 / \zeta'_1$  due to translation invariance on the torus. Since the Green function  $D_{1,1}(q; \zeta)$  is a single-valued function of  $\zeta$  the function  $I_\Gamma(q; \zeta)$  is a single-valued function of  $\zeta$ . Translation invariance implies that the full integral (3.9) is recovered by identifying the points  $\zeta_1$  and  $\zeta'_1$  in the above graph, in other words, by setting  $\zeta = 1$ . This is easily shown using the representation of (3.11) in which the discrete momentum is preserved at each vertex.

**3.4. Proposition.** *Any modular graph function  $I_\Gamma(q)$  may be obtained as the value of a single-valued elliptic multiple polylogarithm  $I_\Gamma(q; \zeta)$  evaluated at the point  $\zeta = 1$ .*

#### 4. EXAMPLES OF MODULAR GRAPH FUNCTIONS

In this section we shall illustrate the structure of modular graph functions  $I_\Gamma(q)$  and their associated single-valued elliptic multiple polylogarithms  $I_\Gamma(q; \zeta)$  on some simple but significant examples. It will be convenient to characterize the families of graphs under consideration by the number of loops  $L$  of the graph  $\Gamma$  in the corresponding modular graph function  $I_\Gamma(q)$ .

**4.1. General one-loop graphs.** The first family of examples is based on modular graph functions  $I_\Gamma(q)$  for graphs  $\Gamma$  with a single loop. The associated single-valued elliptic multiple polylogarithm results from a linear chain graph in which  $a$  Green functions are concatenated and integrated over their  $a - 1$  junction points. Such linear chain graphs are important in their own right, but will also provide natural building blocks for higher graphs that involve linear chains. We introduce the graphic notation given in the figure below

$$\zeta_1 \text{---} \boxed{a} \text{---} \zeta_{a+1} = \zeta_1 \text{---} \underbrace{\text{---} \bullet \text{---} \text{---} \bullet \text{---}}_a \text{---} \zeta_{a+1}$$

Recall that the black dots correspond to integrated vertices, while the white dots are unintegrated and the corresponding vertices are evaluated at the labels of these points. The weight of the graph is  $w = a$ , while its number of vertices is  $N = a + 1$ . The associated elliptic function may be evaluated by performing the  $a - 1$  integrations over the vertices  $\zeta_j$  for  $j = 2, \dots, a$ . A convenient way to carry out this evaluation is with



the help of the Fourier representation of (2.33). By comparing the result with Fourier representation given in (2.33) for  $D_{a,a}(q; \zeta)$ , we readily find

$$\prod_{k=2}^a \int_{\Sigma} \frac{d^2 \log \zeta_k}{4\pi^2 \tau_2} \prod_{j=1}^a D_{1,1}(q; \zeta_{j+1}/\zeta_j) = (-4\pi\tau_2)^{1-a} D_{a,a}(q; \zeta_{a+1}/\zeta_1), \quad (4.1)$$

which is the single-valued elliptic polylogarithm defined in [19]. The single-valuedness in  $\zeta = \zeta_{a+1}/\zeta_1$  of  $D_{a,a}(q; \zeta)$  is again a consequence of this integral representation, as explained in section 3. We may summarize this results by the graphical representation for the linear chain graph in the figure below

$$\circ - \boxed{a} - \circ = (-4\pi\tau_2)^{1-a} D_{a,a}(q; \zeta_{a+1}/\zeta_1) .$$

By evaluating the open linear chain graph when the vertices  $\zeta_1$  and  $\zeta_{a+1}$  coincide yields a one-loop chain graph, depicted in the figure below

$$\boxed{a} \text{ (loop) } = (-4\pi\tau_2)^{1-a} D_{a,a}(q; 1) .$$

Upon setting  $\zeta = 1$  in the argument of  $D_{a,a}(q; \zeta)$  in (2.33) we see that

$$(-4\pi\tau_2)^{1-a} D_{a,a}(q; 1) = E_a(q) = \sum_{(m,n) \neq (0,0)} \frac{\tau_2^a}{\pi^a |m\tau + n|^{2a}} . \quad (4.2)$$

The representation of the real analytic  $E_a$  Eisenstein series as a one-loop graph was used in [9, 16, 17]. The statement that the Eisenstein series is single-valued follows from the fact that it is a special value ( $\zeta = 1$ ) of the single-valued elliptic polylogarithm function  $D_{a,a}(q; \zeta)$ .

**4.2. General two-loop graphs.** The general two-loop graph that produces a modular graph function is given by the modular functions  $I_{\Gamma}(q) = C_{a,b,c}(q)$  introduced and studied extensively in [9].<sup>6</sup> These graphs always have two trivalent vertices, along with  $a + b + c - 3$  bivalent vertices. For any assignment of value of  $a, b, c \geq 1$ , a single-valued elliptic multiple polylogarithm  $I_{\Gamma}(q; \zeta) = C_{a,b,c}(q; \zeta)$  may be constructed from  $\Gamma$  by detaching one edge from one of the two trivalent vertices. By permutation symmetry of the indices  $a, b, c$ , we may detach the edge labeled  $a$  without loss of generality, to obtain the graph shown in the figure below

$$\zeta_1 \text{ (loop with } a, b, c \text{)} \zeta'_1 = C_{a,b,c}(q; \zeta_1/\zeta'_1) .$$

<sup>6</sup>This is a slight change of notation from [9] where we used  $C_{a,b,c}(\tau)$  to represent this function.

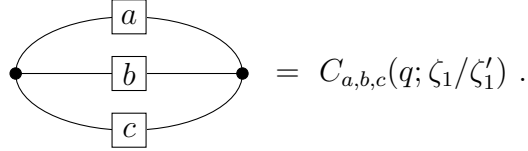
The associated single-valued elliptic multiple polylogarithm,  $C_{a,b,c}(q; \zeta)$  is a function of the ratio  $\zeta = \zeta_1/\zeta'_1$  and has the form

$$C_{a,b,c}(q; \zeta) = (-4\pi\tau_2)^{3-a-b-c} \int_{\Sigma} \frac{d^2 \log \zeta_2}{4\pi^2 \tau_2} D_{a,a}(q; \zeta/\zeta_2) D_{b,b}(q; \zeta_2) D_{c,c}(q; \zeta_2). \quad (4.3)$$

Setting  $\zeta = 1$  leads to the modular graph function  $C_{a,b,c}(q)$

$$C_{a,b,c}(q) = C_{a,b,c}(q; 1), \quad (4.4)$$

associated with the Feynman graph



$$= C_{a,b,c}(q; \zeta_1/\zeta'_1).$$

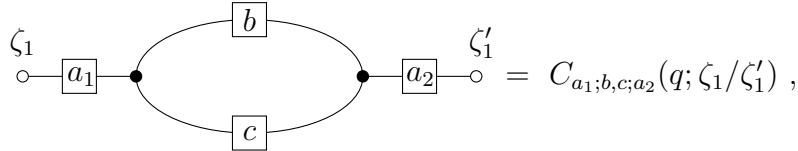
Since  $C_{a,b,c}(q; \zeta)$  is a single-valued function of  $\zeta$  it follows that  $C_{a,b,c}(q)$  is single-valued, which illustrates the general point made in sections 3.3 and 3.4.

Using the lattice sum expression for  $D_{a,a}(q; \zeta)$  given in (2.33) leads to the multiple sum representation studied in [9, 10]

$$C_{a,b,c}(q) = \sum'_{p_1, p_2, p_3 \in \mathbb{Z}\tau + \mathbb{Z}} \left(\frac{\tau_2}{\pi}\right)^{a+b+c} \frac{\delta(p_1 + p_2 + p_3)}{|p_1|^{2a} |p_2|^{2b} |p_3|^{2c}}, \quad (4.5)$$

where  $p_i = m_i\tau + n_i$  parametrizes the momenta on the torus  $\Sigma$ . In [9] linear combinations of the functions  $C_{a,b,c}(q)$  were shown to satisfy Laplace eigenvalue equations with source terms that are quadratic in real analytic Eisenstein series. In principle, these equations determine the form of these functions.

When at least one of the indices  $a, b, c$  is greater than one, the graph  $\Gamma$  contains bivalent vertices in addition to its two trivalent vertices. When this is the case, one may detach the edges meeting at any of the bivalent vertices, to obtain another single-valued elliptic multiple polylogarithm, specified graphically by the figure below



$$= C_{a_1; b, c; a_2}(q; \zeta_1/\zeta'_1),$$

where  $a_1 + a_2 = a$  and  $a_1, a_2 \geq 0$ . This case provides the simplest example of a general phenomenon, namely that to a given modular graph function, there will correspond a number of different and generally inequivalent single-valued elliptic multiple polylogarithm functions. Finally, we note that the above modular two-loop graph functions produce, upon detaching both edges on a single trivalent vertex, a star graph with three branches, and thereby a two-variable single-valued elliptic multiple polylogarithm. We shall examine such graphs further in section 5.1 below.

**4.3. General three-vertex graphs.** The modular graph function associated with an arbitrary graph with a total of three vertices is obtained as a special value of the single-valued elliptic multiple polylogarithm associated with the following graph with two unintegrated vertices

$$\zeta_1 \circ \text{---} (a) \text{---} \bullet \text{---} (b) \text{---} \bullet \text{---} (c) \text{---} \circ \zeta'_1 = D_{a,b,c}(q; \zeta_1/\zeta'_1),$$

where the encircled numbers indicate the number of the Green functions  $D_{1,1}$  which appear in the corresponding link, following our earlier graphical representations. The corresponding single-valued elliptic multiple polylogarithm evaluates as follows

$$D_{a,b,c}(q; \zeta) = \prod_{k=2,3} \int_{\Sigma} \frac{d^2 \log \zeta_k}{4\pi^2 \tau_2} D_{1,1}(q; \zeta/\zeta_2)^a D_{1,1}(q; \zeta_2/\zeta_3)^b D_{1,1}(q; \zeta_3)^c, \quad (4.6)$$

which is a single-valued function of  $\zeta$ . Its restriction to  $\zeta = 1$  gives

$$D_{a,b,c}(q) = D_{a,b,c}(q; 1), \quad (4.7)$$

corresponding to the modular graph function associated with the figure below

$$\begin{array}{c} \bullet \\ / \quad \backslash \\ (a) \quad (c) \\ \backslash \quad / \\ \bullet \text{---} (b) \text{---} \bullet \end{array} = D_{a,b,c}(q).$$

**4.4. General graphs.** The examples of entire classes of modular graph functions, and their associated single-valued elliptic multiple polylogarithms, provided in the earlier paragraphs lend themselves to a natural generalization. From a general Feynman graph  $\Gamma$ , with associated modular graph function  $I_{\Gamma}(q)$ , it is possible to construct at least one, but in general several inequivalent, single-valued elliptic multiple polylogarithms  $I_{\Gamma}(q; \zeta)$  by detaching one or several Green function edges from any one of the vertices in the graph. The original vertex may be placed at an arbitrary reference point, which we choose to be 1, while the vertex of the detached Green function(s) is placed at an arbitrary point  $\zeta$ . These observations show the validity of Proposition 3.4 and provide an explicit construction for  $I_{\Gamma}(q; \eta)$ . The significance of the fact that several inequivalent single-valued elliptic multiple polylogarithms arise in this way from a single modular graph function, remains to be explored in full.

## 5. SINGLE-VALUED ELLIPTIC MULTIPLE POLYLOGARITHMS FOR GRAPHS

In the preceding two sections, we have introduced modular graph functions, and their relation with single-valued elliptic multiple polylogarithms, leading to Proposition 3.4, and a series of concrete examples. In the present section, we shall make progress towards understanding the description in terms of (non-single-valued elliptic) polylogarithms of both modular graph functions and single-valued elliptic multiple

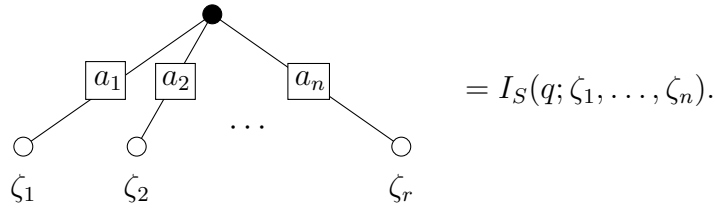
polylogarithms. In particular, we shall advance a *Conjecture* as to this general structure, and offer a proof for a certain infinite subclass of graphs, leaving a complete proof for the general case for future investigations. Assuming the validity of our conjecture, we shall infer important properties for the ring structure over  $\mathbb{Q}$  of the coefficients of the Laurent series in powers of  $\tau_2$  of general modular graph functions in terms of single-valued multiple zeta values. The Laurent series is the dominant contribution to these modular functions in the limit  $\tau_2 \rightarrow \infty$ .

The main conjecture of this paper applies to a modular graph function  $I_\Gamma(q)$  and an associated single-valued elliptic multiple polylogarithm  $I_\Gamma(q; \zeta)$  for a general *single-component graph*. By single-component graph, we mean a connected graph, whose associated modular graph function cannot be reduced to the product of lower weight modular graph functions.

### Conjecture

*The single-valued elliptic multiple polylogarithm  $I_\Gamma(q; \zeta)$  associated with a single component graph  $\Gamma = \Gamma_{w,N}$ , with  $w$  Green functions,  $N + 1$  vertices, and  $L = w - N$  loops, is a linear combination of elliptic multiple polylogarithms of depth at most  $L$  and weight at most  $w + N - 1$ .*

**5.1. Star graphs.** We shall now provide a proof of the above Conjecture for the infinite class of *star graphs*. A graph  $S$  is a star graph provided it contains a single integrated  $n$ -valent vertex with  $n > 2$ , and an arbitrary number of integrated bi-valent vertices. By construction, the graph then has  $n$  unintegrated end points  $\zeta_i$  with  $i = 1, \dots, n$ . To obtain the modular graph function with two  $n$ -valent vertices associated with  $S$  it suffices to set all  $\zeta_i$  equal to 1. To obtain single-valued elliptic multiple polylogarithms from  $S$ , one may set one non-empty subset of end points to 1 while setting the remaining non-empty set of endpoints to  $\zeta$ . A general  $n$ -valent star graph  $S$ , with  $a_i$  Green functions on leg  $i$  of the star, along with its multi-variable elliptic function value  $I_S(q; \zeta_1, \dots, \zeta_n)$ , is represented in the figure below



Its associated multi-variable modular graph function evaluates to

$$I_S(q; \zeta_1, \dots, \zeta_n) = \int_{\Sigma} \frac{d^2 \log \zeta}{4\pi^2 \tau_2} \prod_{i=1}^n D_{a_i, a_i}(q; \zeta / \zeta_i). \quad (5.1)$$

To prove the Conjecture for star graphs, we need to show that the result of performing the integration of the  $n$ -valent vertex  $\zeta$  over the torus  $\Sigma$  produces a linear combination, with rational coefficients, of (non-single-valued elliptic) multiple polylogarithms, resulting in a formula schematically represented as follows

$$\sum_s \sum_{b_1, \dots, b_s} c_{b_1, \dots, b_s} \text{Li}_{b_1, \dots, b_s}(z_1, \dots, z_s) . \quad (5.2)$$

where the coefficients  $c_{b_1, \dots, b_s}$  are rational numbers, the multiple polylogarithm is defined in (2.14), and  $z_i$  are products of the arguments, generically of the form

$$q^m \bar{q}^{m'} \prod_t \zeta_t^{\alpha_t} \bar{\zeta}_t^{\beta_t} , \quad (5.3)$$

where  $\alpha_t, \beta_t$  are integers with  $1 \leq t \leq s$ . We will proceed by considering some algebraic features of the integrand

$$\prod_{i=1}^n D_{a_i, a_i}(q; \zeta/\zeta_i) , \quad (5.4)$$

making use of the fact that the generalised Green functions  $D_{a_i, a_i}(q; \zeta/\zeta_i)$  are themselves elliptic polylogarithm of depth 1 and weight  $a_i$  given in (2.32) via (2.10), and then perform the integrations over the phase  $v$  and the variable  $u$  of  $\zeta = e^{2\pi i(u\tau+v)}$ .

**5.2. Integral over the phase  $v$ .** We begin by recasting the functions  $D_{a_i, a_i}(q; \zeta/\zeta_i)$  in terms of the functions  $D_{a_i, a_i}(q^m(\zeta/\zeta_i)^{\pm 1})$  and their complex conjugate with the help of equation (2.36), and then expressing the latter as a linear combination of polylogarithms  $\text{Li}_k$  using (2.10). In performing the  $v$ -integral, the presence of the Bernoulli polynomials  $B_{2a}(u)$  will be immaterial, since they involve  $u$  but do not depend on  $v$ , and we shall suppress their presence here. Finally, the arguments entering the polylogarithms  $\text{Li}_k$  in  $D_{a_i, a_i}$  depend on  $\zeta$  either through a factor linear in  $\zeta$ , or a factor linear in  $\zeta^{-1}$ . It will be convenient in the sequel to group together the factors in  $\zeta$  and those in  $\zeta^{-1}$ , of which we shall assume there are respectively  $r$  and  $s$ . The maximal values of these parameters are characterized by  $0 \leq r, s$  and  $r + s \leq n$ . Thus, the  $v$ -integrals required to evaluate  $I_S$  take the form

$$J_{\mathbf{a}, \mathbf{b}}(\mathbf{x}, \mathbf{y}) = \int_0^1 dv \prod_{i=1}^r \text{Li}_{a_i}(x_i e^{2\pi i v}) \prod_{j=1}^s \text{Li}_{b_j}(y_j e^{-2\pi i v}) . \quad (5.5)$$

where we have introduced the following notation,

$$\begin{aligned} \mathbf{a} &= (a_1, \dots, a_r) & \mathbf{x} &= (x_1, \dots, x_r) \\ \mathbf{b} &= (b_1, \dots, b_s) & \mathbf{y} &= (y_1, \dots, y_s) . \end{aligned} \quad (5.6)$$

The entries of the composite indices  $\mathbf{a}$  and  $\mathbf{b}$  are positive integers, while the entries of  $\mathbf{x}$  and  $\mathbf{y}$  depend on  $u$ , and on positive powers of  $q$  and positive powers of  $\bar{q}$ , but are independent of  $v$ . The variables obey  $|x_i|, |y_j| < 1$  for all  $i = 1, \dots, r$  and  $j = 1, \dots, s$ .

In view of these bounds, we may use the series representations for the functions Li to perform the integration over  $v$  which results in the insertion into the sum of a Kronecker  $\delta$ -function,

$$J_{\mathbf{a},\mathbf{b}}(\mathbf{x}, \mathbf{y}) = \sum_{0 < m_1, \dots, m_r} \sum_{0 < n_1, \dots, n_s} \delta \left( \sum_{i=1}^r m_i - \sum_{j=1}^s n_j \right) \prod_{i=1}^r \frac{x_i^{m_i}}{m_i^{a_i}} \prod_{j=1}^s \frac{y_j^{n_j}}{n_j^{b_j}}. \quad (5.7)$$

To prove that this constrained sum may be expressed as a linear combination, with rational coefficients, of functions Li evaluated on various combinations of the variables  $\mathbf{x}, \mathbf{y}$ , we shall proceed by induction on  $s$ , at fixed but arbitrary value of  $r$ . For  $r > 0$ , and  $s = 0$ , the integral vanishes identically.

To initiate the induction, consider first the case  $s = 1$ . It is straightforward to solve for  $n_1$  which is given by  $n_1 = m_1 + \dots + m_r$ . The constraint  $n_1 > 0$  imposed by the definition of  $J_{\mathbf{a},\mathbf{b}}$  is automatically satisfied since all the integers  $m_i$  are strictly positive. Thus, we find

$$J_{\mathbf{a},b_1}(\mathbf{x}, y_1) = \sum_{0 < m_1, \dots, m_r} \frac{1}{m_1 + \dots + m_r} \prod_{i=1}^r \frac{(x_i y_1)^{m_i}}{m_i^{a_i}}. \quad (5.8)$$

The summation may be expressed equivalently as the following integral

$$J_{\mathbf{a},b_1}(\mathbf{x}, y_1) = \int_0^1 \frac{dt}{t} \prod_{i=1}^r \text{Li}_{a_i}(x_i y_1 t). \quad (5.9)$$

Using the stuffle relations of (2.17), the product over the Li-functions may be expressed as a linear combination, with rational coefficients, of multiple polylogarithms. The one of highest depth  $r$  is given by  $\text{Li}_{a_1, \dots, a_r}(x_1 y_1 t, \dots, x_r y_1 t)$ , while the one of lowest depth 1 is given by  $\text{Li}_{a_1 + \dots + a_r}(x_1 \dots x_r y_1^r t^r)$ , along with all stuffle combinations with rational coefficients and depths in between 1 and  $r$ . The integral over  $t$  of any of these Li-functions gives back an Li-function of various combinations of the arguments. This concludes the proof of the claim for  $s = 1$ .

Next, we shall assume that the claimed decomposition property of  $J_{\mathbf{a},\mathbf{b}}$  defined in (5.5) holds true for all indices  $\mathbf{a}, \mathbf{b}$  for any value of  $r, s$ . We now wish to prove by induction that the decomposition property will then also hold for all indices  $\mathbf{a}, \mathbf{b}'$  for all  $r$ , and  $s$  replaced by  $s + 1$ . The corresponding  $J$ -function is given by

$$J_{\mathbf{a},\mathbf{b}'}(\mathbf{x}, \mathbf{y}') = \sum_{\substack{0 < m_1, \dots, m_r \\ 0 < n_1, \dots, n_s, n_{s+1}}} \delta \left( \sum_{i=1}^r m_i - \sum_{j=1}^s n_j - n_{s+1} \right) \prod_{i=1}^r \frac{x_i^{m_i}}{m_i^{a_i}} \prod_{j=1}^{s+1} \frac{y_j^{n_j}}{n_j^{b_j}}. \quad (5.10)$$

and we use the notation  $\mathbf{b}' = (b_1, \dots, b_s, b_{s+1})$  and  $\mathbf{y}' = (y_1, \dots, y_s, y_{s+1})$ .

To solve the Kronecker  $\delta$ -function constraint now requires more care than in the case  $s = 1$ . One cannot simply use the solution  $n_{s+1} = m_1 + \dots + m_r - (n_1 + \dots + n_s)$ , since then  $n_{s+1}$  is not guaranteed to obey the constraint that it must be positive. Therefore,

we shall proceed as follows instead. We single out an arbitrary  $m$ -variable, say  $m_1$ , and compare the values of  $n_{s+1}$  allowed by the  $\delta$ -function constraint to the value of  $m_1$ . We have three cases, namely  $n_{s+1} = m_1$ ,  $n_{s+1} < m_1$  or  $n_{s+1} > m_1$ , and split the sum which defines the function  $J$  accordingly into three terms

$$J_{\mathbf{a},\mathbf{b}'}(\mathbf{x}, \mathbf{y}') = J^0 + J^+ + J^-, \quad (5.11)$$

where the terms  $J^0$ ,  $J^+$ , and  $J^-$  respectively include the contributions to the sum from  $n_{s+1} = m_1$ ,  $n_{s+1} < m_1$ , and  $n_{s+1} > m_1$ .

For the contributions to the sum with  $n_{s+1} < m_1$ , which are regrouped in  $J^+$ , we parametrize the difference by  $m_1 = n_{s+1} + k$  with  $k$  ranging over all positive integers. The sum becomes,

$$J^+ = \sum_{\substack{0 < k, m_2, \dots, m_r \\ 0 < n_1, \dots, n_s, n_{s+1}}} \delta \left( k + \sum_{i=2}^r m_i - \sum_{j=1}^s n_j \right) \frac{x_1^{k+n_{s+1}}}{(k+n_{s+1})^{a_1}} \prod_{i=2}^r \frac{x_i^{m_i}}{m_i^{a_i}} \prod_{j=1}^{s+1} \frac{y_j^{n_j}}{n_j^{b_j}}. \quad (5.12)$$

We re-baptize the variable  $k$  as  $m_1$ , setting  $k = m_1$  in the sum. As a result, we have

$$J^+ = \sum_{\substack{0 < m_1, \dots, m_r \\ 0 < n_1, \dots, n_s, n_{s+1}}} \delta \left( \sum_{i=1}^r m_i - \sum_{j=1}^s n_j \right) \frac{x_1^{m_1+n_{s+1}}}{(m_1+n_{s+1})^{a_1}} \prod_{i=2}^r \frac{x_i^{m_i}}{m_i^{a_i}} \prod_{j=1}^{s+1} \frac{y_j^{n_j}}{n_j^{b_j}}. \quad (5.13)$$

Note that the  $\delta$ -function constraint involves only the first  $r+s$  summation variables but does not involve  $n_{s+1}$ . This observation is crucial to make the inductive proof possible. The summations over  $m_1$  and  $n_{s+1}$  are coupled only through the denominator factor  $(m_1+n_{s+1})^{-a_1}$  but not through the Kronecker  $\delta$ -function constraint. To complete the inductive argument, we decouple the summations over  $m_1$  and  $n_{s+1}$  by differentiating in  $x_1$ . The independent sums are then easily regrouped, and lead to the following differential equation

$$\left( x_1 \frac{\partial}{\partial x_1} \right)^{a_1} J^+ = \text{Li}_{b_{s+1}}(x_1 y_{s+1}) \text{Li}_{0, a_2, \dots, a_r, b_1, \dots, b_s}(x_1, \dots, x_r, y_1, \dots, y_s). \quad (5.14)$$

Stuffle identities (2.17) may again be used to decompose this product into a linear combination, with rational coefficients, of Li-functions of depth  $r+s+1$ . The above differential equation for  $J^+$  may then be integrated iteratively using (2.15) in terms of linear combinations of Li-functions of depth  $r+s+1$ . This concludes the part of the proof for the contributions  $J^+$  to  $J_{\mathbf{a},\mathbf{b}}(\mathbf{x}, \mathbf{y})$ .

For the contributions to the sum with  $n_{s+1} > m_1$ , which are regrouped in  $J^-$ , we parametrize the difference by  $m_1 + k = n_{s+1}$  with  $k$  ranging over all positive integers. The arguments now proceed as in the case of  $J^+$  with labels of positive and negative powers of  $\zeta$  reversed.

Finally, for the contributions to the sum with  $n_{s+1} = m_1$ , which are regrouped in  $J^0$ , we have

$$J^0 = \sum_{\substack{0 < m_1, \dots, m_r \\ 0 < n_1, \dots, n_s, n_{s+1}}} \delta \left( \sum_{i=2}^r m_i - \sum_{j=1}^s n_j \right) \frac{(x_1 y_{s+1})^{m_1}}{m_1^{a_1 + b_{s+1}}} \prod_{i=1}^r \frac{x_i^{m_i}}{m_i^{a_i}} \prod_{j=1}^s \frac{y_j^{n_j}}{n_j^{b_j}}. \quad (5.15)$$

The sum over  $m_1$  is now completely decoupled from the other sums, and we have

$$J^0 = \text{Li}_{a_1 + b_{s+1}}(x_1 y_{s+1}) \text{Li}_{a_2, \dots, a_r, b_1, \dots, b_s}(x_2, \dots, x_r, y_1, \dots, y_s) \quad (5.16)$$

Applying the stuffle relations of (2.17) allows us to decompose this product again into a linear combination with rational coefficients of Li-functions. This completes the proof for all contributions to  $J_{\mathbf{a}, \mathbf{b}}(\mathbf{x}, \mathbf{y})$ , and thus of this function itself.

**5.3. Arguments for the  $v$ -integral in the general case.** While the proof of the Conjecture given in the preceding section is complete for star graphs, and graphs derived from it, we need further results to support the Conjecture for general graphs. In this subsection, our aim will not be to provide a complete proof for general graphs, but rather to show that some of the key novel results that are needed for general graphs but were not required for star graphs, can be established in a satisfactory way.

One formulation may be obtained by considering the  $v$ -integral of a general multiple polylogarithm, but with arguments given by higher powers of  $\zeta$  and  $\zeta^{-1}$  (for star graphs only arguments linear in  $\zeta$  and  $\zeta^{-1}$  were required). This will lead to a general  $v$ -integral of the form

$$\mathcal{J}_{\mathbf{a}}^{\gamma}(\mathbf{z}) = \int_0^1 dv \text{Li}_{\mathbf{a}}(z_1 e^{2\pi i \gamma_1 v}, \dots, z_r e^{2\pi i \gamma_r v}), \quad (5.17)$$

where the exponent  $\gamma = (\gamma_1, \dots, \gamma_r)$  has integer entries, while  $\mathbf{a} = (a_1, \dots, a_r)$  has positive integer entries, and the arguments are again bounded by  $|z_i| < 1$  for all  $i = 1, \dots, r$ . By expanding the Li-function in a power series in  $z_i$ , we may readily perform the integral over  $v$  to obtain

$$\mathcal{J}_{\mathbf{a}}^{\gamma}(\mathbf{z}) = \sum_{0 < \mu_1 < \dots < \mu_r} \delta \left( \sum_{i=1}^r \gamma_i \mu_i \right) \prod_{i=1}^r \frac{z_i^{\mu_i}}{\mu_i^{a_i}}, \quad (5.18)$$

where the summation is over ordered positive integers  $\mu_1, \dots, \mu_r$ . The case considered earlier corresponds to having  $\gamma_i = \pm 1$ , so that the Kronecker  $\delta$ -function constraint may be solved to provide an iterative proof as we did in the preceding subsection. When  $\gamma_i \neq \pm 1$ , however, one needs to deal with the issues of divisibility in solving the constraint equation imposed by the Kronecker  $\delta$ -function.

We begin by parametrizing the ordering conditions on the indices  $\mu_i$  by setting  $\mu_0 = 0$  and introducing the following change of variables

$$\mu_i = \mu_{i-1} + m_i \quad 0 < m_i \quad 1 \leq i \leq r. \quad (5.19)$$



The variables  $m_i$  are now constrained only to be positive, but no further ordering requirement between them is being imposed. Importantly, this change of variables has the effect of decoupling the ordering conditions between different  $m_i$ . Next, it will be convenient to express  $\gamma$  and  $\mathbf{z}$  in this new basis and we define

$$\alpha_i = \sum_{j=i}^r \gamma_j \quad x_i = \prod_{j=i}^r z_j. \quad (5.20)$$

As a result, the function  $\mathcal{J}_{\mathbf{a}}^\gamma(\mathbf{z})$  admits the following sum representation

$$\mathcal{J}_{\mathbf{a}}^\gamma(\mathbf{z}) = \sum_{0 < m_1, \dots, m_r} \delta \left( \sum_{i=1}^r \alpha_i m_i \right) \prod_{i=1}^r \frac{x_i^{m_i}}{\mu_i^{\alpha_i}}. \quad (5.21)$$

To save unnecessary extra notation, the denominator factors have been left in terms of  $\mu_i$  and should be viewed as given in terms of the functions of the  $m_i$  by (5.19). Without loss of generality, we shall assume that all  $\alpha_i$  are non-vanishing; in the contrary case, the corresponding summation indices do not enter into the Kronecker  $\delta$ -function constraint and the sum reduces to a case with lower value of  $r$ .

Next, we partition the set of indices  $\{1, 2, \dots, r\}$  into two sets  $I_+$  and  $I_-$  according to whether  $\alpha_i$  is positive or negative

$$\begin{aligned} \alpha_i > 0 &\implies i \in I_+ \\ \alpha_i < 0 &\implies i \in I_- . \end{aligned} \quad (5.22)$$

It will be convenient to work with positive exponents only, so we shall set  $\beta_i = -\alpha_i$  whenever  $i \in I_-$ , so that  $\beta_i > 0$ . The summation may then be rearranged as follows

$$\mathcal{J}_{\mathbf{a}}^\gamma(\mathbf{z}) = \sum_{0 < m_1, \dots, m_r} \delta \left( \sum_{i \in I_+} \alpha_i m_i - \sum_{i \in I_-} \beta_i m_i \right) \prod_{i=1}^r \frac{x_i^{m_i}}{\mu_i^{\alpha_i}}. \quad (5.23)$$

We now proceed iteratively on the value of  $n$ , and assume that the extended array  $\gamma' = (\gamma_1, \dots, \gamma_r, \gamma_{r+1})$  produces one additional entry belonging to  $I_-$ , which corresponds to an additional exponent  $\beta_{i_-} > 0$  and an additional summation variable  $m_{i_-} > 0$ .

We now wish to reduce the Kronecker  $\delta$ -function involving  $n+1$  summation variables  $m$  to one that involves only  $n$  summation variables, imitating the procedure we followed for star graphs. To this end, we single out an index  $i_+$  in  $I_+$  corresponding to  $m_{i_+} \alpha_{i_+}$ . Thus, we need to reduce the following constraint on  $n+1$  summation variables  $m_1, \dots, m_{n+1}$

$$\sum_{i \in I_+ \setminus \{i_+\}} \alpha_i m_i - \sum_{i \in I_-} \alpha_i m_i + \alpha_{i_+} m_{i_+} - \beta_{i_-} m_{i_-} = 0, \quad (5.24)$$

to a summation involving only  $n$  summation variables. We proceed as in the case of star graphs, but take careful account of the divisibility issues. We shall distinguish contributions for which  $\alpha_{i_+} m_{i_+} = \beta_{i_-} m_{i_-}$  from those for which  $\alpha_{i_+} m_{i_+} > \beta_{i_-} m_{i_-}$  and

those for which  $\alpha_{i_+} m_{i_+} < \beta_{i_-} m_{i_-}$ . We shall refer to these partial sums respectively as  $\mathcal{J}^0$ ,  $\mathcal{J}^+$ , and  $\mathcal{J}^-$ .

We shall concentrate here on the contribution  $\mathcal{J}^+$ , arising from the contributions satisfying  $\alpha_{i_+} m_{i_+} > \beta_{i_-} m_{i_-}$ . The other cases are analogous, following the model of the star graphs. The integers  $\alpha_{i_+}$  and  $\beta_{i_-}$  are both positive, and have been assumed to be non-zero. We shall denote their greatest common divisor by  $\alpha_\star > 0$ . There now exists a unique solution  $m_{i_\pm}^0$  to the Bézout equation

$$\alpha_{i_+} m_{i_+}^0 - \beta_{i_-} m_{i_-}^0 = \alpha_\star \quad (5.25)$$

with the properties  $0 \leq m_{i_+}^0 < \beta_{i_-}/\alpha_\star$  and  $0 \leq m_{i_-}^0 < \alpha_{i_+}/\alpha_\star$ . Note that this solution is determined completely by the nature of  $\alpha_{i_+}$  and  $\beta_{i_-}$  and does not in any way involve the summation variables  $m$ . Having determined  $m_{i_\pm}^0$ , we may now recast the constraint as follows

$$\sum_{i \in I_+ \setminus \{i_+\}} \alpha_i m_i - \sum_{i \in I_-} \alpha_i m_i + \alpha_\star k = 0 \quad (5.26)$$

for an integer  $k$  which is positive by the assumption  $\alpha_{i_+} m_{i_+} > \beta_{i_-} m_{i_-}$ . As a result, the Kronecker  $\delta$ -function constraint now involves only  $n$  summation variables. The complete parametrization of the summation variables  $m_{i_\pm}$  is given in terms of the integer  $k > 0$ , and an additional integer summation variable  $\ell$ , as follows

$$\begin{aligned} m_{i_+} &= k m_{i_+}^0 + \ell \frac{\beta_{i_-}}{\alpha_\star} \\ m_{i_-} &= k m_{i_-}^0 + \ell \frac{\alpha_{i_+}}{\alpha_\star} \end{aligned} \quad (5.27)$$

The variable  $\ell$  does not enter into the constraint. In general, the range of  $\ell$  will include both positive, zero, and negative integers, and must itself be subdivided into these three sub-ranges. Within each sub-range, the summation over  $\ell$  may be performed, after sufficient derivatives in the external variables  $x_i$  have been taken to deal with the denominator factors  $\mu_i^{a_i}$ , in a manner analogous to how we proceeded for star graphs. This concludes the outline of the proof necessary to show that the integral over  $v$  produces a linear combination, with rational coefficients, of Li-functions.

**5.4. Integral over the variable  $u$ .** The second type of integral involves integration over one of the variables  $u_i$ , leading to the integrals of the type

$$\mathcal{K}_{a,\mathbf{a}}^\alpha(\mathbf{x}) = \int_0^1 du u^a \text{Li}_{a_1, \dots, a_r} (x_1 e^{-2\pi\tau_2 \alpha_1 u}, \dots, x_r e^{-2\pi\tau_2 \alpha_r u}) , \quad (5.28)$$

where the variables  $x_i$  are independent of  $u$  and bounded by  $|x_i| < 1$ , while  $a$ ,  $a_i$  and  $\alpha_i$  are positive or zero integers for all  $i = 1, \dots, r$ . We have introduced the notation  $\alpha = (\alpha_1, \dots, \alpha_r)$ . Actually, the integrals make perfect sense when  $\alpha_i$  is a positive integer, and it will sometimes be convenient for performing intermediate calculations, to let  $\alpha_i$  be real. The multiple polylogarithm arises from stuffle relations

in the expansion of products of Green functions on the torus, while the pre-factor  $u^a$  arises from expanding the Bernoulli polynomials in  $D_{a_i, a_i}(\zeta)$ .

It is sufficient to study the integrals  $\mathcal{K}_{0, \mathbf{a}}^\alpha$ , since cases with  $a > 0$  may be obtained by differentiating with respect to one the arguments  $\alpha_i$  (temporarily allowing it to be a real variable instead of an integer) of the polylogarithms using the differentiation rule of the multiple polylogarithms in (2.15). Using the sum representation of the multiple polylogarithms and the integral

$$\int_0^1 du e^{-2\pi\tau_2 u \alpha} = \frac{1}{2\pi\tau_2} \frac{1}{\alpha} (1 - e^{-2\pi\tau_2 \alpha}) \quad (5.29)$$

gives

$$\mathcal{K}_{0, \mathbf{a}}^\alpha(\mathbf{x}) = \frac{1}{2\pi\tau_2} \sum_{0 < m_1 < \dots < m_r} \prod_{i=1}^r \frac{x_i^{m_i}}{m_i^{\alpha_i}} \frac{1 - e^{-2\pi\tau_2 \sum_{i=1}^r \alpha_i m_i}}{\sum_{i=1}^r \alpha_i m_i}. \quad (5.30)$$

By differentiation we obtain

$$\left( \sum_{i=1}^r \alpha_i x_i \frac{d}{dx_i} \right) \mathcal{K}_{0, \mathbf{a}}^\alpha(\mathbf{x}) = \frac{1}{2\pi\tau_2} \left( \text{Li}_{a_1, \dots, a_r}(x_1, \dots, x_r) - \text{Li}_{a_1, \dots, a_r}(x_1 e^{-2\pi\tau_2 \alpha_1}, \dots, x_r e^{-2\pi\tau_2 \alpha_r}) \right). \quad (5.31)$$

Again, using the differential relations of multiple polylogarithms (2.15) one concludes that  $\mathcal{K}_{0, \mathbf{a}}^\alpha$  is a combination of multiple polylogarithm of depth  $r$  and weight  $\sum_{i=1}^r a_i + 1$ . The same argument implies that  $\mathcal{K}_{a, \mathbf{a}}^\alpha$  is a linear combination of multiple polylogarithms of depth  $r$  and weight  $\sum_{i=1}^r a_i + a$ .

## 6. FOURIER COEFFICIENTS OF MODULAR GRAPH FUNCTIONS

Having proved the Conjecture for the case of star graphs, and graphs derived from star graphs, and presented an outline of the evidence for the validity of the Conjecture for general graphs, we now proceed to discuss some implications of the Conjecture, assuming it holds true for general graphs.

**6.1. Implications for the constant Fourier mode.** Assuming the validity of the conjecture implies the following corollary of the conjecture concerning the constant Fourier mode of the Feynman integral  $I_\Gamma(q; 1) = I_\Gamma(q)$ .

### Corollary of the Conjecture

Let  $I_\Gamma^0(\tau_2)$  denote the constant Fourier mode of the modular graph function  $I_\Gamma(q)$  associated with a graph  $\Gamma = \Gamma_{w, N}$  with  $w$  Green functions and  $N$  vertices

$$I_\Gamma^0(\tau_2) = \int_{-\frac{1}{2}}^{\frac{1}{2}} d\tau_1 I_\Gamma(q). \quad (6.1)$$

- (1) *The constant Fourier mode  $I_\Gamma^0(\tau_2)$  has an expansion for large  $\tau_2$  consisting of a Laurent polynomial with a term of highest degree  $\tau_2^w$  and a term of lowest degree  $\tau_2^{1-w}$  with exponentially suppressed corrections of order  $\mathcal{O}(e^{-2\pi\tau_2})$ ;*
- (2) *The coefficient of the  $(\pi\tau_2)^w$  term in the Laurent polynomial is a rational number [9, 16, 17, 40];*
- (3) *The coefficients of the  $(\pi\tau_2)^k$  terms in the Laurent polynomial are single-valued multiple zeta values for  $1 - w \leq k \leq w - 1$ .*

The first two items are straightforward to establish. The third item follows from the Conjecture since the modular graph function  $I_\Gamma(q)$  results from evaluating a single-valued elliptic multiple polylogarithm at unit argument, and the single-valued elliptic multiple polylogarithm itself admits a decomposition as a linear combination, with rational coefficients, of single-valued multiple polylogarithms. Since single-valued multiple polylogarithms evaluated at unit argument yield single-valued multiple zeta functions, the validity of the Conjecture supports the validity of item (3) above.

We shall now provide examples of graphs, and their associated modular graph functions, for which these Laurent polynomials of the constant Fourier mode have been calculated, or for which the general structure of the corresponding Laurent polynomial is understood on general grounds.

**6.2. The Eisenstein series.** The simplest case where the implications of the Corollary on the constant Fourier mode may be easily verified is the non-holomorphic Eisenstein series itself. The constant Fourier mode consists of a Laurent polynomial in  $\tau_2$  with just two terms. Throughout, it will be convenient to give the Laurent series in powers of  $y = \pi\tau_2$  instead of  $\tau_2$ , given the structure of the coefficients announced in item (3) of the implications. For the Eisenstein series, we have

$$\int_{-\frac{1}{2}}^{\frac{1}{2}} d\tau_1 E_a(q) = (-1)^{a-1} \frac{B_{2a}}{(2a)!} (4y)^a + \frac{4(2a-3)!}{(a-2)!(a-1)!} \zeta(2a-1) (4y)^{1-a}. \quad (6.2)$$

The coefficients of powers of  $y$  are rational numbers multiplying odd  $\zeta$ -values which are single-valued according the definition used in [6–8]. The leading coefficient is a rational number involving the Bernoulli number  $B_{2a}$ , which is single-valued. This exemplifies the relation between the world-sheet graphs and the single-valued elliptic multiple polylogarithms (sveMPL's) in the very simplest cases.

**6.3. The  $C_{a,b,c}(q)$  modular graph functions.** Although explicit expressions for the Laurent series of the general two-loop modular graph  $C_{a,b,c}$  are not available in closed form, it is not difficult to verify the compatibility of their structure with the Corollary given above. First of all, the Laurent series of every specific  $C_{a,b,c}$  function evaluated explicitly in [9] is of the form predicted by the Corollary, with the highest power being  $y^w$  and the lowest  $y^{1-w}$ . The simplest case is  $C_{1,1,1} = E_3 + \zeta(3)$ , which is manifestly

compatible, in view of the Fourier mode expansion of the Eisenstein series given in the preceding section.

The cases of weight 4 and 5 worked out in [9] give the following Laurent expansions

$$\begin{aligned} \int_{-\frac{1}{2}}^{\frac{1}{2}} d\tau_1 C_{2,1,1}(q) &= \frac{2y^4}{14175} + \frac{\zeta(3)y}{45} + \frac{5\zeta(5)}{12y} - \frac{\zeta(3)^2}{4y^2} + \frac{9\zeta(7)}{16y^3} + \mathcal{O}(e^{-2y}), \\ \int_{-\frac{1}{2}}^{\frac{1}{2}} d\tau_1 C_{3,1,1}(q) &= \frac{2y^5}{155925} + \frac{2\zeta(3)y^2}{945} - \frac{\zeta(5)}{180} + \frac{7\zeta(7)}{16y^2} - \frac{\zeta(3)\zeta(5)}{2y^3} + \frac{43\zeta(9)}{64y^4} + \mathcal{O}(e^{-2y}), \end{aligned} \quad (6.3)$$

and these expressions are again compatible with the structure predicted by the Corollary since the coefficients only involve polynomials in odd  $\zeta$  values, which are MZVs.

This structure of the Laurent series is easily seen to follow from the structure of nested differential equations satisfied by the  $C_{a,b,c}(q)$  functions. Indeed, every Laplace-eigenvalue equation of [9] for linear combinations (with rational coefficients)  $\mathfrak{C}_{w,s,p}$  of  $C_{a,b,c}$ -functions of weight  $w = a + b + c$  and eigenvalue  $s = 1, 2, \dots, w - 2$ , has an inhomogeneous part consisting of a term linear in the Eisenstein series  $E_w$  multiplied by a rational number plus terms of the form  $E_{w_1}E_{w_2}$  with  $w_1, w_2 \geq 2$  and  $w = w_1 + w_2$ , and multiplied by rational coefficients. Therefore, the inhomogeneous part involves a Laurent polynomial with powers  $y^k$  ranging from  $1 - w \leq k \leq w$ , the coefficient of the term  $y^w$  being rational, and the remaining coefficients being odd  $\zeta$ -values, which by definition are proportional to single-valued zeta values with rational coefficients. We may represent this schematically by

$$\left( y^2 \frac{\partial^2}{\partial y^2} - s(s-1) \right) \mathfrak{C}_{w,s,p} \in \mathbb{Q}E_w + \sum_{\substack{w_1, w_2 \geq 2 \\ w_1 + w_2 = w}} \mathbb{Q}E_{w_1}E_{w_2}. \quad (6.4)$$

The operator on the left hand side of the equation has a kernel consisting of the Laurent terms proportional to  $y^s$  and  $y^{1-s}$ . The Laplace eigenvalue equations therefore imply that almost all the coefficients in the Laurent polynomial for  $\mathfrak{C}_{w,s,p}$ , except for the highest coefficient of  $y^w$ , are odd zeta values or products thereof, as indeed predicted by the Corollary.

**6.4. Examples of higher cases.** Some Laurent expansions of modular graph functions that do not belong to the family of graphs  $C_{a,b,c}$  were also evaluated in [9]. Here, we give just two examples of Laurent series at weight 5, with  $y = \pi\tau_2$

$$\begin{aligned} \int_{-\frac{1}{2}}^{\frac{1}{2}} d\tau_1 D_{3,1,1}(q) &= \frac{2y^5}{22275} + \frac{y^2\zeta(3)}{45} + \frac{11\zeta(5)}{60} + \frac{105\zeta(7)}{32y^2} - \frac{3\zeta(3)\zeta(5)}{2y^3} + \frac{81\zeta(9)}{64y^4} + \mathcal{O}(e^{-2y}), \\ \int_{-\frac{1}{2}}^{\frac{1}{2}} d\tau_1 D_{2,2,1}(q) &= \frac{8y^5}{467775} + \frac{4\zeta(3)y^2}{945} + \frac{13\zeta(5)}{45} + \frac{7\zeta(7)}{8y^2} - \frac{\zeta(3)\zeta(5)}{y^3} + \frac{9\zeta(9)}{8y^4} + \mathcal{O}(e^{-2y}). \end{aligned} \quad (6.5)$$

In each case, only odd zeta values, or products thereof, occur all of which are single-valued multi-zeta values. The preceding expressions do not involve non-trivial svMZV's (i.e. ones which are not reducible to polynomials in odd zeta values) since these only arise at weights  $\geq 11$ .

Explicit computations of Laurent series at higher weight have been performed recently by Zerbini [40]. These strikingly confirm the conjectured structure of the coefficients of the Laurent polynomials. For example, Zerbini's evaluation of the constant Fourier mode of  $D_{1,1,5}(q)$  (which is a special case of the class of star graphs) is given by

$$\begin{aligned} \frac{1}{47} \int_{-\frac{1}{2}}^{\frac{1}{2}} d\tau_1 D_{1,1,5}(q) &= \frac{62 y^7}{10945935} + \frac{\zeta_{sv}(3)}{243} y^4 + \frac{119}{648} \zeta_{sv}(5) y^2 + \frac{11}{54} \zeta_{sv}(3)^2 y \\ &+ \frac{21}{32} \zeta_{sv}(7) + \frac{23 \zeta_{sv}(3) \zeta_{sv}(5)}{6y} + \frac{7115 \zeta_{sv}(9) - 1800 \zeta_{sv}(3)^2}{576 y^2} \\ &+ \frac{1245 \zeta_{sv}(3) \zeta_{sv}(7) - 150 \zeta_{sv}(5)^2}{64 y^3} \\ &+ \frac{288 \zeta_{sv}(3, 5, 3) - 4080 \zeta_{sv}(5) \zeta_{sv}(3)^2 - 9573 \zeta_{sv}(11)}{256 y^4} \\ &+ \frac{2475 \zeta_{sv}(5) \zeta_{sv}(7) + 1125 \zeta_{sv}(9) \zeta_{sv}(3)}{128 y^5} - \frac{1575}{64} \frac{\zeta_{sv}(13)}{y^6} + \mathcal{O}(e^{-2y}). \quad (6.6) \end{aligned}$$

The fact that the coefficient of  $y^{-4}$  is a linear combination of depth 3 single-valued multiple zeta-values is a non-trivial result. Similarly, the coefficients of  $y^{-3}$  and  $y^{-5}$  in  $D_{1,1,6}$  are linear combinations with rational coefficients of svMZV of weight 11 and 13 respectively [40].

Another important aspect of the structure described in this paper is that the modular graph functions that arise at special values of the svMPL's can be expressed as integrals of sums of single-valued multiple polylogarithms. This means that their functional form can be determined by using the algebraic properties of multiple polylogarithms that were described in section 2. Although this is an interesting comment on the structure of these functions, it is quite complicated to evaluate the functions in this manner, even for the simplest nontrivial function which is  $C_{1,1,1}(q)$ , which is evaluated explicitly in appendix B. The result has previously been obtained by Zagier (unpublished) and in [9] by other methods.

## 7. SUMMARY AND THOUGHTS ON THE BASIS OF MODULAR GRAPH FUNCTIONS

The preceding discussion concerns properties of modular graph functions that are associated with Feynman diagrams for a free conformal scalar field on a torus. A general graph, denoted  $\Gamma = \Gamma_{w,N}$ , has  $w$  Green functions joining  $N$  vertices at positions  $\zeta_i$  ( $i = 1, \dots, N$ ) that are integrated over the torus. Each Green function is itself a

single-valued elliptic polylogarithm,  $D_{1,1}(q, \zeta_i/\zeta_j)$ . This procedure defines a modular function,  $I_\Gamma(q)$ , associated with the graph as in (3.9).

We also considered the graph with  $N + 1$  vertices that is obtained from  $\Gamma = \Gamma_{w,N}$  by integrating over only  $N - 1$  vertices and separating the Green functions joined to the  $N$ -th vertex into two groups that end at the points  $\zeta_N$  and  $\zeta_{N+1}$ . This is associated with the function  $I_\Gamma(q; \zeta)$  (3.13), which is a single-valued function of the argument  $\zeta = \zeta_{N+1}/\zeta_N$ . This is a particular example of a single-valued elliptic multiple polylogarithm. The value of this function at the point  $\zeta = 1$  is a modular graph function,  $I_\Gamma(q; 1) = I_\Gamma(q)$ .

As in the case of multiple zeta values the question of determining a basis of modular graph functions seems daunting. In considering this question we should bear in mind that the context that motivated our analysis was limited. The general Feynman diagrams with  $N$  vertices and  $w$  Green functions that were considered in this paper and in [9, 10] generalise those obtained by expanding the four-graviton superstring scattering amplitude, which only generates the diagrams with  $N \leq 4$ . This does not provide a complete set of basis functions. In fact, the relations between the functions that were proved or conjectured in [9] explicitly involve relations between functions with  $p \leq 4$  and those with  $p > 4$ .

The preceding comments parallel known properties of the expansion of the tree-level  $N$ -graviton amplitude, which generates a subset of single-valued multi-zeta values at any weight. For example for  $N = 4$  the expansion coefficients are polynomial in ordinary odd Riemann zeta values. In order to access the complete basis of MZVs it is necessary to expand the tree amplitudes for general  $N$  [4, 30]. By analogy, we would expect that a more complete understanding of the basis for the space of modular graph functions requires analysis of the low energy expansion of  $N$ -graviton genus-one amplitudes, which is presently rather limited. It is known [39] that for  $N = 5$  new kinds of modular graph functions appear for weights  $w > 4$ . These are graphs in which  $2N - 8$  Green functions contain a holomorphic or anti-holomorphic numerator momentum factor, which have the form

$$D_{1,0}(q; \zeta) = -\zeta \partial_\zeta D_{1,1}(q; \zeta), \quad D_{0,1}(q; \zeta) = -\bar{\zeta} \partial_{\bar{\zeta}} D_{1,1}(q; \zeta). \quad (7.1)$$

These functions transform as modular forms of non-zero modular weight as seen from (2.35). Such graphs are integrals of products of  $D_{a,b}(q; \zeta)$  with  $a, b = 0, 1$ , with the constraint that the integrand has zero net modular weight. For example the integrands of the modular graph functions of this type in the  $N = 5$  case have one  $D_{0,1}$  factor and one  $D_{1,0}$  factor.

Since the relationships between different modular graph functions are difficult to determine with present techniques, it is not clear how many independent modular weight functions there are at any weight.

The fact that the coefficients of the Laurent polynomials of the constant Fourier mode of modular graph functions are single-valued multiple zeta values suggests that

they may be related to coefficients in the low energy expansion of tree-level amplitudes. After all, in the limit  $\tau_2 \rightarrow \infty$  the genus-one  $N$ -graviton amplitude degenerates to a  $N + 2$ -particle tree amplitude at special values of the momenta. It would therefore be interesting to connect these coefficients to the svMZVs that appear in the low energy expansion of tree-level closed string amplitudes [4, 5].

A related issue is a possible connection between modular graph functions and holomorphic elliptic multiple polylogarithms of [13]. The latter arise in the low energy expansion of open string theory, as discussed in [14]. Such a connection would be analogous to the relation of ordinary multiple polylogarithms to single-valued multiple polylogarithms. Within string theory this corresponds to the KLT relation that builds the tree level closed string amplitudes from the open string tree amplitudes. There are strong hints that the integrands of the open string one-loop amplitude and the closed string genus-one amplitude considered in this paper are related in an analogous manner. Determining the details of such a genus-one ‘‘KLT’’ relation would be very enlightening.

#### ACKNOWLEDGMENTS

We would like to thank Erik Panzer and Don Zagier for many useful discussions and exchanges. We are very grateful to Federico Zerbini for providing us with his evaluation of the Laurent polynomials of various graphs prior to publication.

The research of ED was supported in part by a grant from the National Science Foundation (NSF) PHY-13-13986. The research of PV has received funding the ANR grant reference QST 12 BS05 003 01, and the CNRS grants PICS number 6430. The research of OG has received funding by European Research Council (Programme Ideas ERC-2012-AdG 320769 AdS-CFT-solvable) and by the ANR grant StrongInt (BLANC-SIMI-4-2011). PV is partially supported by a fellowship funded by the French Government at Churchill College, Cambridge. MBG and PV acknowledge funding from the European Research Council under the European Community’s Seventh Framework Programme (FP7/2007-2013) ERC grant agreement no. [247252].

#### APPENDIX A. POLYLOGARITHMS AND EICHLER INTEGRALS

We will here demonstrate an interesting connection between elliptic polylogarithms and integrals of holomorphic Eisenstein series. Part of this discussion is contained in [41, §7].

To begin with we will introduce a set of derivatives that shows how the  $D_{a,b}(q; \zeta)$  functions of section 2.7 can be expressed as derivatives of the polylogarithms and holomorphic Eisenstein series. Consider the differential operator

$$\mathbb{D}_a = (-1)^{a-1} \sum_{k=0}^{a-1} \binom{2a-2-k}{a-1} \frac{(-\log|x|^2)^k}{k!} \left( \frac{d}{d \log x} \right)^k. \quad (\text{A.1})$$



Making use of the differentiation relation (2.6) one can express  $D_{a,a}(x)$  as

$$D_{a,a}(x) = 2\Re(\mathbb{D}_a \text{Li}_{2a-1}(x)) . \quad (\text{A.2})$$

The action of  $\mathbb{D}_a$  can be promoted to the action of a covariant derivative

$$\mathcal{D}_a = (-1)^{a-1} \sum_{k=0}^{a-1} \binom{2a-2-k}{a-1} \frac{(4\pi\tau_2)^k}{k!} \left( \frac{d}{d \log q} \right)^k , \quad (\text{A.3})$$

which maps weight  $\ell$  modular forms to weight  $k + \ell$  modular forms. This differential operator can be expressed as the product of covariant derivatives [41, §7]

$$\mathcal{D}_a = (4\pi\tau_2)^{m-1} (\partial_{-2} \circ \partial_{-4} \circ \cdots \circ \partial_{2-2a}) , \quad (\text{A.4})$$

where  $\partial_k = \frac{1}{2i\pi} \partial_\tau - \frac{k}{4\pi\tau_2}$ . Equipped with  $\mathcal{D}_a$  one can rewrite (A.2) in the form

$$D_{a,a}(q; \zeta) = 2\Re \mathcal{D}_a \left[ \tilde{G}_{2a-1}(q; \zeta) \right] . \quad (\text{A.5})$$

where

$$\tilde{G}_{2a-1}(q; \zeta) = \sum_{n \geq 0} (\text{Li}_{2a-1}(q^n \zeta) + \text{Li}_{2a-1}(q^{n+1}/\zeta)) - \frac{(\log q)^{2a-1}}{(2a)!} B_{2a}(u) . \quad (\text{A.6})$$

Setting  $\zeta = 1$  and using  $B_{2a}(1) = -2a \zeta(1-2a)$  leads to an expression for the Eisenstein series as

$$E_a(q) = \frac{1}{(-4\pi\tau_2)^{a-1}} 2\Re \mathcal{D}_a \left[ \tilde{G}_{2a-1}(q; 1) \right] , \quad (\text{A.7})$$

where the function  $\tilde{G}_{2a-1}(q; 1)$  is given by

$$\tilde{G}_{2a-1}(q) = \zeta(1-2a) \frac{(\log q)^{2a-1}}{(2a-1)!} + \zeta(2a-1) + 2 \sum_{n=1}^{\infty} \text{Li}_{2a-1}(q^n) . \quad (\text{A.8})$$

In fact,  $\tilde{G}_{2a-1}(q)$  is an Eichler integral<sup>7</sup> that satisfies

$$\left( \frac{d}{d \log q} \right)^{2a-1} \tilde{G}_{2a-1}(q) = 2G_{2a}(q) . \quad (\text{A.9})$$

The quantity  $G_{2a}(q)$  weight  $2a$  holomorphic Eisenstein series for  $SL(2, \mathbb{Z})$ , defined by

$$G_{2a}(q) = \frac{(2a-1)!}{2(2i\pi)^{2a}} \sum_{(m,n) \neq (0,0)} \frac{1}{(m\tau + n)^{2a}} = \frac{1}{2} \zeta(1-2a) + \sum_{n=1}^{\infty} n^{2a-1} \frac{q^n}{1-q^n} . \quad (\text{A.10})$$

Since  $\text{Li}_0(x) = \frac{x}{1-x}$  the  $q$ -expansion is given by a weight 0 polylogarithm as implied by the relation (A.9). Combining (A.5) with (A.9) leads to a relation between the non-holomorphic Eisenstein series  $E_a(q)$  and the holomorphic Eisenstein series,  $G_{2a}(q)$ .

<sup>7</sup> There is a polynomial ambiguity  $\sum_{i=0}^{2a-2} a_i (\log q)^i$  where  $a_i$  are constants. This polynomial is the period of the modular form as discussed in [42].

At the end of the next appendix we will use a similar argument to obtain relations between both the modular functions  $E_3(q)$  and  $C_{1,1,1}(q)$  and the holomorphic Eisenstein series  $G_6(q)$ .

### APPENDIX B. THE $C_{1,1,1}(q)$ MODULAR GRAPH FUNCTION

In this section we illustrate how the connection between modular graph functions and single-valued multiple polylogarithms can be used to evaluate their functional form. We will here consider the first nontrivial example, which is the function  $C_{1,1,1}(q)$  that is defined by By definition

$$C_{1,1,1}(q) = \int_{\Sigma} \frac{d^2 \log \zeta}{4\pi^2 \tau_2} D_{1,1}(q; \zeta)^3, \quad (\text{B.1})$$

with  $D_{1,1}(q; \zeta)$  defined in (3.8) and  $\zeta = q^u e^{2i\pi v}$  with  $u, v \in [0, 1]$ .

We now expand the integrand using the stuffle relations of polylogarithms obtained using the series representation

$$\begin{aligned} \prod_{i=1}^3 \text{Li}_{a_i}(x_i) &= \text{Li}_{a_1+a_2+a_3}(x_1 x_2 x_3) + \sum_{\sigma \in \mathfrak{S}_3} \text{Li}_{a_{\sigma(1)}, a_{\sigma(2)}, a_{\sigma(3)}}(x_{\sigma(1)}, x_{\sigma(2)}, x_{\sigma(3)}) \\ &\quad + \sum_{i=1}^3 \delta_{\{i,j,k\}=\{1,2,3\}} \left( \text{Li}_{a_i, a_j+a_k}(x_i, x_j x_k) + \text{Li}_{a_j+a_k, a_i}(x_j x_k, x_i) \right), \quad (\text{B.2}) \end{aligned}$$

where  $x_i$  are the arguments of the polylogarithms that are of the form in the various terms of (3.8). The symbol  $\mathfrak{S}_3$  denotes the set of permutations of three elements, and  $\delta_{\{i,j,k\}=\{1,2,3\}}$  enforces the constraint that  $(i, j, k)$  is a permutation of  $(1, 2, 3)$ . We apply this identity with  $a_i = 1, 2$

As an illustration we evaluate the contribution from the term in the integrand with no factors of  $Li_2$ . This has the form

$$\begin{aligned} I_{2,2,2} &= \frac{\tau_2^3}{\pi^3} \int_{\Sigma} \frac{d^2 \log \zeta}{4\pi^2 \tau_2} (\text{Li}_2(e^{2i\pi u}) + \text{Li}_2(e^{-2i\pi u}))^3 \\ &= \frac{6\tau_2^3}{\pi^3} \Re \int_0^1 du \text{Li}_2(e^{2i\pi u})^2 \text{Li}_2(e^{-2i\pi u}). \quad (\text{B.3}) \end{aligned}$$

Substituting the expansion (B.2)

$$\begin{aligned} \text{Li}_2(e^{2i\pi u})^2 \text{Li}_2(e^{-2i\pi u}) &= \text{Li}_6(e^{2i\pi u}) + 2\text{Li}_{2,2,2}(e^{2i\pi u}, e^{2i\pi u}, e^{-2i\pi u}) \\ &\quad + 2\text{Li}_{2,2,2}(e^{2i\pi u}, e^{-2i\pi u}, e^{2i\pi u}) + 2\text{Li}_{2,2,2}(e^{-2i\pi u}, e^{2i\pi u}, e^{2i\pi u}) \\ &\quad + 2\text{Li}_{2,4}(e^{2i\pi u}, 1) + 2\text{Li}_{4,2}(1, e^{2i\pi u}) + \text{Li}_{2,4}(e^{-2i\pi u}, e^{4i\pi u}) + \text{Li}_{4,2}(e^{4i\pi u}, e^{-2i\pi u}), \quad (\text{B.4}) \end{aligned}$$

and using the series expression for the multiple polylogarithms we find that the only non vanishing contributions are

$$\begin{aligned} I_{2,2,2} &= \frac{6\tau_2^3}{\pi^3} \Re \int_0^1 du \left( 2\text{Li}_{2,2,2}(e^{2i\pi u}, e^{2i\pi u}, e^{-2i\pi u}) + \text{Li}_{4,2}(e^{4i\pi u}, e^{-2i\pi u}) \right) \\ &= \frac{6\tau_2^3}{\pi^3} \Re \left( 2 \sum_{0 < m_1 < m_2} \frac{1}{m_1^2 m_2^2 (m_1 + m_2)^2} + \frac{1}{4} \text{Li}_6(1) \right). \end{aligned} \quad (\text{B.5})$$

The expression in parenthesis is real and is a special value of the multiple sum

$$L_{a,b;c}(x, y) = \sum_{m_1, m_2 \geq 1} \frac{x^{m_1} y^{m_2}}{m_1^a m_2^b (m_1 + m_2)^c}. \quad (\text{B.6})$$

since

$$\begin{aligned} L_{a,b;c}(x, y) &= \frac{1}{2^c} \sum_{0 < m} \frac{(xy)^m}{m^{a+b+c}} \\ &\quad + \sum_{0 < m_1 < m_2} \frac{x^{m_1} y^{m_2}}{m_1^a m_2^b (m_1 + m_2)^c} + \sum_{0 < m_2 < m_1} \frac{x^{m_1} y^{m_2}}{m_1^a m_2^b (m_1 + m_2)^c}. \end{aligned} \quad (\text{B.7})$$

The multiple sums  $L_{a,b;c}(x, y)$  may be reduced to a linear combination of multiple polylogarithms as shown in appendix C, with the result

$$L_{2,2,2}(x, y) = \text{Li}_{2,4}\left(\frac{y}{x}, x\right) + \text{Li}_{2,4}\left(\frac{x}{y}, y\right) + 2\text{Li}_{1,5}\left(\frac{y}{x}, x\right) + 2\text{Li}_{1,5}\left(\frac{x}{y}, y\right), \quad (\text{B.8})$$

and the value at  $x = y = 1$  is easily obtained using `Hyperint` routines of [23]

$$L_{2,2,2}(1, 1) = 2\text{Li}_{2,4}(1, 1) + 4\text{Li}_{1,5}(1, 1) = \frac{\zeta(6)}{3}. \quad (\text{B.9})$$

Therefore

$$I_{2,2,2} = 2\zeta(6) \frac{\tau_2^3}{\pi^3}. \quad (\text{B.10})$$

The remaining integrals are performed similarly. Collecting everything we have

$$C_{1,1,1}(q) = \frac{2\zeta(6)}{\pi^3} \tau_2^3 + \varphi_0(q) + \frac{\varphi_1(q)}{\tau_2} + \frac{\varphi_2(q)}{\tau_2^2}, \quad (\text{B.11})$$

where

$$\varphi_0 = \zeta(3) + 4 \sum_{n \geq 1} \Re \text{Li}_3(q^n) + 12 \sum_{n_1, n_2, n_3 \geq 1} \Re L_{1,1,1}(q^{n_1+n_3+1}, q^{n_2+n_3+1}), \quad (\text{B.12})$$

and

$$\begin{aligned} \varphi_1(q) &= -\frac{3}{\pi} \sum_{n_1, n_2 \geq 0} \Re(\text{Li}_4(q^{n_1} \bar{q}^{n_2}) + \text{Li}_4(q^{n_1+1} \bar{q}^{n_2+1})) \\ &\quad + \frac{3}{\pi} \sum_{n_1, n_2, n_3 > 0} \Re(L_{1,1;2}(q^{n_1} \bar{q}^{n_3}, q^{n_2} \bar{q}^{n_3}) - L_{1,1;2}(q^{n_1+1} \bar{q}^{n_3+1}, q^{n_2+1} \bar{q}^{n_3+1})) \\ &\quad + \frac{6}{\pi} \sum_{n_1, n_2, n_3 \geq 0} \Re(L_{2,1;1}(q^{n_1} \bar{q}^{n_3}, \bar{q}^{n_2+n_3+1}) - L_{2,1;1}(q^{n_1+1} \bar{q}^{n_3+1}, \bar{q}^{n_2+n_3+1})) , \end{aligned} \quad (\text{B.13})$$

and finally

$$\varphi_2(q) = -\frac{3}{4\pi^2} \zeta(5) + \frac{3}{2\pi^2} \sum_{n \geq 0} \Re \text{Li}_5(q^n) . \quad (\text{B.14})$$

We note that the Eisenstein series  $E_3(q)$  has an expansion that follows from (A.7) of the form

$$\begin{aligned} E_3(q) &= \frac{2\zeta(6)}{\pi^3} \tau_2^3 + \frac{3\zeta(5)}{4\pi^2 \tau_2^2} + \sum_{n \geq 1} \left( 2n^2 \Re \text{Li}_3(q^n) + \frac{3n}{\pi \tau_2} \Re \text{Li}_4(q^n) \right. \\ &\quad \left. + \frac{3}{2\pi^2 \tau_2^2} \Re \text{Li}_5(q^n) \right) . \end{aligned} \quad (\text{B.15})$$

Putting everything together we have

$$C_{1,1,1}(q) - E_3(q) - \zeta(3) = \Re \left( 2\phi_0(q) + \frac{3}{\pi} \frac{\phi_1(q)}{\tau_2} \right) , \quad (\text{B.16})$$

where

$$\phi_0(q) = 6 \sum_{n_1, n_2, n_3 \geq 0} L_{1,1;1}(q^{n_1+n_3+1}, q^{n_2+n_3+1}) + \sum_{n \geq 1} (1-n^2) \text{Li}_3(q^n) . \quad (\text{B.17})$$

and

$$\begin{aligned} \phi_1(q) &= -\sum_{n \geq 1} n \text{Li}_4(q^n) + \sum_{n_1, n_2 \geq 0} (\text{Li}_4(q^{n_1} \bar{q}^{n_2}) + \text{Li}_4(q^{n_1+1} \bar{q}^{n_2+1})) \\ &\quad - \sum_{n_1, n_2, n_3 \geq 0} (L_{1,1;2}(q^{n_1} \bar{q}^{n_3}, q^{n_2} \bar{q}^{n_3}) - L_{1,1;2}(q^{n_1+1} \bar{q}^{n_3+1}, q^{n_2+1} \bar{q}^{n_3+1})) \\ &\quad - 2 \sum_{n_1, n_2, n_3 \geq 0} (L_{2,1;1}(q^{n_1} \bar{q}^{n_3}, \bar{q}^{n_2+n_3+1}) - L_{2,1;1}(q^{n_1+1} \bar{q}^{n_3+1}, \bar{q}^{n_2+n_3+1})) . \end{aligned} \quad (\text{B.18})$$

It is striking that these identities are not reducible to standard polylogarithm identities and that they mix the  $q$  and  $\bar{q}$  expansion as is seen in the expression for  $\phi_1(q)$ .

Using the expressions for  $L_{a,b,c}(x, y)$  derived in appendix C

$$\begin{aligned} L_{1,1;1}(x, y) &= \text{Li}_{1,2}(y/x, x) + \text{Li}_{1,2}(x/y, y) \\ L_{1,1;2}(x, y) &= \text{Li}_{1,3}(y/x, x) + \text{Li}_{1,3}(x/y, y) \end{aligned}$$

$$L_{2,1;1}(x, y) = \text{Li}_{2,2}(x/y, y) + \text{Li}_{1,3}(y/x, x) + \text{Li}_{1,3}(x/y, y) . \quad (\text{B.19})$$

one can easily check that  $\phi_0(q) = 0 = \phi_1(q)$  to an arbitrary order in the  $q$  expansion using `HyperInt` [23]. It then follows from (B.16) that

$$C_{1,1,1}(q) = E_3(q) + \zeta(3) . \quad (\text{B.20})$$

**B.1. Remark on Eichler integrals.** Introducing the covariant derivative

$$\mathcal{D}_2 = -2 - 4\pi\tau_2 \frac{d}{d \log q} , \quad (\text{B.21})$$

we find that

$$E_3(q) = \frac{1}{(-4\pi\tau_2)^2} 2\Re \mathcal{D}_2 \tilde{G}_5(q) \quad (\text{B.22})$$

and

$$C_{1,1,1}(q) = -\frac{1}{(-4\pi\tau_2)^2} 2\Re \mathcal{D}_2 \hat{G}_5(q) , \quad (\text{B.23})$$

where

$$\tilde{G}_5(q) = \zeta(-5) \frac{(\log q)^6}{5!} + \zeta(5) + 2 \sum_{n=1}^{\infty} \text{Li}_5(q^n) \quad (\text{B.24})$$

and

$$\hat{G}_5(q) = \tilde{G}_5(q) + \frac{1}{2} \pi^3 \zeta(3) (\log q)^2 . \quad (\text{B.25})$$

Both of these function satisfy

$$\left( \frac{d}{d \log q} \right)^5 \tilde{G}_5(q) = \left( \frac{d}{d \log q} \right)^5 \hat{G}_5(q) = 2G_6(q) . \quad (\text{B.26})$$

where  $G_6(q)$  is the holomorphic Eisenstein series

$$G_6(q) = \frac{60}{(2i\pi)^6} \sum_{(m,n) \neq (0,0)} \frac{1}{(m\tau + n)^6} . \quad (\text{B.27})$$

This means that  $E_3(q)$  and  $C_{1,1,1}(q)$  are related to two Eichler integrals of the holomorphic weight 6 Eisenstein series for  $SL(2, \mathbb{Z})$ . The difference between  $\tilde{G}_5(q)$  and  $\hat{G}_5(q)$  is the polynomial ambiguity related to the period polynomial arising when integrating the holomorphic Eisenstein series. We refer back to the appendix A for a review of this construction.

## APPENDIX C. REDUCTION OF MULTIPLE SUMS TO MULTIPLE POLYLOGARITHMS

In this appendix we will reduce various multiple constrained sums to multiple polylogarithms, making use of the partial fractions identity given in [19]

$$\frac{1}{m^a n^b} = \sum_{r=b}^{a+b-1} \frac{\binom{r-1}{b-1}}{(m+n)^r m^{a+b-r}} + \sum_{r=a}^{a+b-1} \frac{\binom{r-1}{a-1}}{(m+n)^r n^{a+b-r}}. \quad (\text{C.1})$$

These reductions have been checked with the program `HyperInt` [23].<sup>8</sup>

We will now reduce the multiple constrained sum  $L_{a,b;c}(x, y)$  that arose in the evaluation of  $C_{1,1,1}(q)$  to a sum of multiple polylogarithms, which is given by

$$L_{a,b;c}(x, y) = \sum_{m_1, m_2 \geq 1} \frac{x^{m_1} y^{m_2}}{m_1^a m_2^b (m_1 + m_2)^c} \quad (\text{C.2})$$

can be expressed, when  $a, b > 0$ , in the form

$$L_{a,b;c}(x, y) = \sum_{\substack{r+s=a+b \\ r, s > 0}} \left( \binom{r-1}{a-1} \text{Li}_{s,c+r}(y/x, x) + \binom{r-1}{b-1} \text{Li}_{s,c+r}(x/y, y) \right). \quad (\text{C.3})$$

Clearly  $L_{a,b;c}(x, y) = L_{b,a;c}(y, x)$  and

$$L_{0,b;c}(x, y) = \text{Li}_{b,c}(y/x, x) \quad (\text{C.4})$$

$$L_{a,b;0}(x, y) = \text{Li}_a(x) \text{Li}_b(y). \quad (\text{C.5})$$

One can integrate the differential equation

$$\left( x \frac{d}{dx} + y \frac{d}{dy} \right)^c L_{a,b;c}(x, y) = \text{Li}_a(x) \text{Li}_b(y). \quad (\text{C.6})$$

to give

$$L_{a,b;c}(x, y) = \frac{1}{(c-1)!} \int_0^\infty d\alpha \alpha^{c-1} \text{Li}_a(xe^{-\alpha}) \text{Li}_b(ye^{-\alpha}). \quad (\text{C.7})$$

This integral representation can then be efficiently integrated with `HyperInt` [23]. The expression can be written in different ways as a consequence of the shuffle algebra. For example

$$L_{1,1,1}(x, y) = \text{Li}_{1,2}\left(\frac{y}{x}, x\right) + \text{Li}_{1,2}\left(\frac{x}{y}, y\right), \quad (\text{C.8})$$

which is easily checked with `HyperInt` to be equal to

$$L_{1,1,1}(x, y) = \text{Li}_3(x) + (\text{Li}_2(x) - \text{Li}_2\left(\frac{x}{y}\right)) \text{Li}_1(y) + \text{Li}_{1,2}\left(y, \frac{x}{y}\right). \quad (\text{C.9})$$

Equating these two expressions evaluated at  $x = y = 1$  implies Euler's famous relation  $\zeta(1, 2) = \zeta(3)$ .

<sup>8</sup> We thank Erik Panzer for help in performing these checks.

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